

Neutrinos as Probes of New Physics

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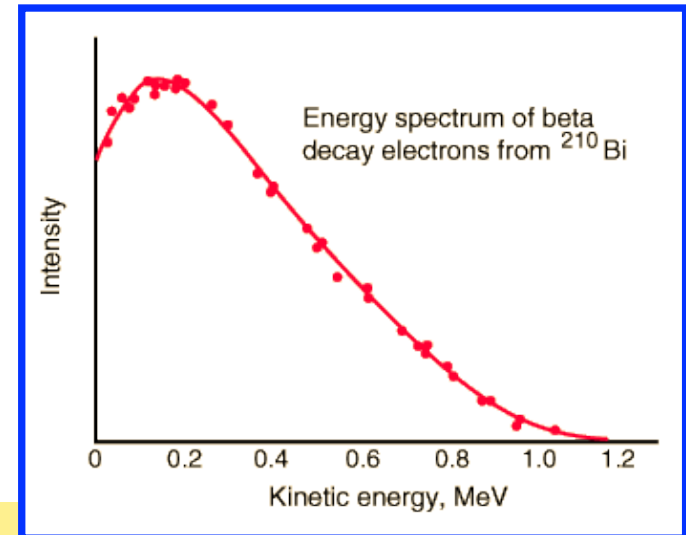


The Birth of the Neutrino



energy-momentum conservation:

- postulate new particle
- invisible, since $Q=0$
- spin $\frac{1}{2}$, ...



W. Pauli: Letter to DPG meeting in Tübingen

Abschrift/15.12.56

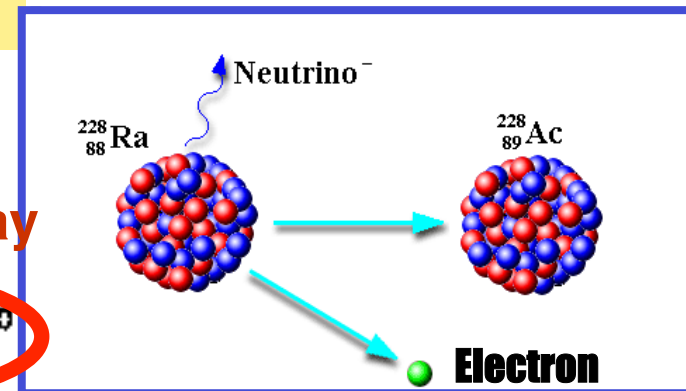
Offener Brief an die Gruppe der Radioaktiven bei der Gauvereins-Tagung zu Tübingen.

Abschrift

Physikalisches Institut
der Eidg. Technischen Hochschule
Zürich

happy 80th birthday

Zürich, 4. Dez. 1930
Usterstrasse

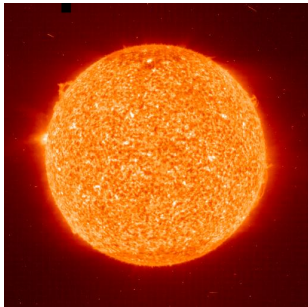


... never detectable!

Liebe Radioaktive Damen und Herren,

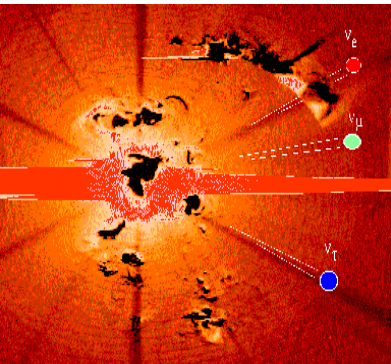
Wie der Überbringer dieser Zeilen, den ich halbvollst anhören bitte, Ihnen das näherem auseinandersetzen wird, bin ich angesichts der "falschen" Statistik der N - und Li-6 Kerne, sowie des kontinuierlichen beta-Spektrums auf einen verweifelten Ausweg verfallen um den "Wechselst" (1) der Statistik und den Energiesatz zu retten. Mithin die Mangelhaftigkeit an Momenten elektrisch neutralen

New Physics: Neutrino Sources



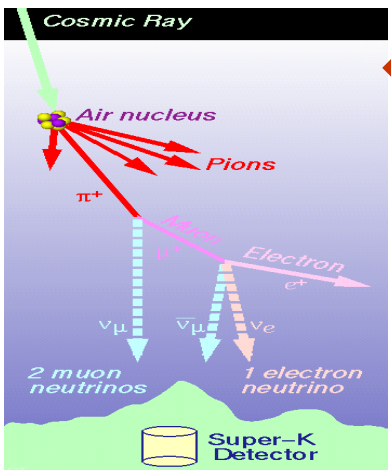
← Sun

Astronomy: →
Supernovae
GRBs
UHE ν 's



← **Cosmology**

Reactors →

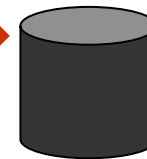


← **Atmosphere**

Accelerators →



β -Sources →



← **Earth**

Fermion Mass Terms in SM

SM: Fermion masses via Yukawa Couplings

Field	$SU(3)_C$	$SU(2)_L$	$U(1)_Y$
$L_Q = \begin{pmatrix} l_u \\ l_d \end{pmatrix}$	3	2	1/3
r_u	3	1	4/3
r_d	3	1	-2/3
$L_L = \begin{pmatrix} l_\nu \\ l_e \end{pmatrix}$	1	2	-1
$r_\nu ???$	1	1	0
r_e	1	1	-2

← ???

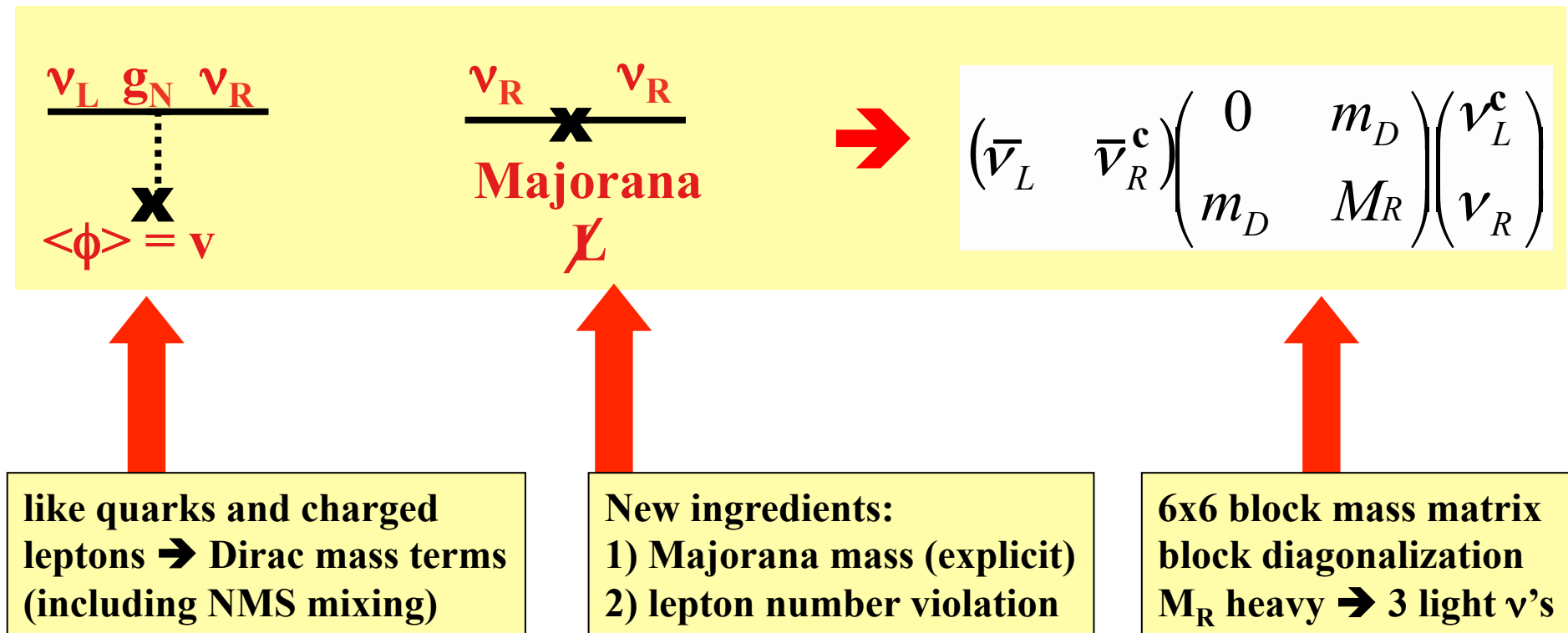
Mass terms: $\overline{L} r + r L$

$L=2, r=1 \rightarrow$ no singlet mass terms

$\Phi=2 \rightarrow \overline{L} \Phi r + r \Phi L \rightarrow$ singlet mass via $\langle \Phi \rangle = v$

