



CENTRO DE ESTUDIOS DE FÍSICA DEL COSMOS
DE ARAGÓN
(CEFCA)



Imprints on the CMB from the dark ages

Carlos Hernández-Monteagudo

CEFCA

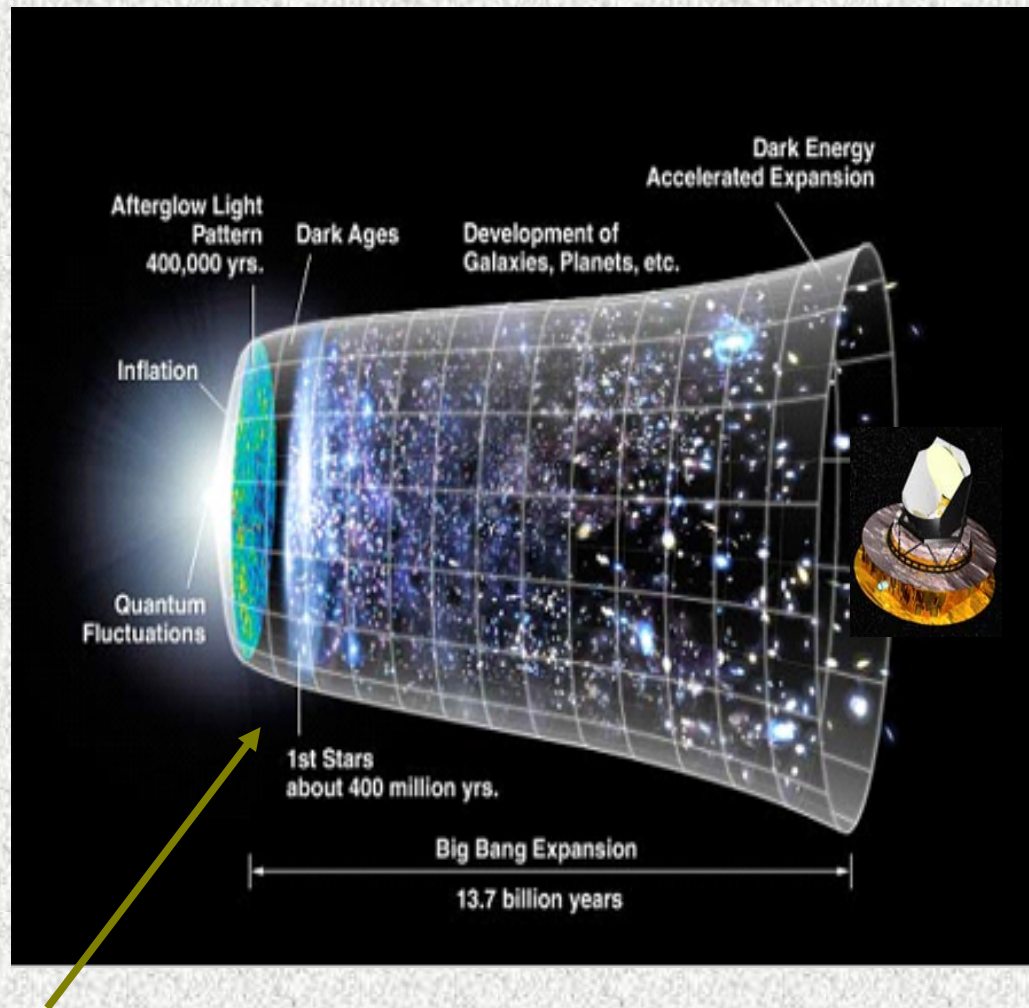


Microwave Spectral
Polarimetry

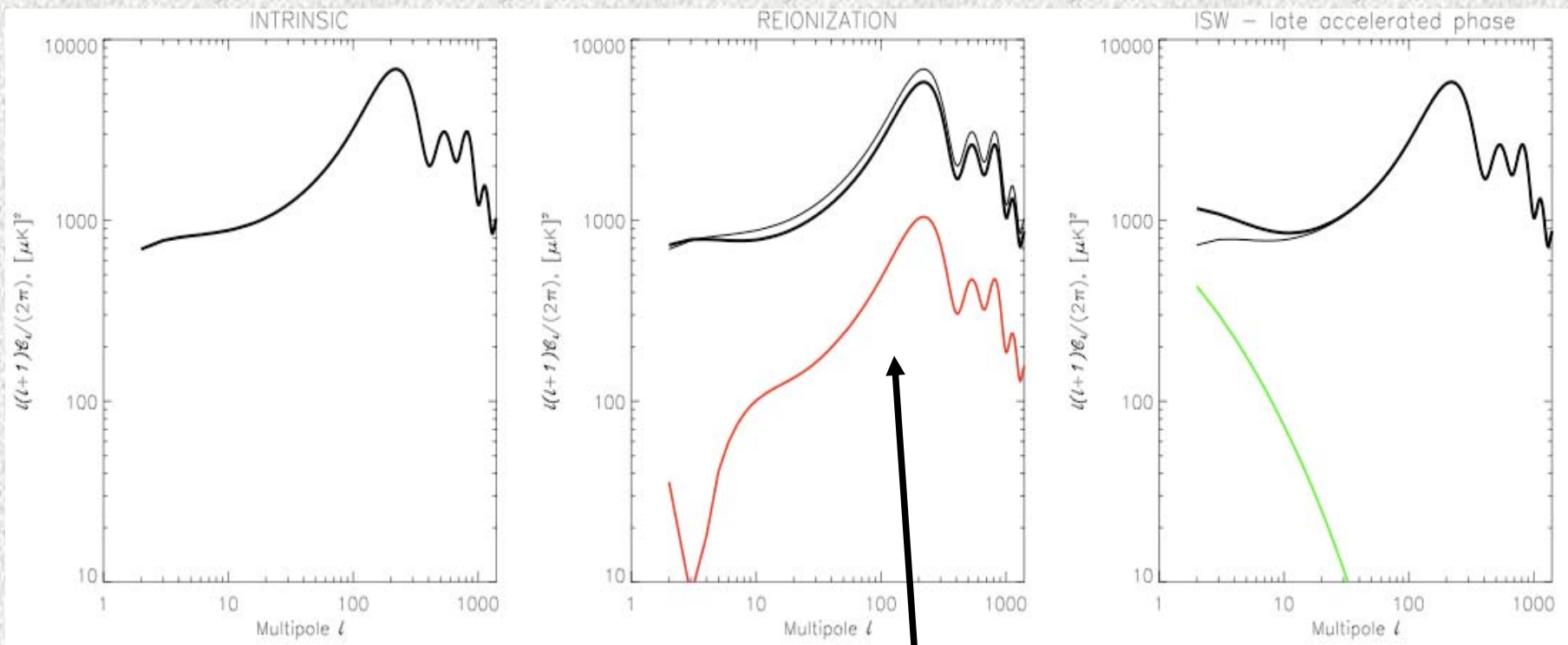
APC, Université Paris 7, Dec. 11th 2012

Outline

- *Standard* CMB secondary anisotropies generated during the end of the dark ages and beginning of reionization
- Physics of the resonant scattering of CMB photons off metallic resonant transitions
- Collisional emission in rotational levels of CO molecules at moderate to high redshift
- The Wouthuysen-Feld effect in OI: UV induced CMB distortion in the 63μ OI transition



Let us focus on *Reionization*...



$$\delta C_l \approx -2 \tau_T C_l$$

Biggest effect is *frequency independent*...

FREQUENCY DEPENDENT SECONDARY EFFECTS DURING REIONIZATION:

Effects preserving the CMB black body spectrum:

- * Resonant scattering on fine structure **lines** of ions, metals and molecules

*Effects **distorting** the CMB black body spectrum:*

- * Collisional emission on fine structure **lines** of ions, metals and molecules
- * Wouthuysen-Field effect (UV pumping) in OI 63.2 micron **transition**
- * Reprocessed UV emission in the microwave range by dust particles

The presence of **LINES** in the CMB:

- If a given species X has a resonant transition with a given resonant frequency, the observed CMB will interact with that species at a redshift

$$\nu_{X_i} / \nu_{obs} = 1 + z_{X_i}$$

- In *multifrequency* experiments like Planck, a low frequency channel can be used as **reference**, since it will probe **highest redshifts** (for which the species **abundance** should be **lowest**)

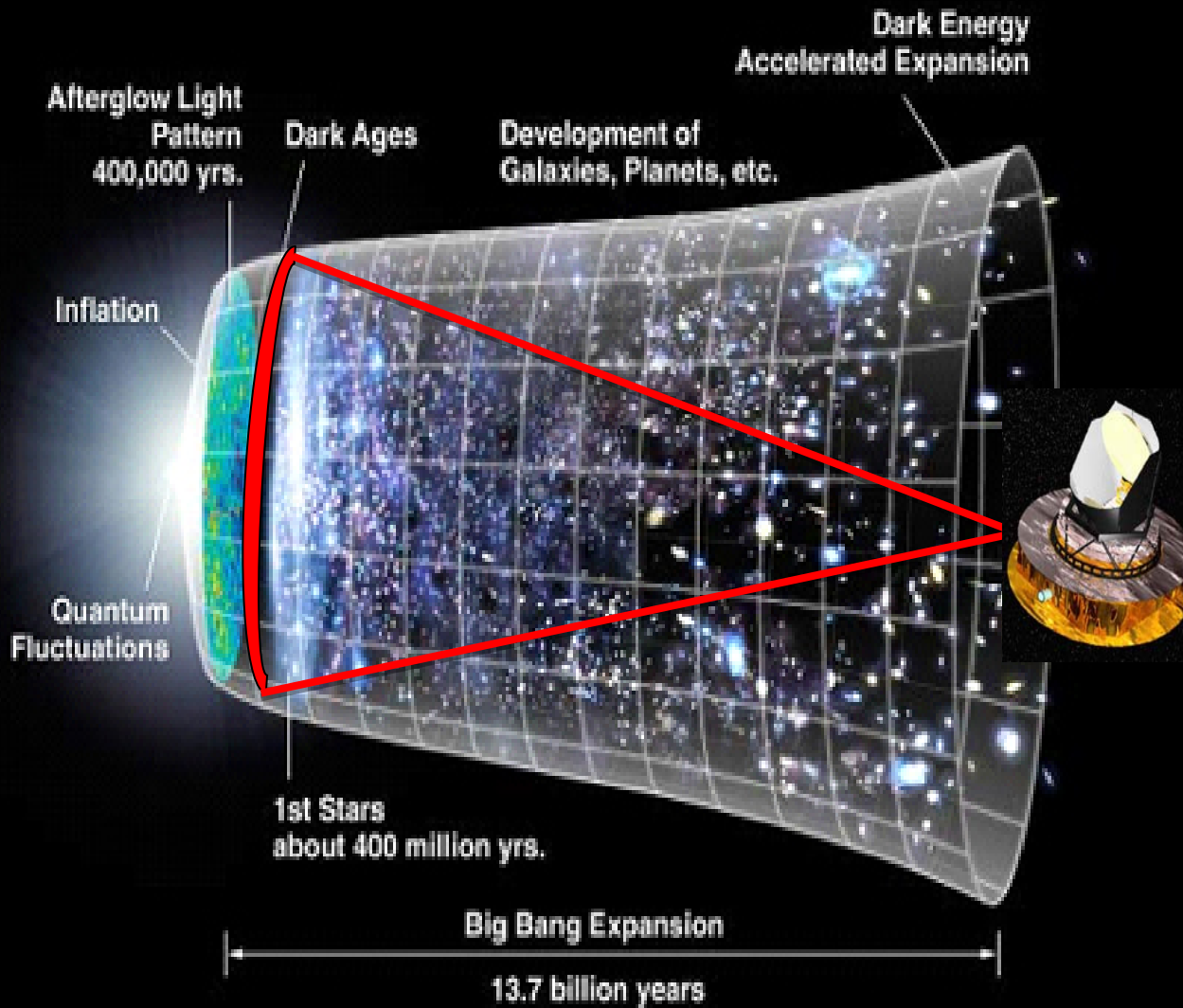
$$\tau_{X_i}(z) = f_i \frac{\pi e^2}{m_e c} \frac{\lambda_i n_{X_i}(z)}{H(z)}$$

- By changing the observing frequency, one can perform **tomography** of the species during reionization

Basu, CHM & Sunyaev (A&A, 2004), CHM & Sunyaev (MNRAS, 2005), CHM, Rubiño-Martín & Sunyaev (A&A 2005), CHM, Verde & Jimenez (ApJ, 2006)