New Physics: Neutrino Mass Terms

1) Simplest possibility: add 3 right handed neutrino fields



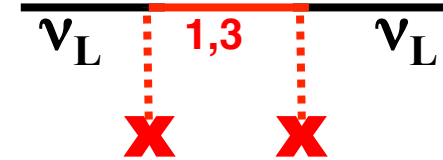
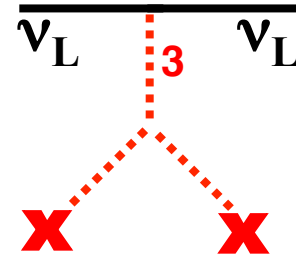
NEW ingredients, 9 parameters \rightarrow SM+

2) Maybe 3+N right handed neutrino fields

→ (6+N) x (6+N) mass matrix

→ how many of the 6+N eigenvalues are light (also for N=0)

3) new: scalar triplets (3_L) or fermionic 1_L or 3_L



→ left-handed Majorana mass term:

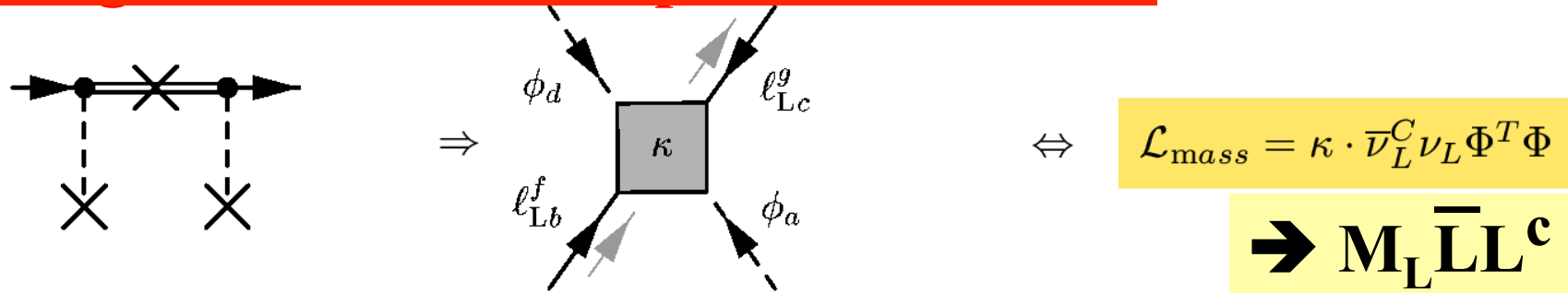
$$\rightarrow M_L \bar{L} L^c$$

4) Both ν_R and new singlets / triplets:

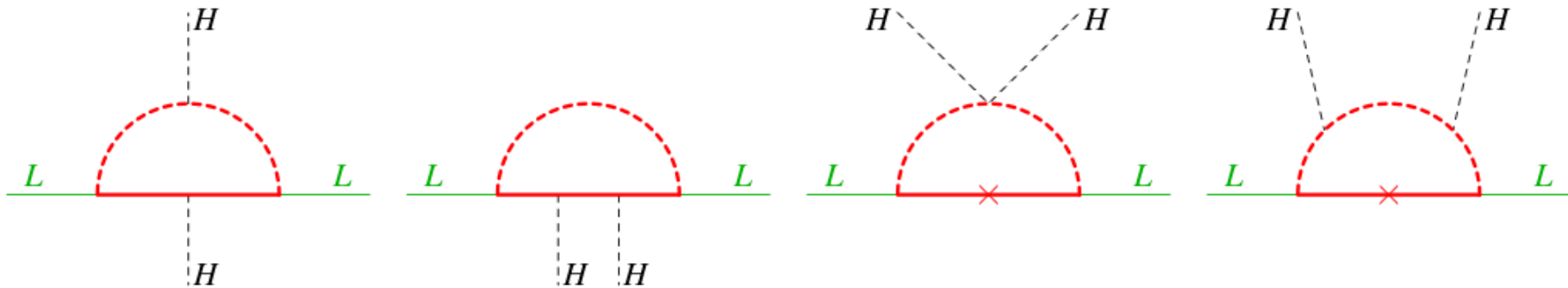
→ see-saw type II, III

$$m_\nu = M_L - m_D M_R^{-1} m_D^T$$

5) Higher dimensional operators: $d=5, \dots$



6) Radiative neutrino mass generation



7-N) SUSY, extra dimensions, ...

\rightarrow so many options...

Other effective Operators Beyond the SM

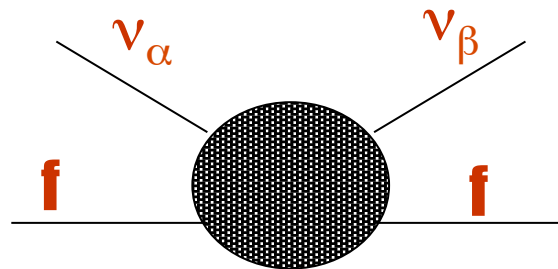
→ effects beyond 3 flavours

→ **Non Standard Interactions = NSIs** → effective 4f operators

$$\mathcal{L}_{NSI} \simeq \epsilon_{\alpha\beta} 2\sqrt{2}G_F (\bar{\nu}_{L\beta} \gamma^\rho \nu_{L\alpha}) (\bar{f}_L \gamma_\rho f_L)$$

• **integrating out heavy physics (c.f. $G_F \leftrightarrow M_W$)**

$$|\epsilon| \simeq \frac{M_W^2}{M_{NSI}^2}$$



Grossman, Bergmann+Grossman, Ota+Sato, Honda et al., Friedland+Lunardini, Blennow +Ohlsson+Skrotzki, Huber+Valle, Huber+Schwetz+Valle, Campanelli+Romanino, Bueno et al., Barranco+Miranda+Rashba, Kopp+ML+Ota, ...