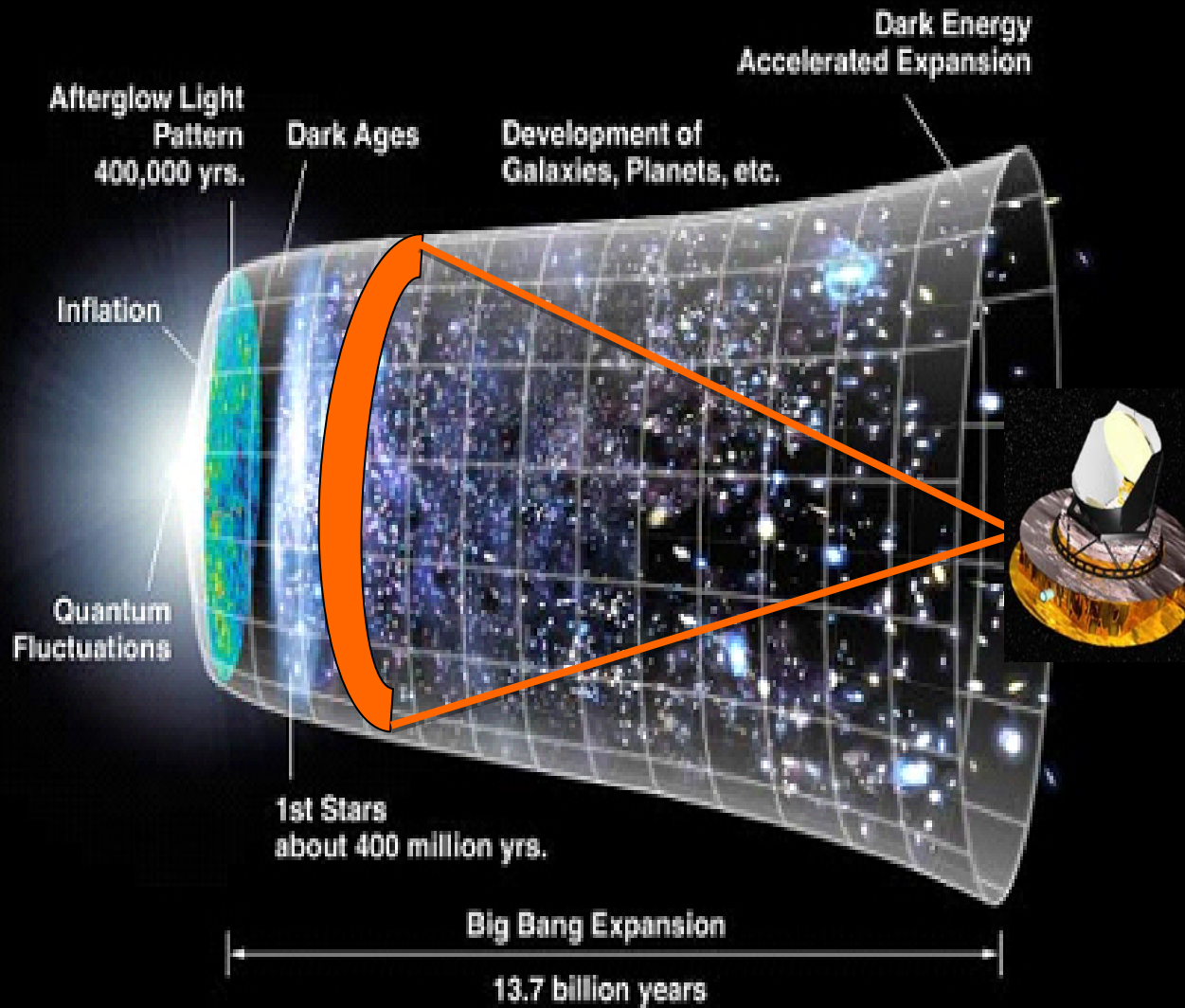
$$\nu_{X_i} / \nu_{obs} = 1 + z_{X_i}$$

@ 100 GHz



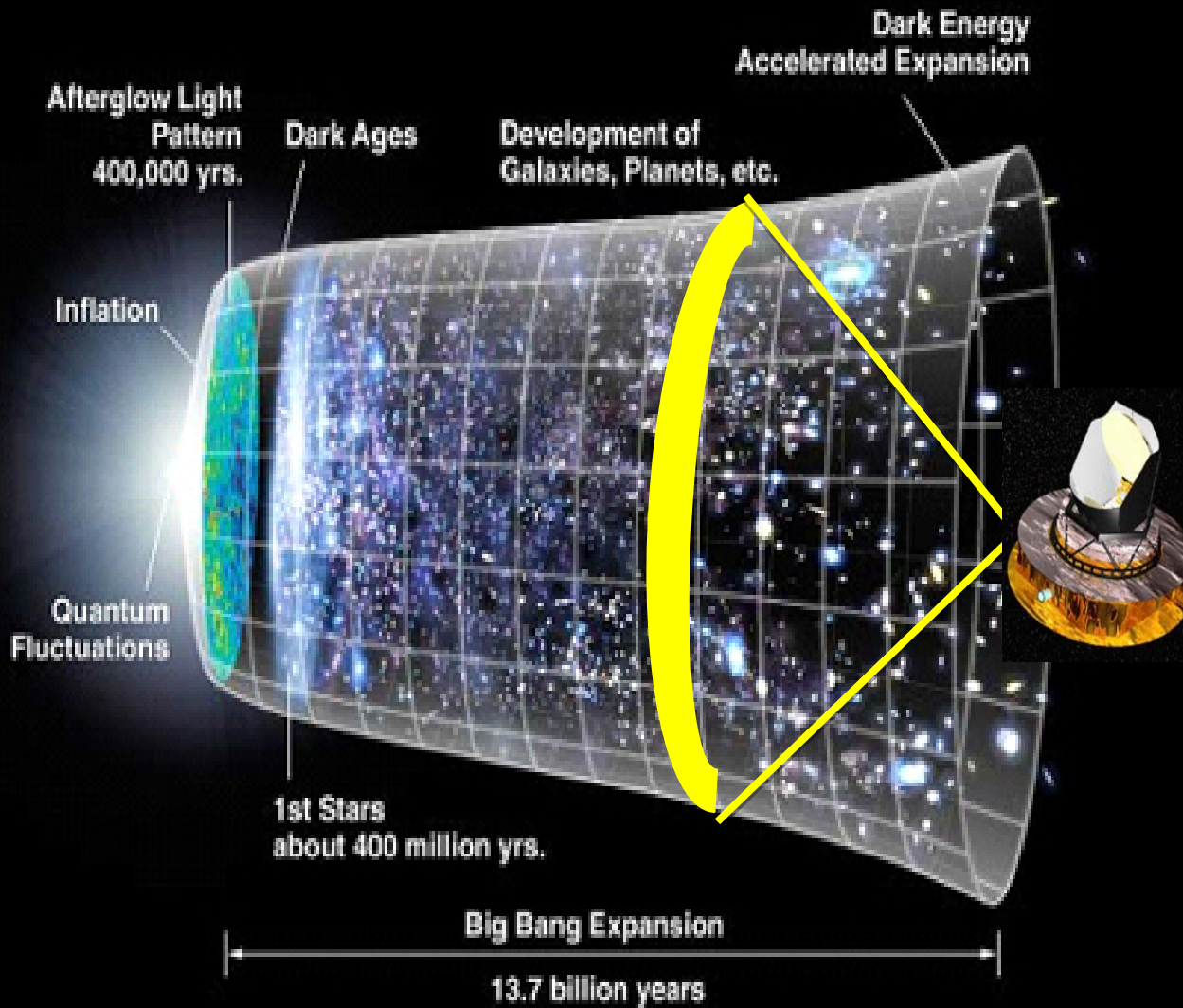
$$\nu_{X_i} / \nu_{obs} = 1 + z_{X_i}$$

@ 143 GHz



$$\nu_{X_i} / \nu_{obs} = 1 + z_{X_i}$$

@ 217 GHz



End of the Dark Ages and beginning of Reionization

Imprint of metals and molecules synthesized during reionization on the CMB via *frequency dependent signals*:

- **Resonant scattering** on fine structure lines of metal and ionic species
- **Collisional emission** on those same lines and species + on molecules like CO
- **UV pumping** on the Balmer-alpha line of OI at 63.2 micron

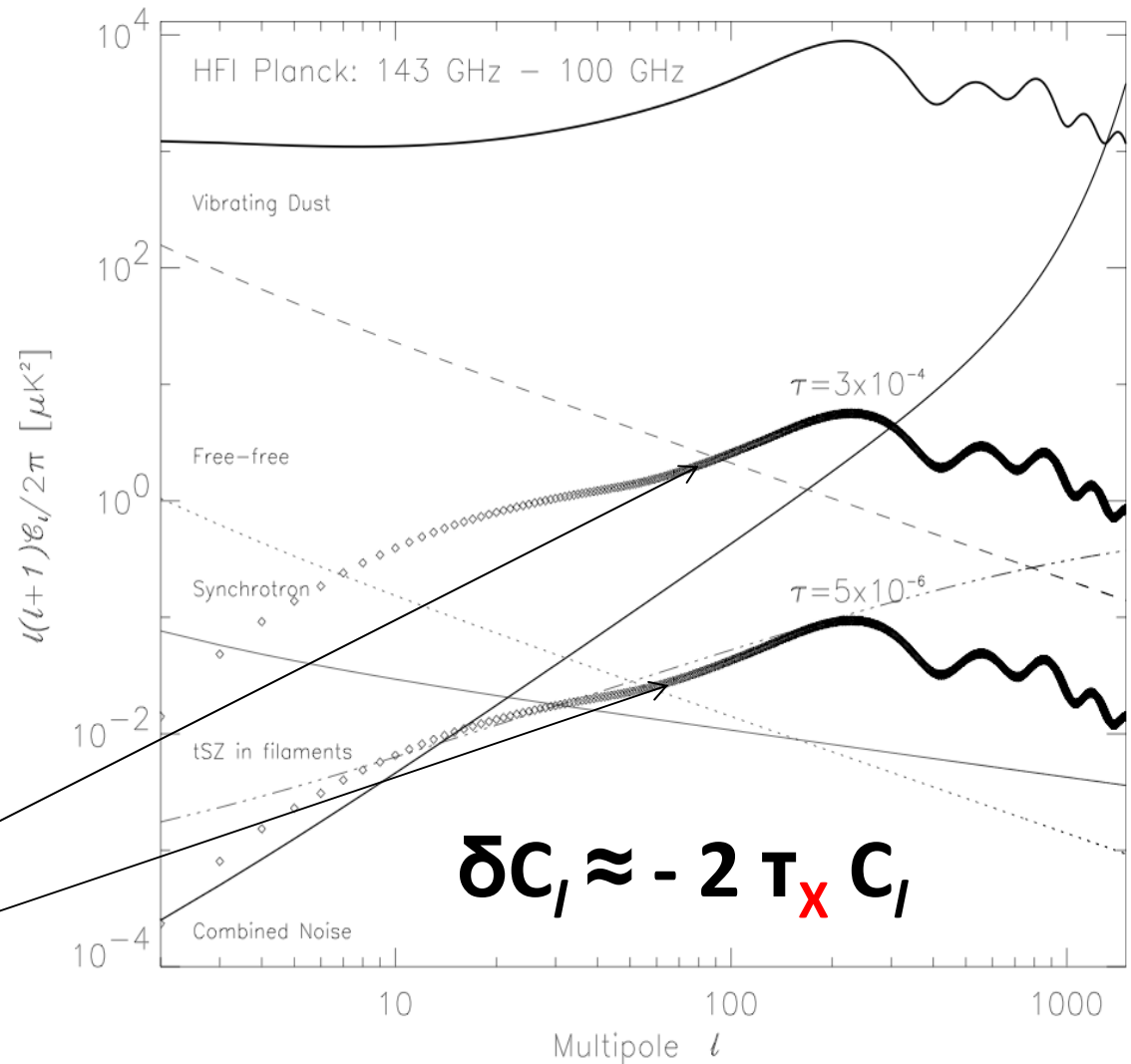
Resonant scattering

IMPACT OF **RESONANT SCATTERING** ON CMB POWER SPECTRUM

The biggest effect is a mere **frequency dependeng blurring** of the intrinsic CMB intensity anisotropies !! (almost like in CMB reionization ...)

$$\tau_{X_i}(z) = f_i \frac{\pi e^2}{m_e c} \frac{\lambda_i n_{X_i}(z)}{H(z)}$$

Basu, CHM & Sunyaev (04)



$$\delta C_l \approx - 2 \tau_X C_l$$

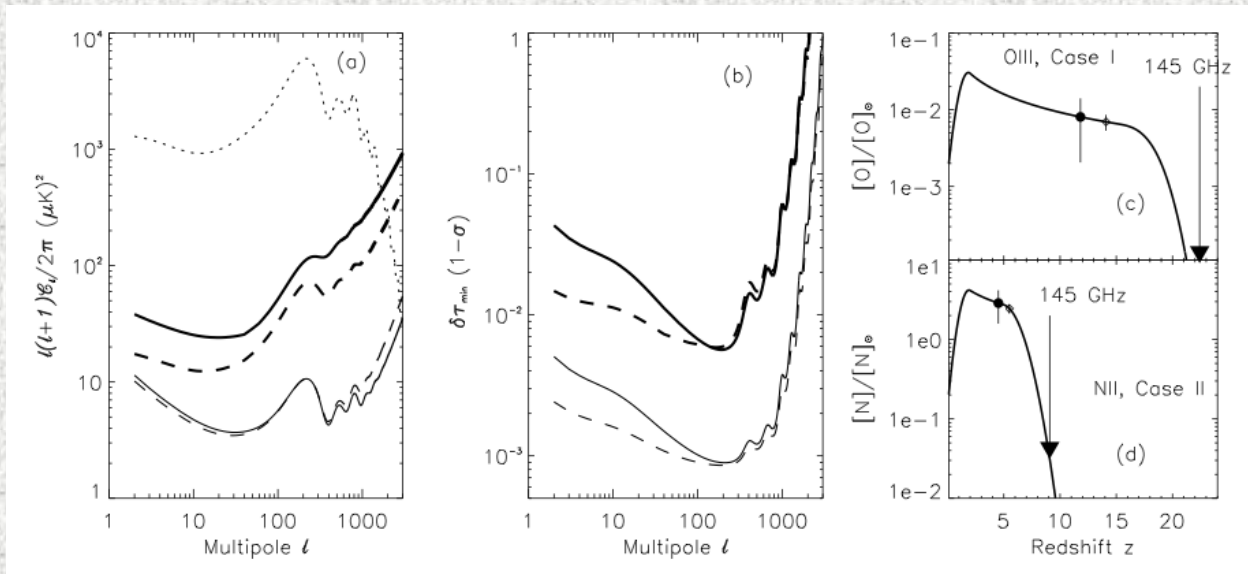
Change in angular power spectrum (C_l) ~ *Optical depth in the line*

~ Abundance of metals at resonant redshift ~ **PLANCK will set constraints**

on metal/ion abundances at very high redshift ($z \sim 8-50$) with **CMB observations**

IMPACT OF **RESONANT SCATTERING** ON CMB POWER SPECTRUM

Technical requirements



$$\tau_{X_i}(z) = f_i \frac{\pi e^2}{m_e c} \frac{\lambda_i n_{X_i}(z)}{H(z)}$$

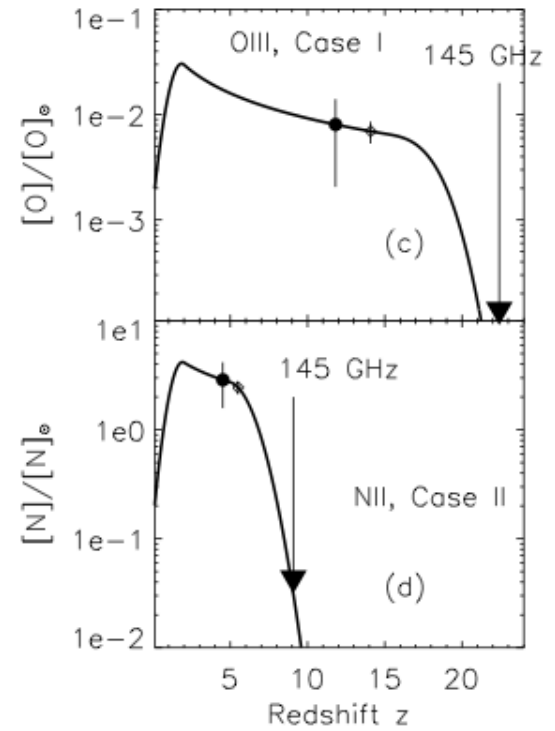
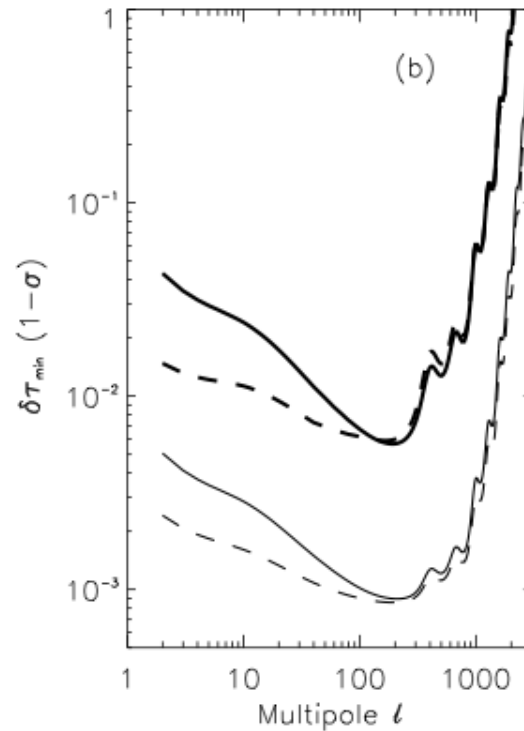
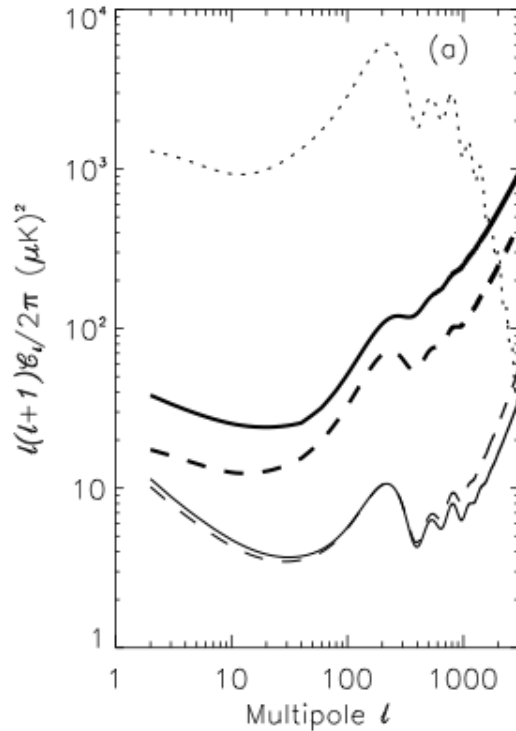
$$a_{lm}^j \equiv \frac{\tilde{a}_{lm}^j}{(W_{lm}^j)_{\text{model}}} = f_j \left[B_{lm}^j \left(a_{lm}^{\text{CMB}} + g_{\text{tSZ},j} a_{lm}^{\text{tSZ}} + \sum_M g_{M,j} a_{lm}^M + \sum_F g_{F,j} a_{lm}^F \right) + N_{lm}^j \right]$$

$$\begin{aligned} |a_{lm}^j|^2 - |a_{lm}^i|^2 &\simeq 2 |a_{lm}^{\text{CMB}}|^2 \left(\text{Re}(\delta f_{ji} + \delta \epsilon_{lm}^{ji}) - \delta \tau_{\text{eff}}^{ji} \right) \\ &+ |a_{lm}^{\text{tSZ}}|^2 \Delta[(g_{\text{tSZ}}^{ji})^2] + \sum_F |a_{lm}^F|^2 \Delta[(g_F^{ji})^2] \\ &+ 2 \sum_{F_u > F_s} \text{Re} [a_{lm}^{F_u} (a_{lm}^{F_s})^*] \Delta(g_{F_u}^{ji} g_{F_s}^{ji}) \\ &+ |N_j|^2 - |N_i|^2 + 2 \text{Re} [a_{lm}^{\text{CMB}} (a_{lm}^{\text{tSZ}})^*] \Delta g_{\text{tSZ}}^{ji} \\ &+ 2 \sum_F \text{Re} [a_{lm}^{\text{CMB}} (a_{lm}^F)^*] \Delta g_F^{ji} \end{aligned}$$

CHM, Verde & Jiménez (06)