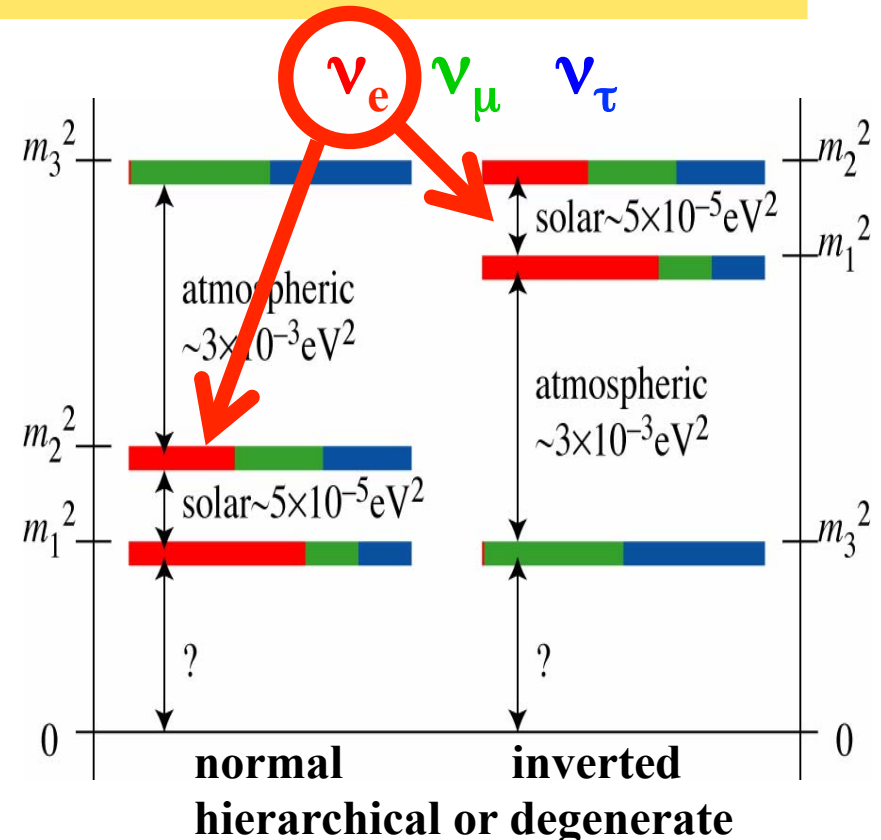
3 Light Neutrinos (...assumed)

Mass & mixing parameters: m_1 , Δm_{21}^2 , $|\Delta m_{31}^2|$, $\text{sign}(\Delta m_{31}^2)$

$$U = \begin{pmatrix} c_{12}c_{13} & s_{12}c_{13} & s_{13}e^{-i\delta} \\ -s_{12}c_{23} - c_{12}s_{23}s_{13}e^{i\delta} & c_{12}c_{23} - s_{12}s_{23}s_{13}e^{i\delta} & s_{23}c_{13} \\ s_{12}s_{23} - c_{12}c_{23}s_{13}e^{i\delta} & -c_{12}s_{23} - s_{12}c_{23}s_{13}e^{i\delta} & c_{23}c_{13} \end{pmatrix} \text{diag}(e^{i\alpha}, e^{i\beta}, 1)$$

questions:

- Dirac / Majorana
- mass scale: m_1
- mass ordering: $\text{sgn}(\Delta m_{31}^2)$
- how small is θ_{13} , θ_{23} maximal?
- leptonic CP violation
- 3 flavour unitarity?
- why 3 generations, why $d=4$, ...



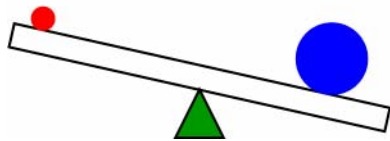
Suggestive Seesaw Features

QFT: natural value of mass operators \leftrightarrow scale of symmetry

$m_D \sim$ electro-weak scale

$M_R \sim$ L violation scale $\leftarrow? \rightarrow$ embedding (GUTs, ...)

See-saw mechanism (type I)



$$m_\nu = m_D M_R^{-1} m_D^T$$

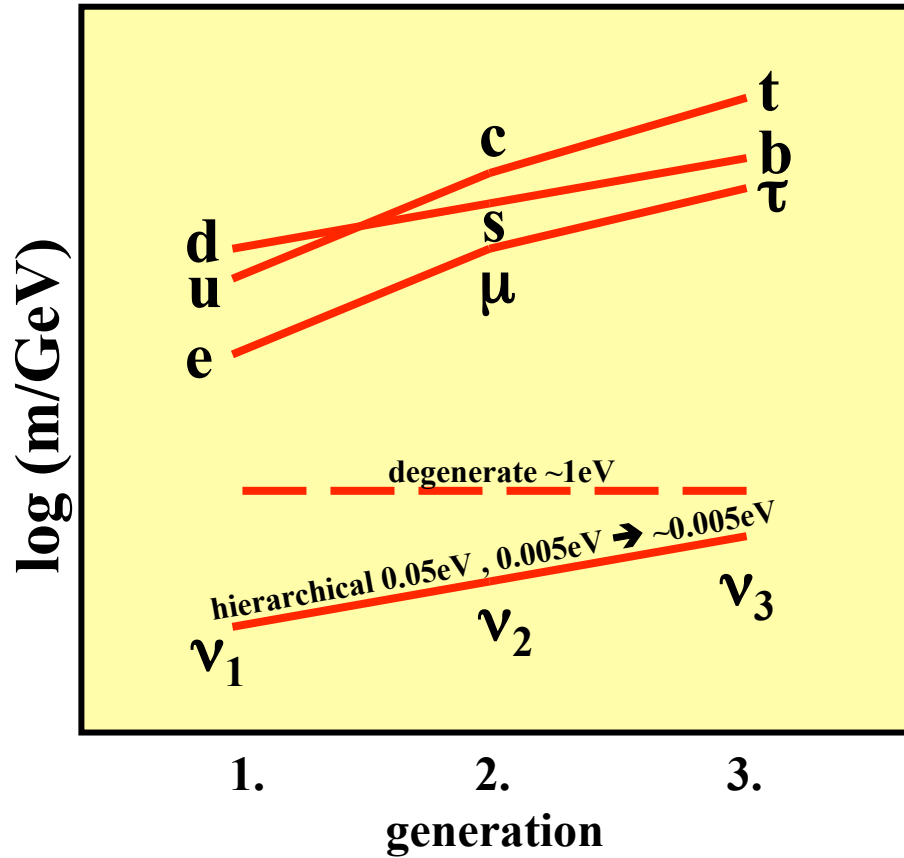
$$m_h = M_R$$

Numerical hints:

For $m_3 \sim (\Delta m_{\text{atm}}^2)^{1/2}$, $m_D \sim$ leptons $\rightarrow M_R \sim 10^{11} - 10^{16} \text{ GeV}$
 $\rightarrow \nu$'s are **Majorana particles**, m_ν probes \sim GUT scale physics!
 \rightarrow smallness of $m_\nu \leftrightarrow$ high scale of L , symmetries of m_D, M_R

2nd Look Questions

Quarks & charged leptons → hierarchical masses → neutrinos?



Quarks and charged leptons:

$$m_D \sim H^n ; n = 0, 1, 2 \rightarrow H \geq 20 \dots 200$$

Neutrinos:

$$m_\nu \sim H^n \rightarrow H \leq \sim 10$$

See-saw:

$$m_\nu = -m_D^T M_R^{-1} m_D$$

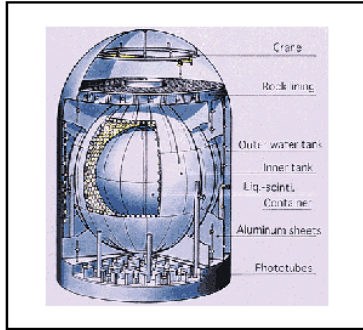
	↑	↑	↑	↑
H	~10	≥20	?	≥20

- » less hierarchy in m_D or corr. hierarchy in M_R ? → theoretically not connected!
- » other version of see-saw? → type II, III, ...?
- » Dirac masses?

Neutrino Mass Determinations

Subtitle: What do we actually measure?