Constraints on τ are mostly limited by uncertainties in the **inter-channel calibration and PSF characterisation**

$$\tau_{X_i}(z) = f_i \frac{\pi e^2 \lambda_i n_{X_i}(z)}{m_e c H(z)}$$



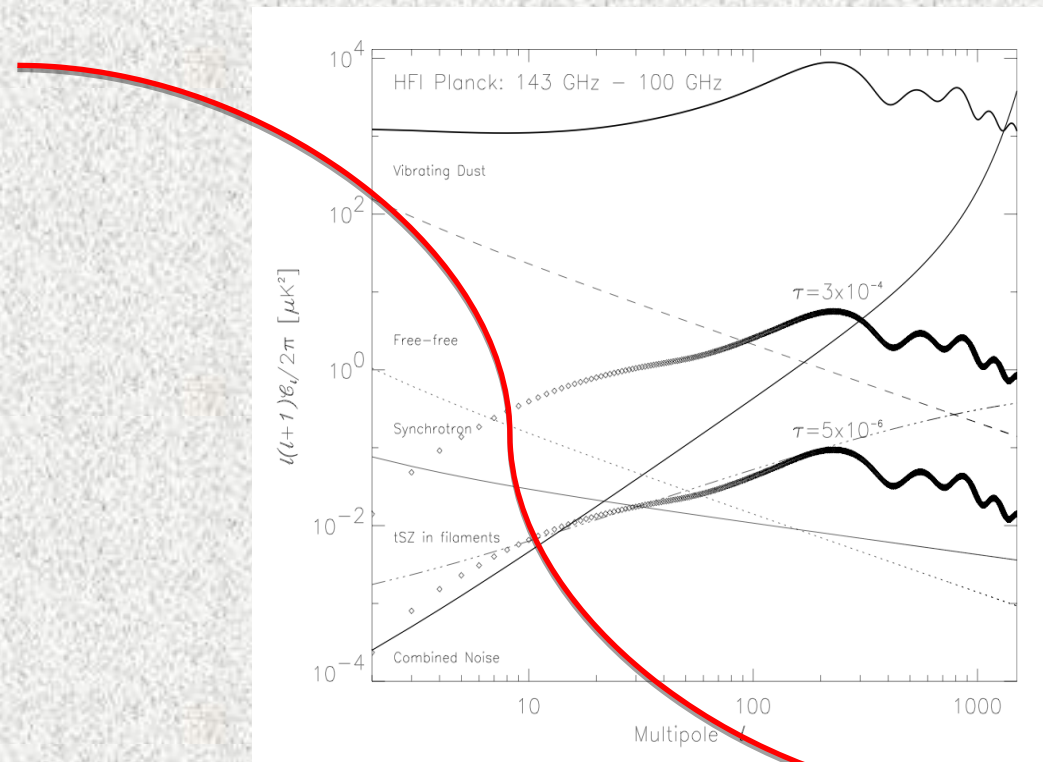
There are signatures on **polarization** as well, see talk by José Alberto Rubiño-Martín tomorrow ...

CHM, Verde & Jiménez (06)

IMPACT OF **RESONANT SCATTERING** ON CMB POWER SPECTRUM

Second order effects

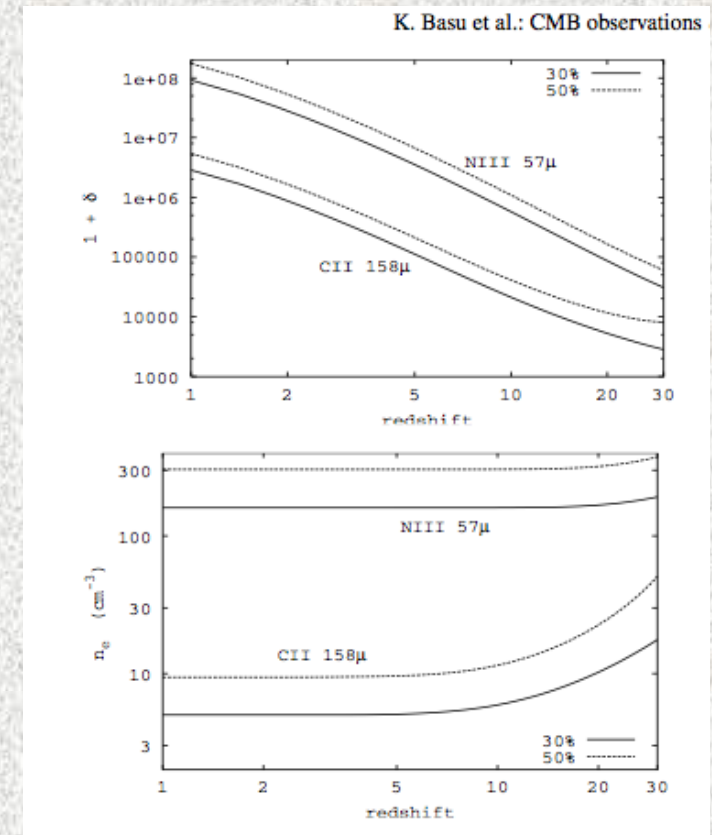
- The **clustering** of bubbles containing metals and the **peculiar motions** of the scatterers introduce **second order** anisotropies
- However, these are **not linear** in the **optical depth (τ)**, but **quadratic (τ^2)**, thus yielding much smaller amplitudes (just like for *patchy reionization signatures in the high- l regime....*)



Collisional emission

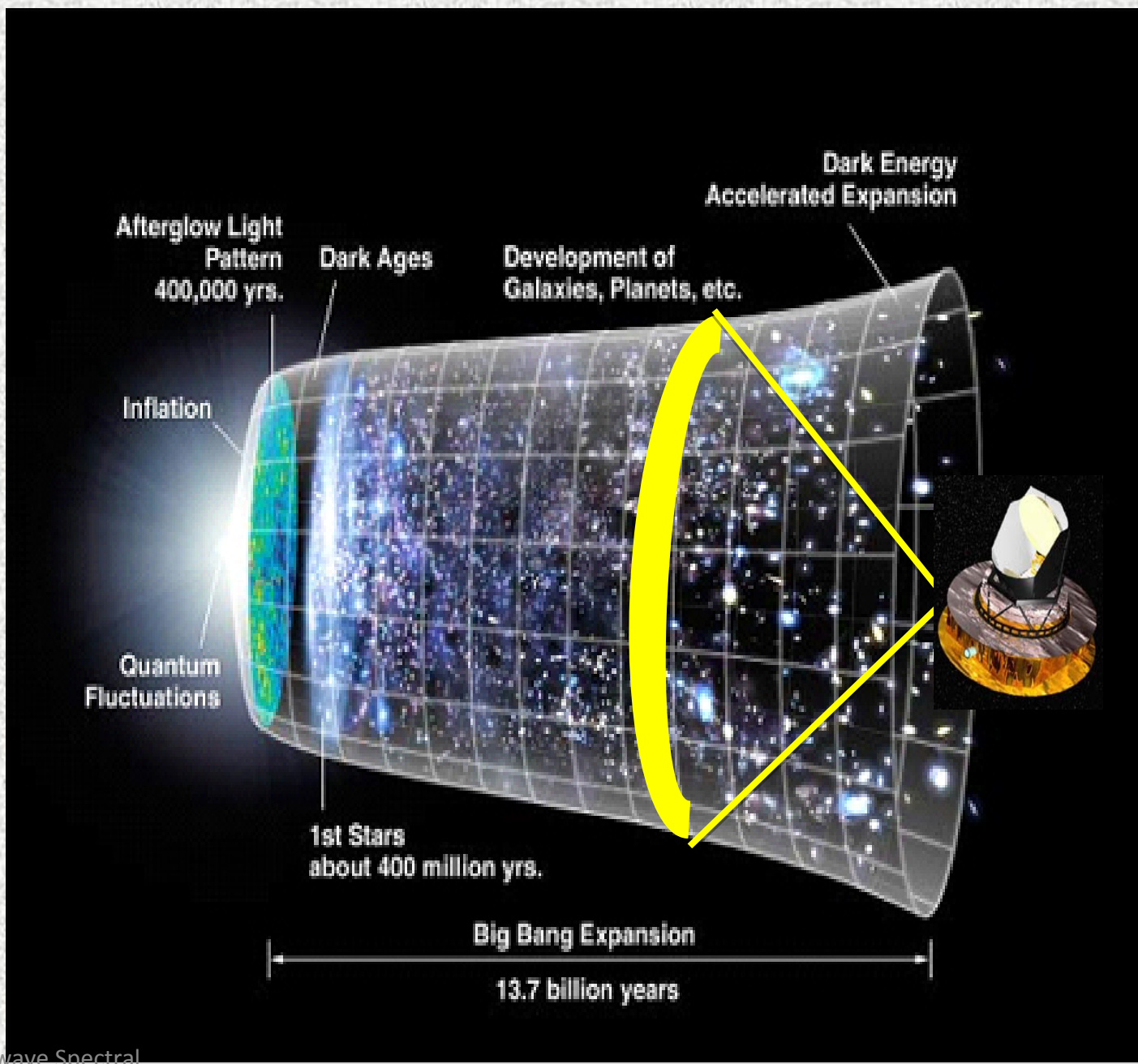
Collisionally induced **emission** in dense regions: the clustering of star forming regions

- One needs to go to **highly overdense** regions in order to change significantly the lower level population in the fine structure transitions of interest and hence modify the optical depth ($\delta \sim 1e4 -- 1e7$)
- In all other regions the **resonant scattering** is the dominant process
- However, the anisotropy pattern of line emission allows playing with the **radial coordinate – frequency resolution**

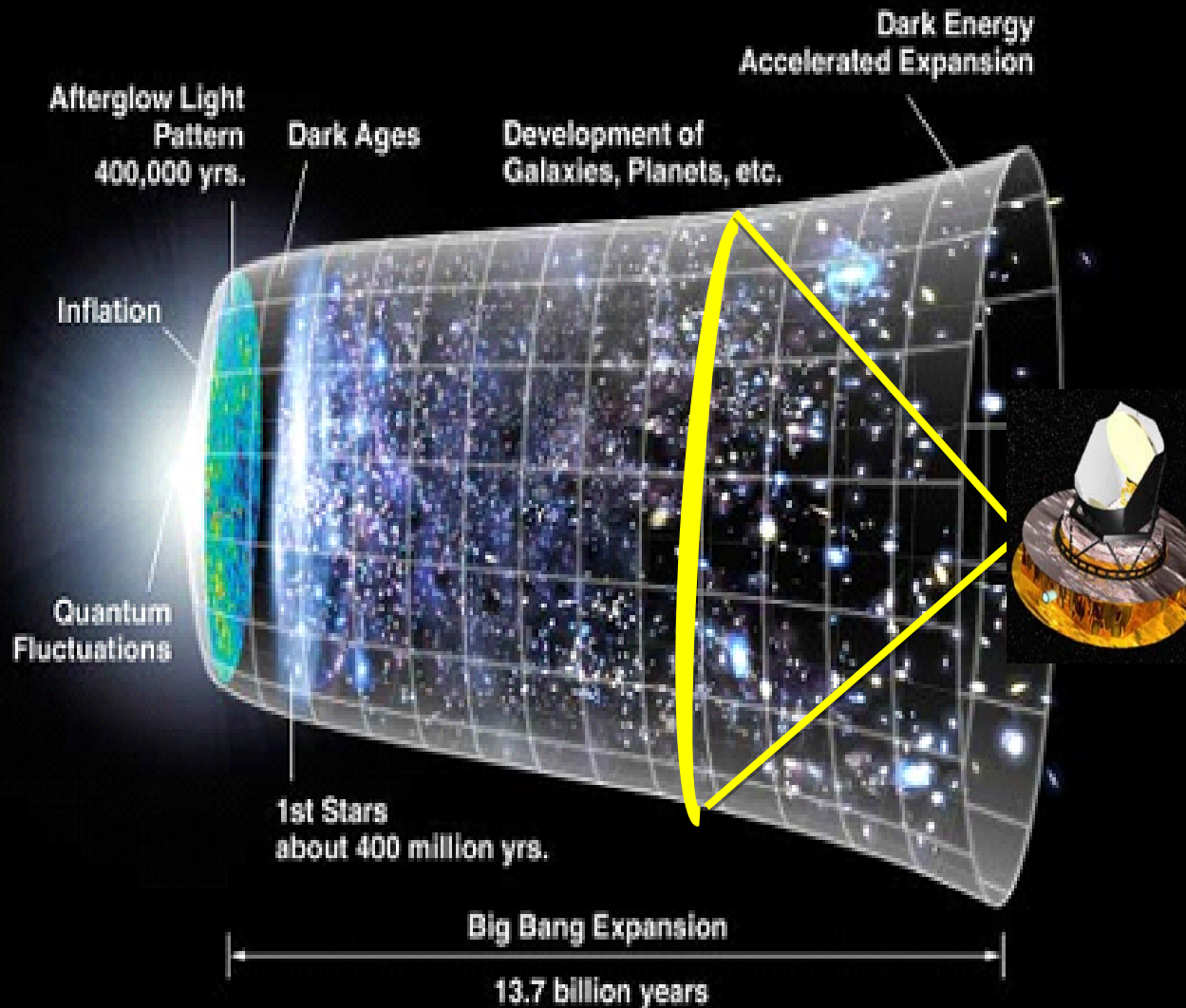


Basu, CHM & Sunyaev, (A&A 2004)

$$\frac{\Delta\nu}{\nu_{obs}} = 0.2$$



$$\frac{\Delta\nu}{\nu_{obs}} = 10^{-3}$$

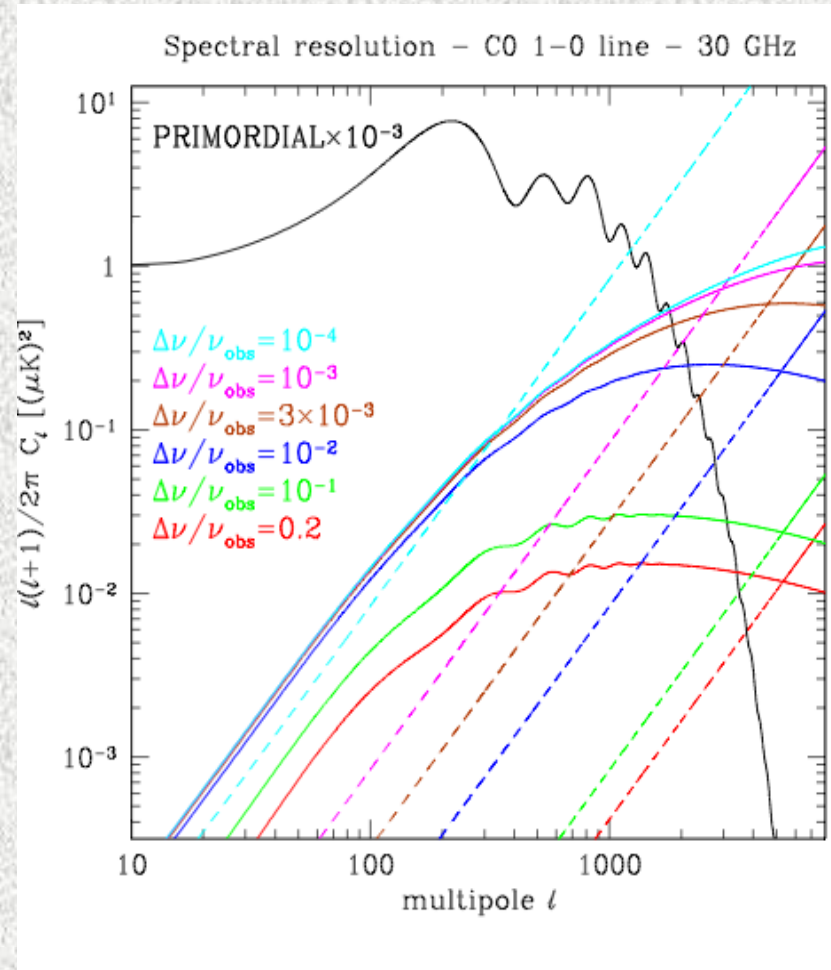


Collisionally induced **emission** in dense regions: the clustering of star forming regions

- The *spectral width* of the observing experiment conditions the *width of the redshift shell* interacting with the CMB: **this allows projecting out all “non-line” foregrounds (including CMB!)**

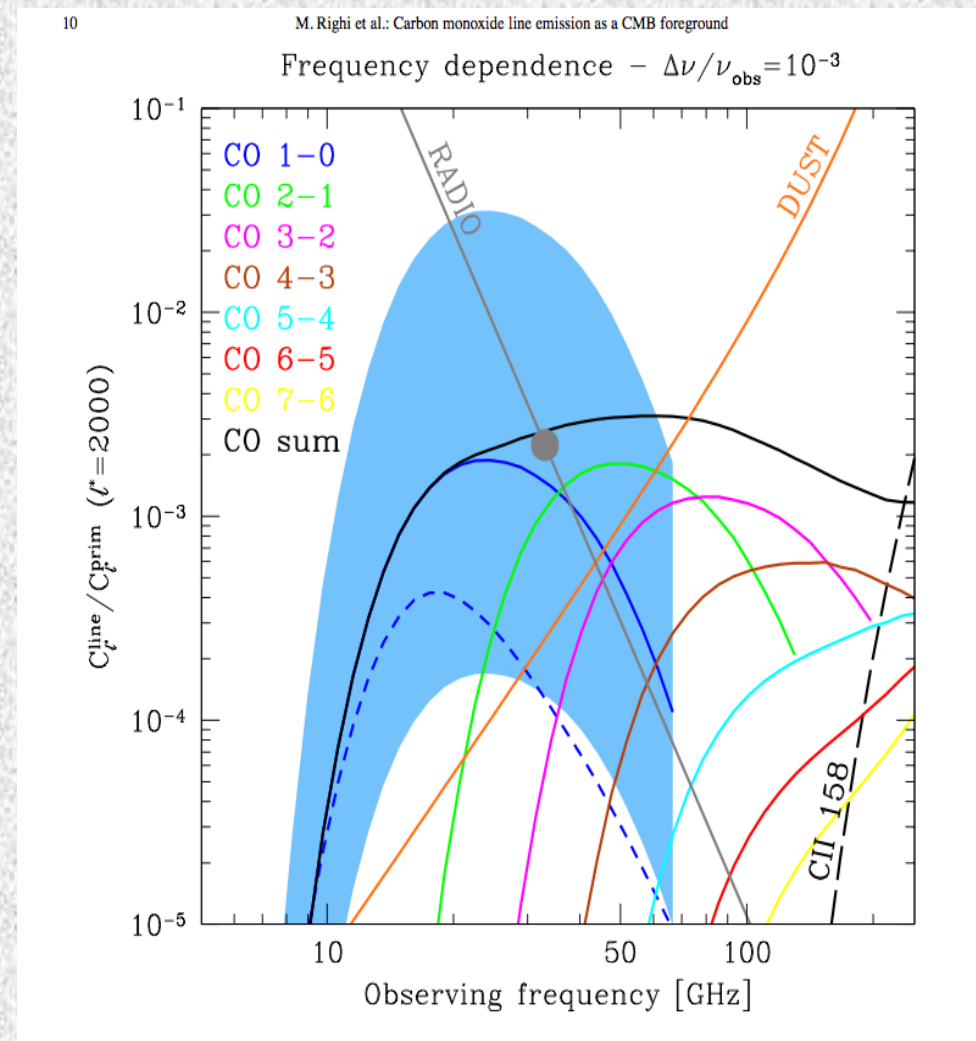
Collision induced CO 1- \rightarrow 0 line emission at 30 GHz and resulting anisotropy, after implementing a Kennicutt-law type relation between CO luminosity and Star Formation Rate in nearby galaxies and extrapolate it to higher-z using a Lacey & Cole approach