Status of Neutrino Oscillations



Reactors: KAMLAND

**Beams: K2K → MINOS
→ T2K, OPERA**

improved results

$$\Delta m_{21}^2 = 7.59 \pm 0.20 \begin{pmatrix} +0.61 \\ -0.69 \end{pmatrix} \times 10^{-5} \text{ eV}^2$$

$$\Delta m_{31}^2 = \begin{cases} -2.40 \pm 0.11 \begin{pmatrix} +0.37 \\ -0.39 \end{pmatrix} \times 10^{-3} \text{ eV}^2 \text{ (inv)} \\ +2.51 \pm 0.12 \begin{pmatrix} +0.39 \\ -0.36 \end{pmatrix} \times 10^{-3} \text{ eV}^2 \text{ (nor)} \end{cases}$$

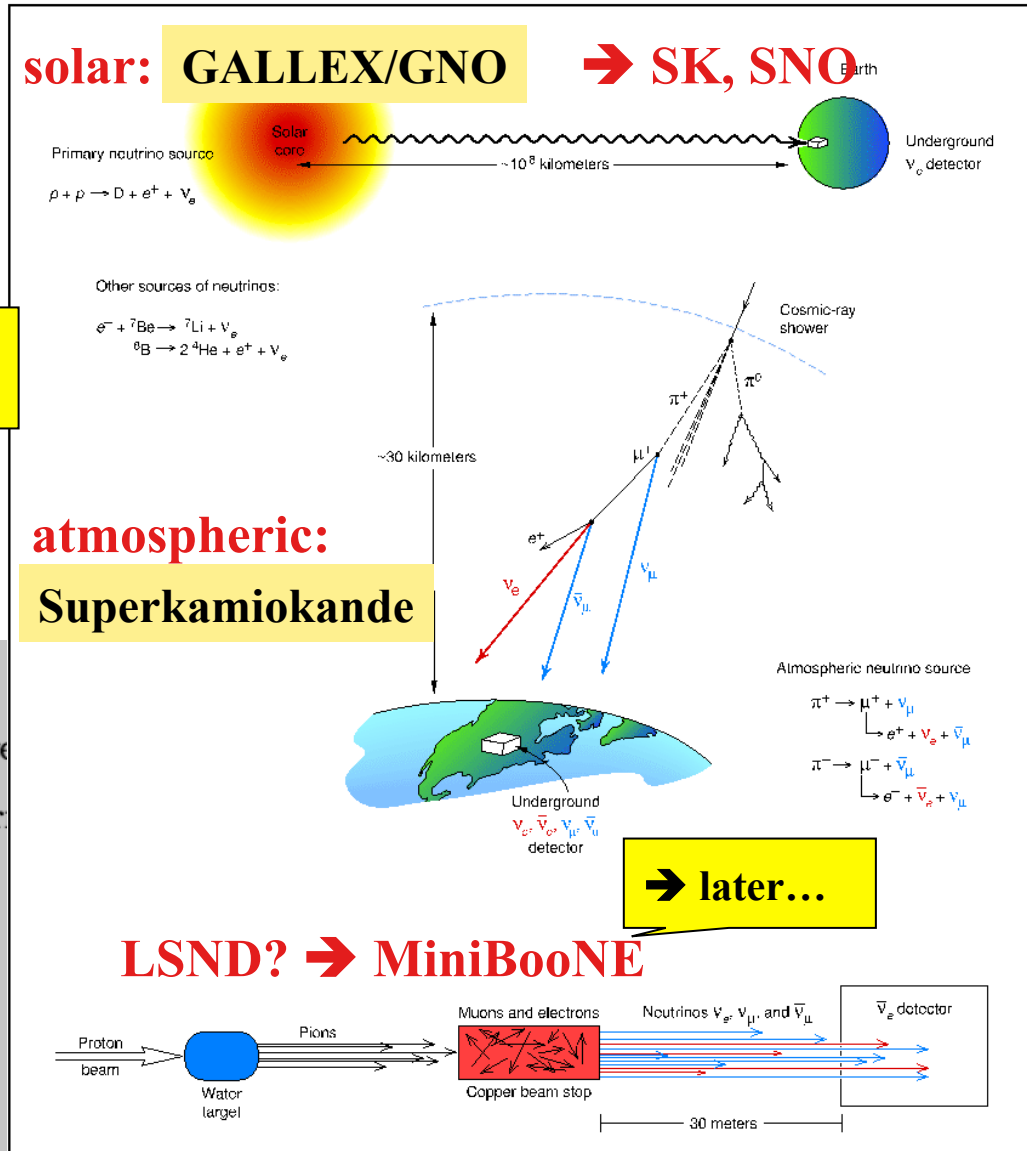
$$\theta_{12} = 34.4 \pm 1.0 \begin{pmatrix} +3.2 \\ -2.9 \end{pmatrix}$$

$$\theta_{23} = 42.3 \begin{matrix} +5.3 \\ -2.8 \end{matrix} \begin{pmatrix} +11.4 \\ -7.1 \end{pmatrix}$$

$$\theta_{13} = 6.8 \begin{matrix} +2.6 \\ -3.6 \end{matrix} (\leq 13.2)$$

$$[\sin^2 \theta_{13} = 0.014 \begin{matrix} +0.013 \\ -0.011 \end{matrix} (\leq 0.052)]$$

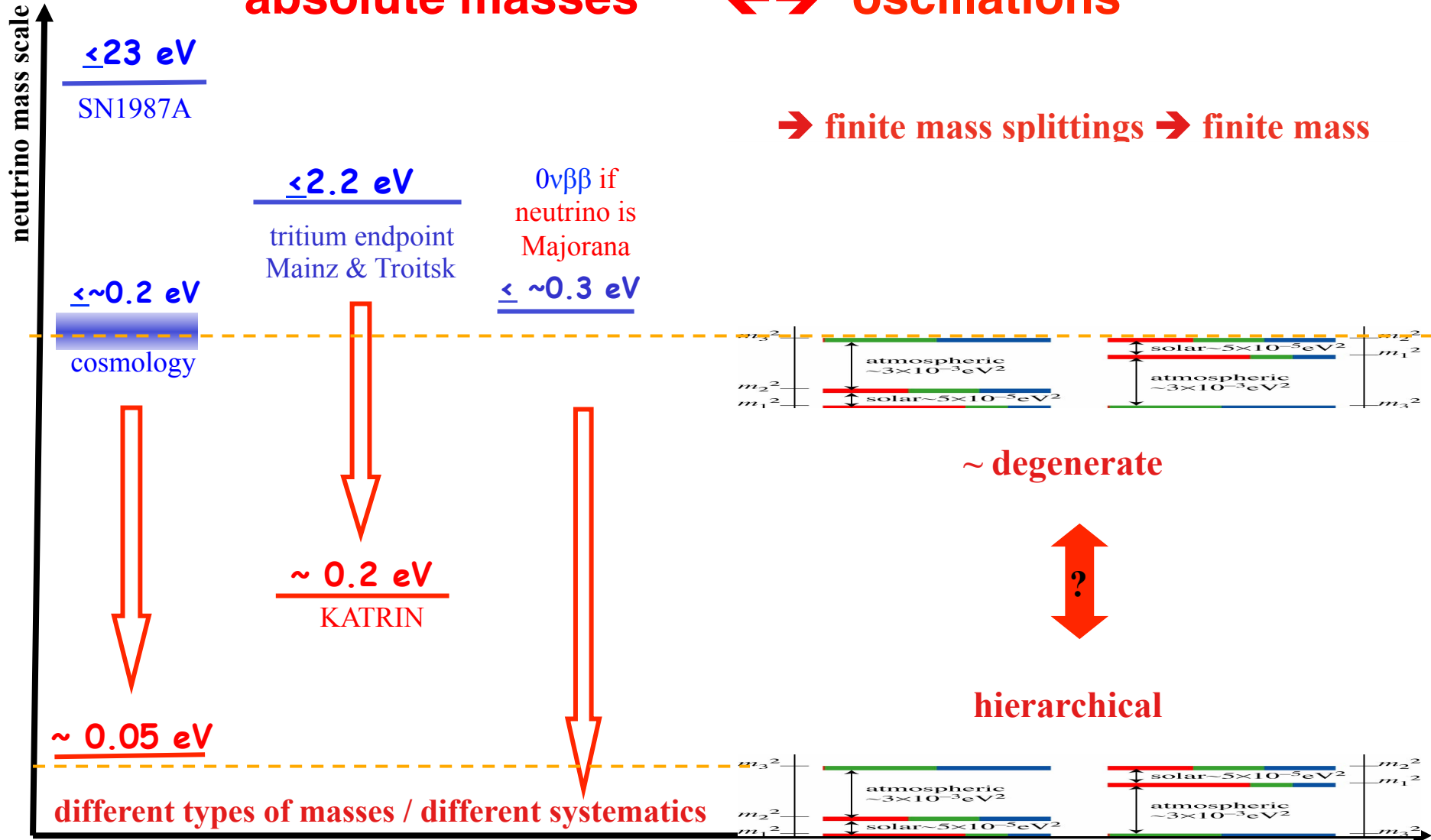
$$\delta_{CP} \in [0, 360]$$



What is special about 50meV ?

absolute masses

↔ oscillations



Four Methods of Mass Determination

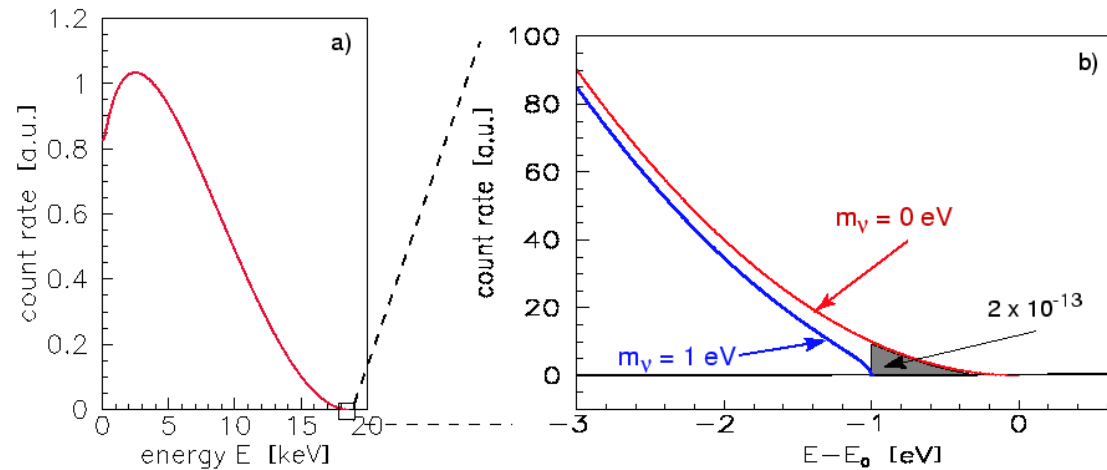
- **kinematical**
- **lepton number violation**
 \leftrightarrow **Majorana nature**
- **oscillations**
- **astrophysics & cosmology**

Kinematical Mass Determination

Relativistic kinematics:

$$E^2 = p^2 + m^2; \quad \sum p_i^\mu = \sum p_f^\mu$$

Endpoint of decays:



Bounds:

“Elektron-Neutrino”: $m < 2.2 \text{ eV}$ (Mainz, Troitsk)
 “Muon-Neutrino”: $m < 170 \text{ keV}$
 “Tau-Neutrino”: $m < 15.5 \text{ MeV}$

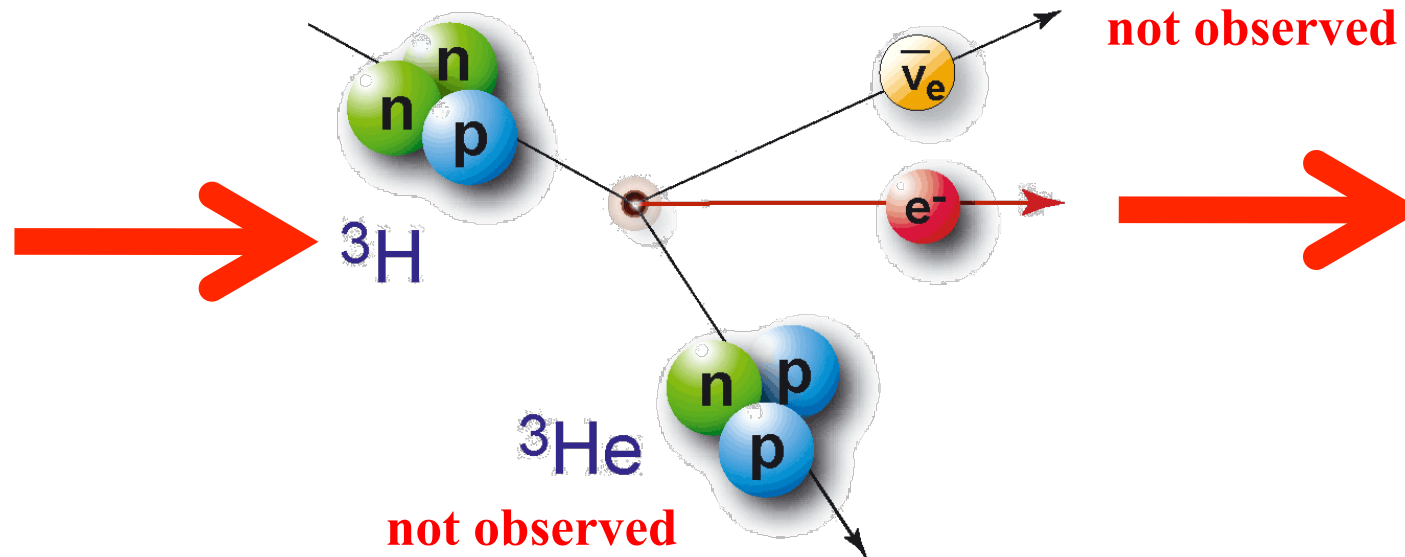
Sensitivity \Leftrightarrow degenerate ν -spectrum

$$\Rightarrow \text{Oscillations: } \Delta m_{ij}^2 \ll m_i^2 \Rightarrow \sum m_i^2 |U_{ei}|^2 < (2.2 \text{ eV})^2$$

Future: KATRIN \rightarrow 0.20 eV

\leftrightarrow c.f. cosmological bounds

Other Options?