Righi, CHM & Sunyaev, (A&A 2008)



For relative spectral width of $1e-3$, the CO J=1-0 115 GHz is the **largest** foreground at $\ell \sim 2,000$ between 30 – 70 GHz

Future NIR and FIR surveys may be able to detect the presence of these collisions at moderate to high redshifts **on different transitions in OI, NII, OIII ...**



Righi, CHM & Sunyaev, (A&A 2008)

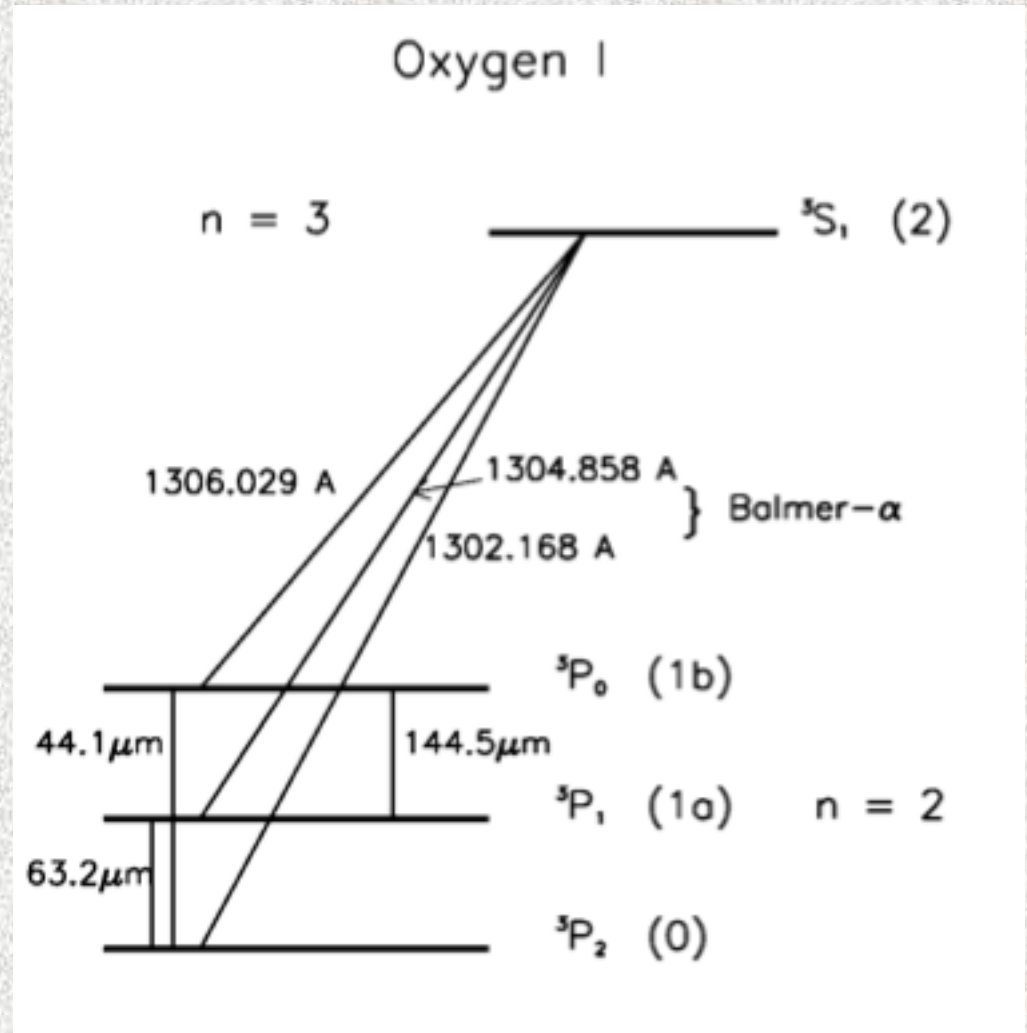
Wouthuysen-Feld effect on OI

The **Wouthuysen-Feld effect** or **UV pumping** in OI and the resulting induced distortion on the CMB

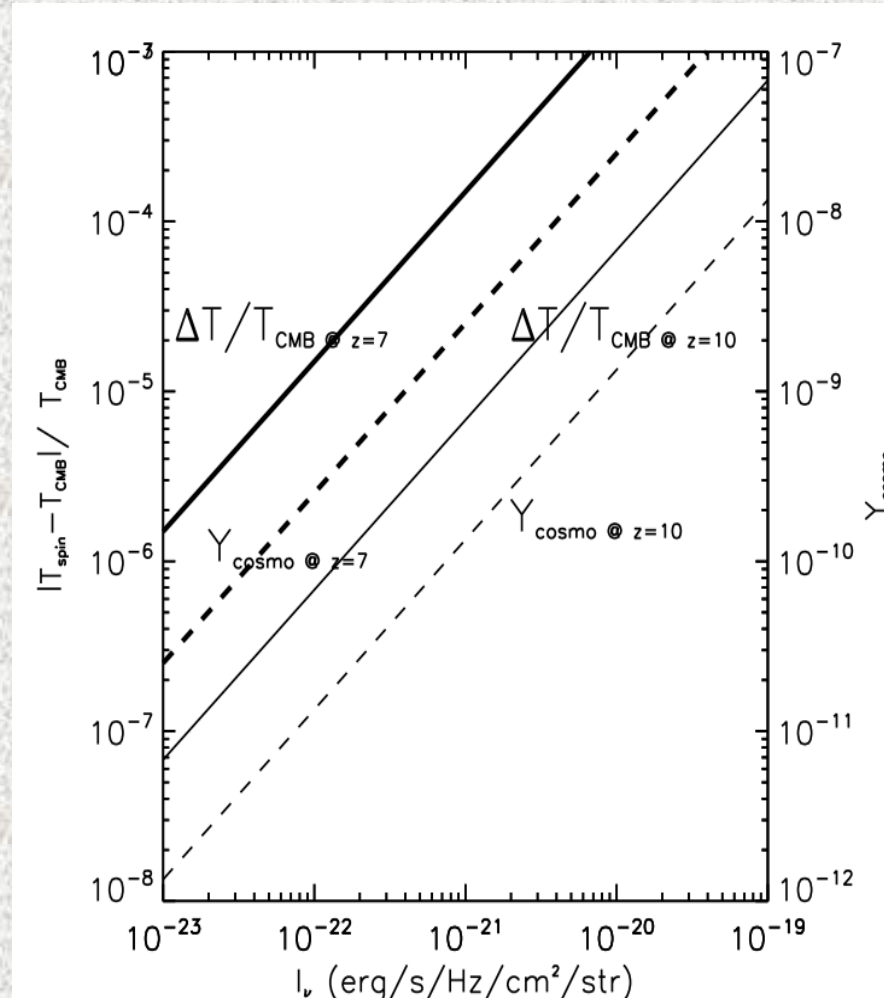
UV radiation pumps systems from (0-1a-1b) to (2), **distorting the CMB-induced equilibrium** between the (0-1a-1b) levels

$$\frac{n_j}{n_i} = \frac{P_{ij}^{UV} + B_{ij}I_\nu}{P_{ji}^{UV} + B_{ji}I_\nu + A_{ji}} \equiv \frac{g_j}{g_i} \exp\left(\frac{-T_{*,j}}{T_{S,j}}\right).$$