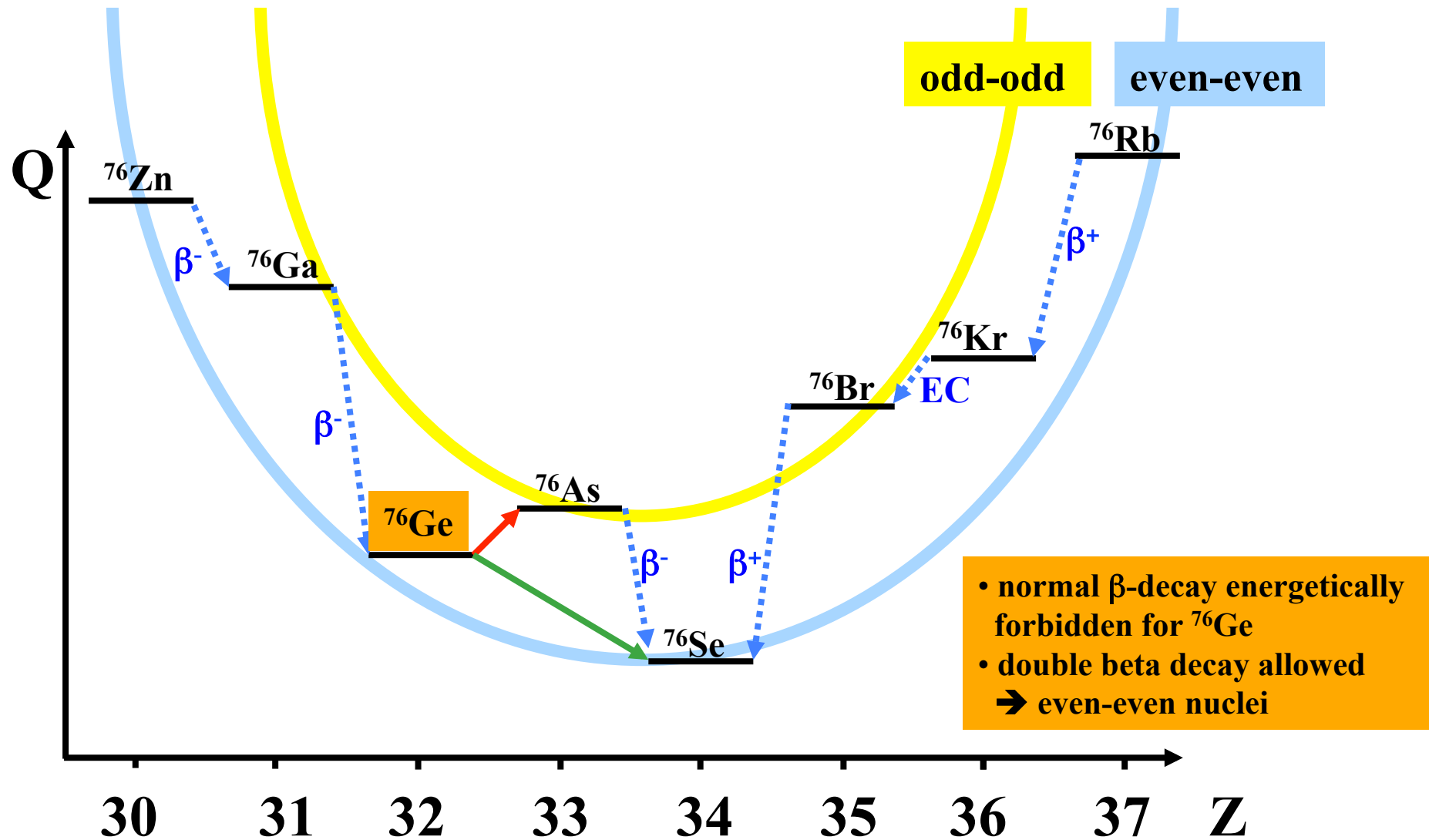


**Other new physics which modifies kinematics of the endpoint?
→ ...**

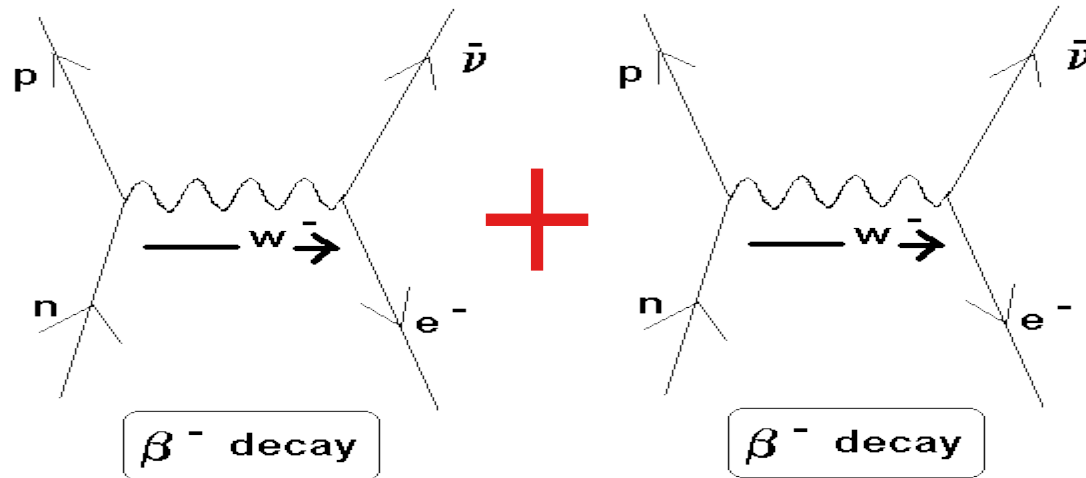
Four Methods of Mass Determination

- **kinematical**
- **lepton number violation**
 ↔ Majorana nature
- **oscillations**
- **astrophysics & cosmology**

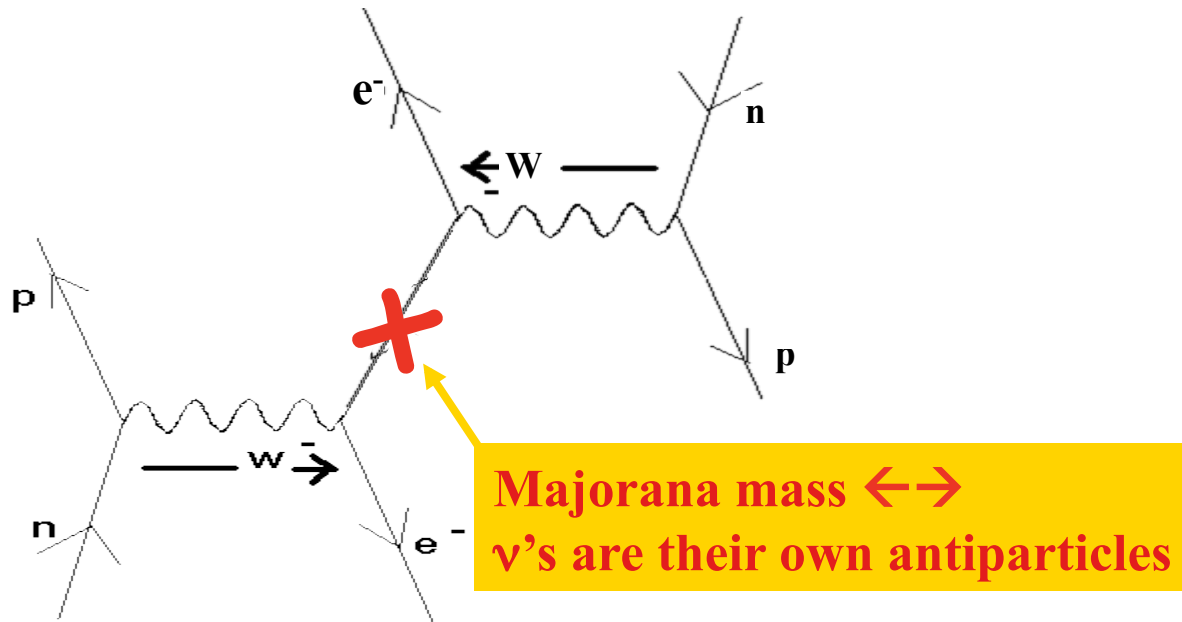
Double Beta Decay: Mass Parabolas



Double Beta Decay

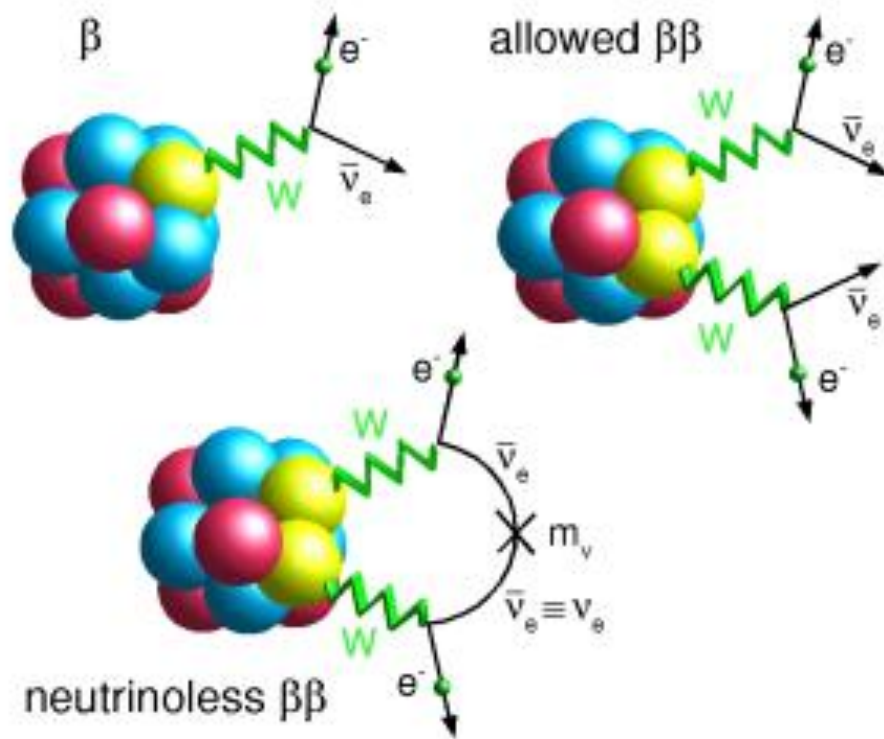


→ 2 neutrinos
plus 2 electrons

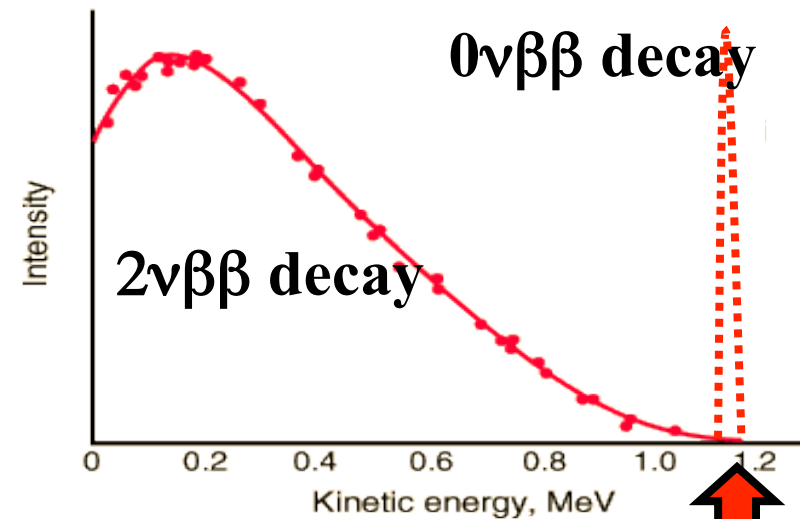


→ no neutrinos
just 2 electrons

$0\nu\beta\beta$ Decay Kinematics



$2\nu\beta\beta$ decay of ^{76}Ge observed:
 $\tau = 1.5 \times 10^{21}$ y



Majorana $\nu \rightarrow 0\nu\beta\beta$ decay

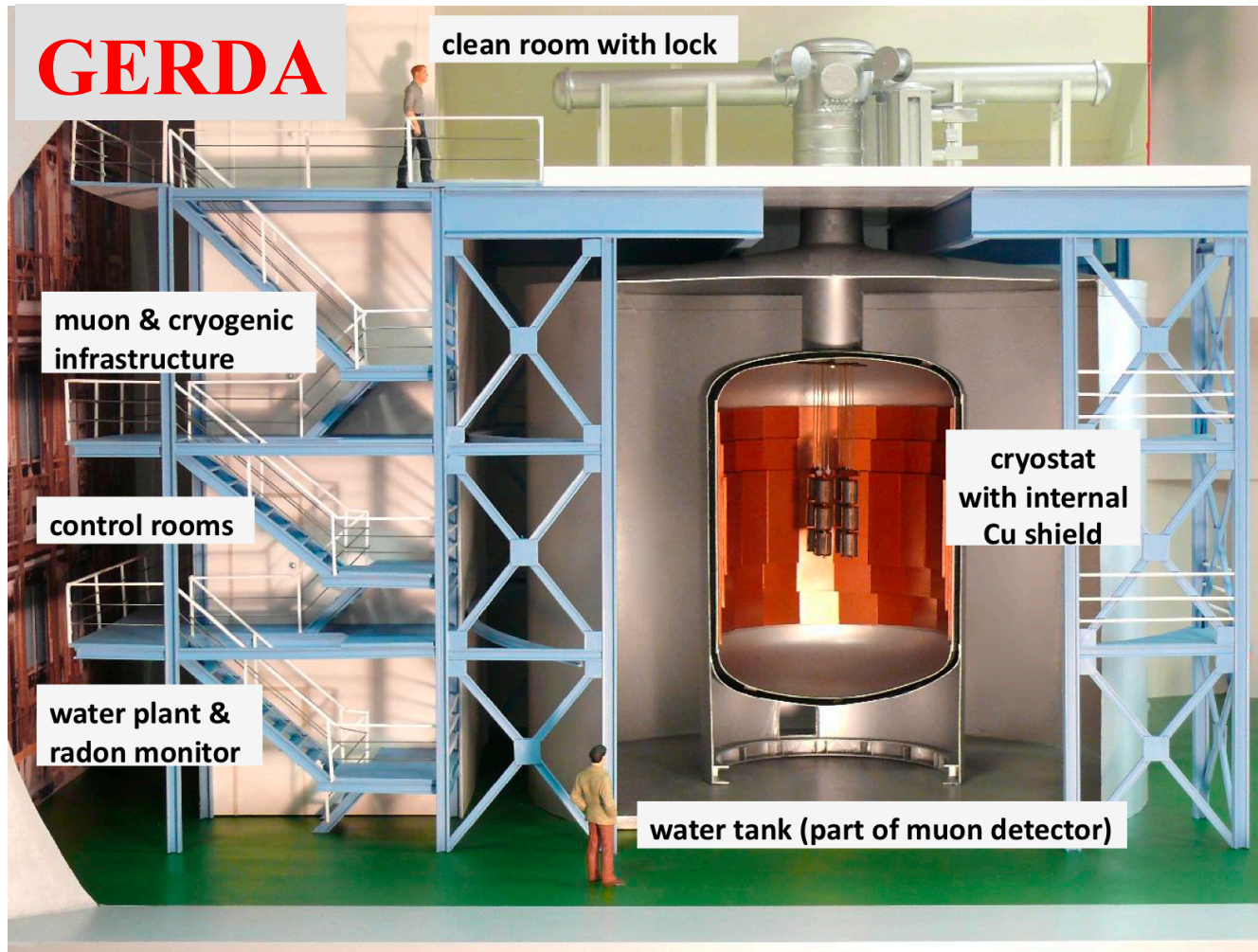
warning:

other lepton number violating processes...