CHM, Haiman, Verde & Jimenez, ApJLetters 2007

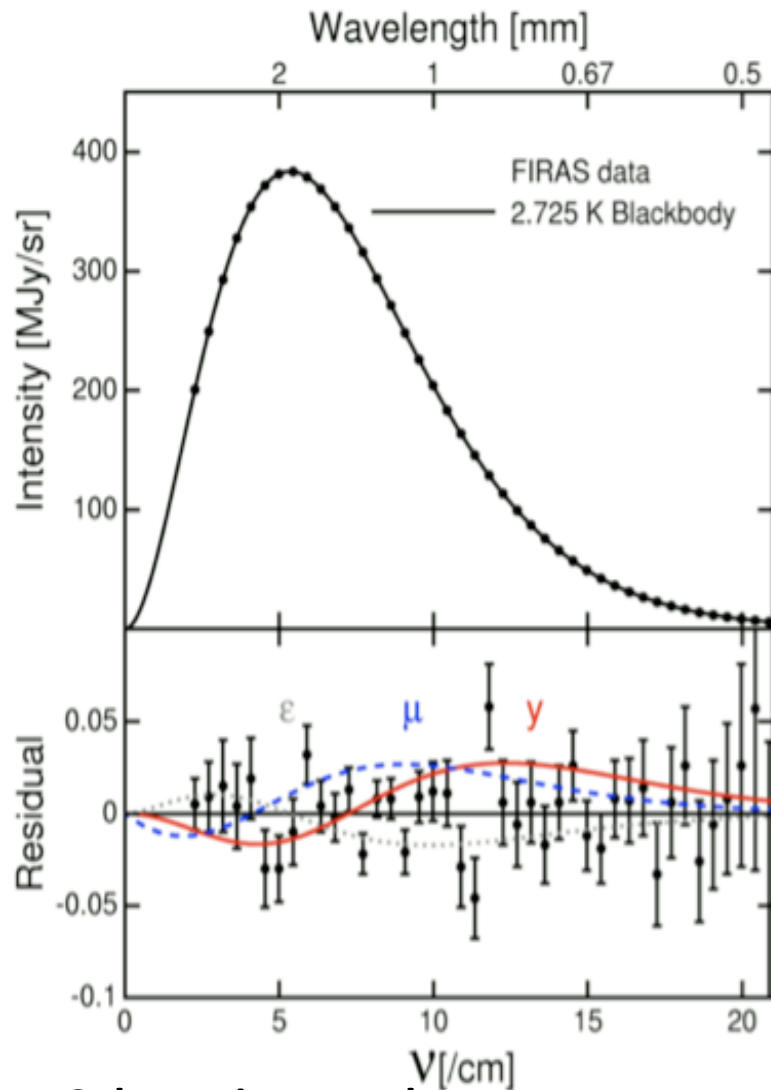


UV pumping induced distortion of the CMB



CHM, Haiman, Verde & Jimenez, ApJLetters 2007

UV background generated by first stars changes T_{spin} in fine structure transitions of **neutral oxygen** OI (64micron, 148micron) ~ **distortion** in CMB spectrum and **temperature anisotropies** ~ **constraints** on **neutral oxygen** during reionization.



FIRAS data, Fixen et al.

- FIRAS data can already constrain OI abundance to be less than $5 - 40 Z_{\odot}$ at $z > 10$ \Rightarrow **first constraints on OI at reionization!**
- Future experiments (improved versions of FIRAS) should be able to set tight constraints on OI during the reionization epoch.

CHM, Haiman, Verde & Jimenez, ApJLetters 2007

What are the effects of OI clustering ??

$$y_{\text{eff}}(\nu_{\text{obs}}, \hat{n}) \approx \int dM \left[B * \left(\frac{dn}{dM} * W_b \right) \right] \hat{y}.$$

As OI gets expelled in **bubbles** from SN winds, these grow in size ... the total contribution along a given line of sight must account for the **halo mass function** (giving rise to the bubbles) and the **instrumental frequency** and **angular responses**

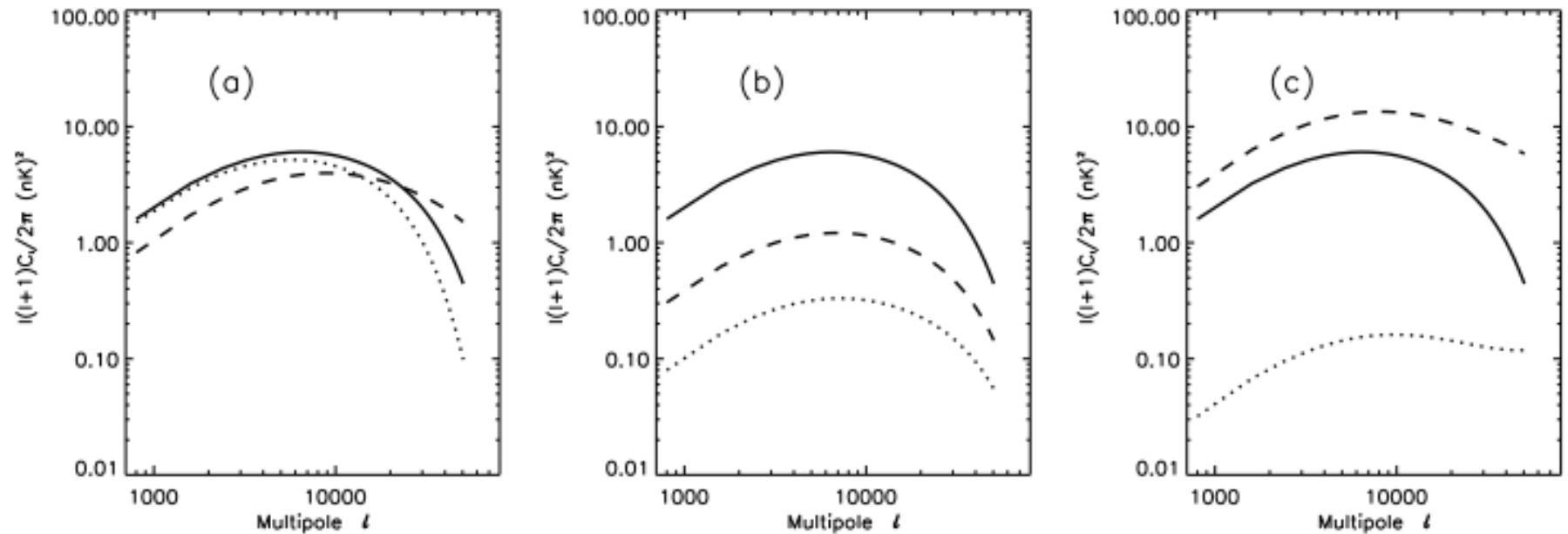
Since halos are **correlated**, so these bubbles will, and this can be computed ...

What are the effects of OI clustering ??

$$y_{\text{eff}}(\nu_{\text{obs}}, \hat{\mathbf{n}}) \approx \int dM \left[B * \left(\frac{dn}{dM} * W_b \right) \right] \hat{y}.$$

$$\langle y_{\text{eff}}(\hat{\mathbf{n}}_1) y_{\text{eff}}(\hat{\mathbf{n}}_2) \rangle = \int dM_1 dM_2 \langle \tilde{n}_1 \tilde{n}_2 \rangle \hat{y}_1 \hat{y}_2$$

$$= \int dM_1 dM_2 \tilde{n}_1 \tilde{n}_2 \hat{y}_1 \hat{y}_2 [1 + \tilde{\xi}_{hh}(\mathbf{r}_1, M_1, \mathbf{r}_2, M_2)]$$



CHM, Haiman, Verde & Jiménez (08)

At the nK level ... ALMA??

Summary

- Metals at the end of the Dark Ages can interact with CMB photons via *resonant scattering, collisional emission* and the *Wouthuysen-Feld effect*
- The **most frequent** mechanism is the **resonant scattering**, which affect all species outside the most overdense regions. ***Constraints on metal production during the dark ages may be obtained after accurate inter channel calibration and PSF characterisation***
- At the same time, **collisional emission** can be separated from other dominant signals due to **its dependence with the observing spectral resolution**
- Distortion of the CMB due to the **Wouthuysen-Feld effect on OI** lies at the **nK level** and may provide an **independent way** to set constraints on OI at the end of reionization by upcoming measurements of CMB distortions.