- signal at known Q-value
- $2\nu\beta\beta$ background (resolution)
- nuclear backgrounds
- ➔ use different nuclei

New Experiments

GERDA



- GERDA (Ge76)
- EXO (Xe136)
- KamLAND-zen (Xe)

Construction:

- CUORE (Te-130)
- Majorana (Ge76)
- Super-NEMO (mult)



- ➔ filled, commissioning with non-enriched Ge
- ➔ Start of phase I: Nov. 1, 2011(enriched)

Decay Rates

$2\nu 2\beta$:

$$\left[T_{\frac{1}{2}}^{2\nu}(0^+ \rightarrow 0^+) \right]^{-1} = G^{2\nu}(E_0, Z) \left| M_{GT}^{2\nu} - \frac{g_V^2}{g_A^2} M_F^{2\nu} \right|^2$$

Phase space

nuclear matrix element

$0\nu 2\beta$:

(assuming that leading term is due to exchange of light Majorana-neutrinos)

$$\left[T_{\frac{1}{2}}^{0\nu}(0^+ \rightarrow 0^+) \right]^{-1} = G^{0\nu}(E_0, Z) \left| M_{GT}^{0\nu} - \frac{g_V^2}{g_A^2} M_F^{0\nu} \right|^2 \langle m_\nu \rangle^2$$

Warning: other...

effective Majorana mass

Relating Rates / Lifetimes to Neutrino Masses

rate of $0\nu\beta\beta$

phase space

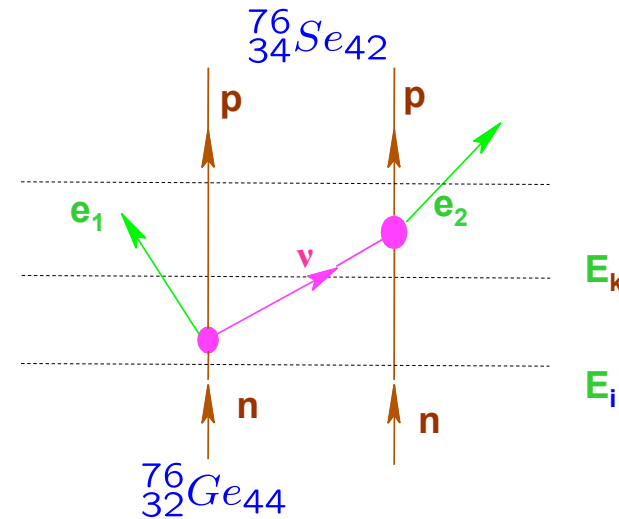
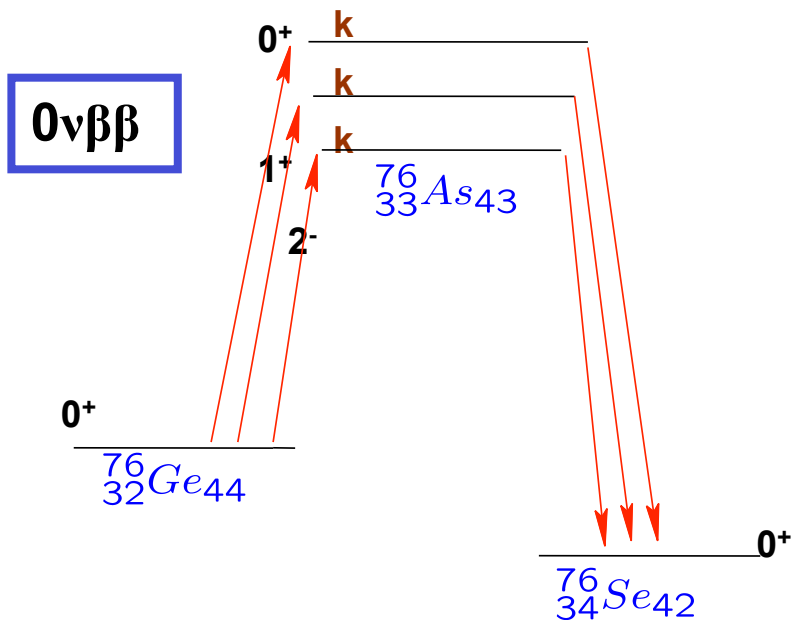
nuclear matrix elements

effective Majorana neutrino mass

$$1/\tau = G(Q,Z) |M_{\text{nucl}}|^2 \langle m_{ee} \rangle^2$$

nuclear matrix elements: Fäßler et al., ...

→ virtual excitations of intermediate states



$$T = \sum_k \frac{\langle f | \hat{H}_W | k \rangle \langle k | \hat{H}_W | i \rangle}{E_i - E_k}$$

progress in TH errors → reduced uncertainties

Nuclear Matrix Elements

$0\nu 2\beta$ half-lives in units of 10^{26} years for $\langle m_\nu \rangle = 50$ meV for nuclear matrix for different methods and authors

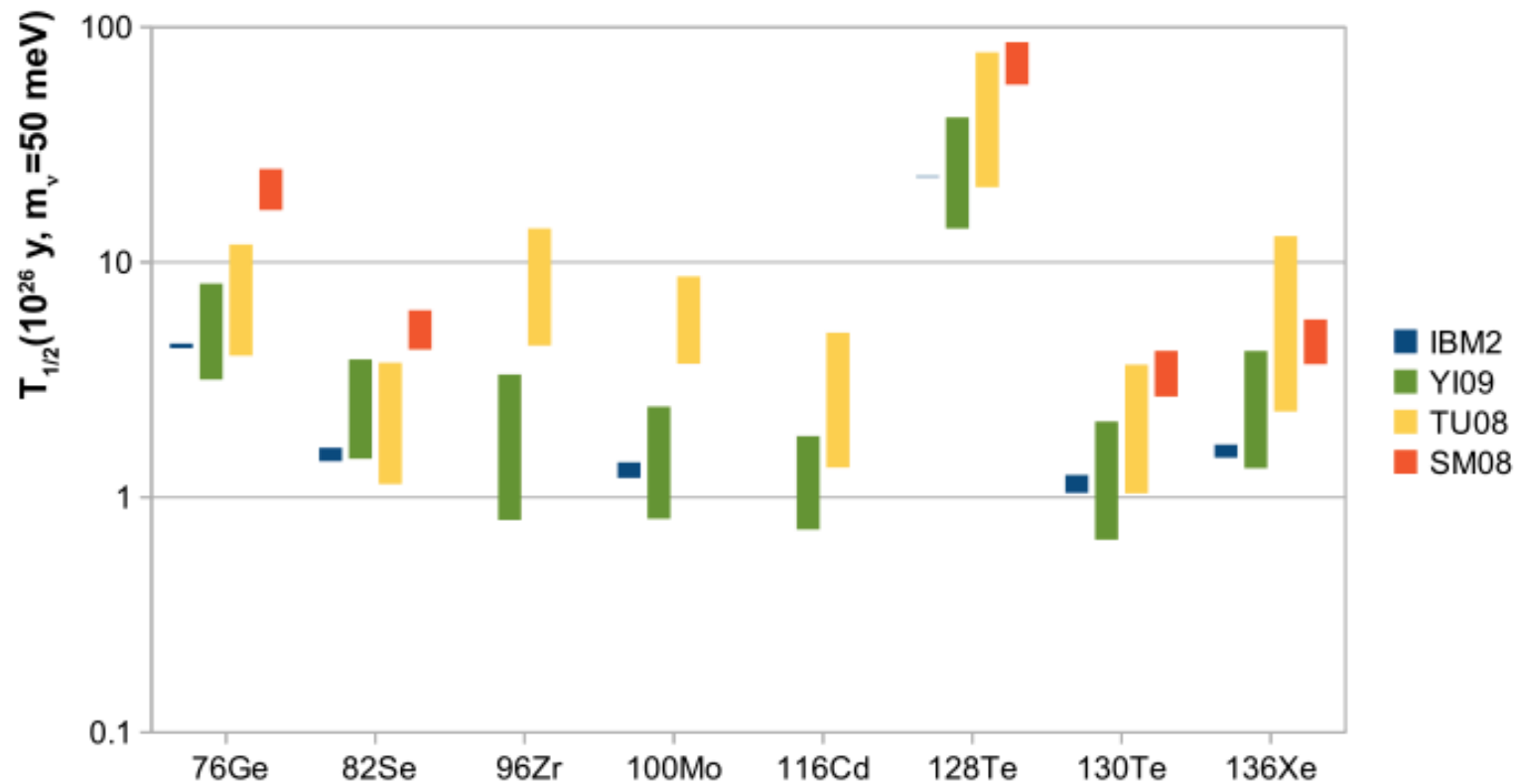
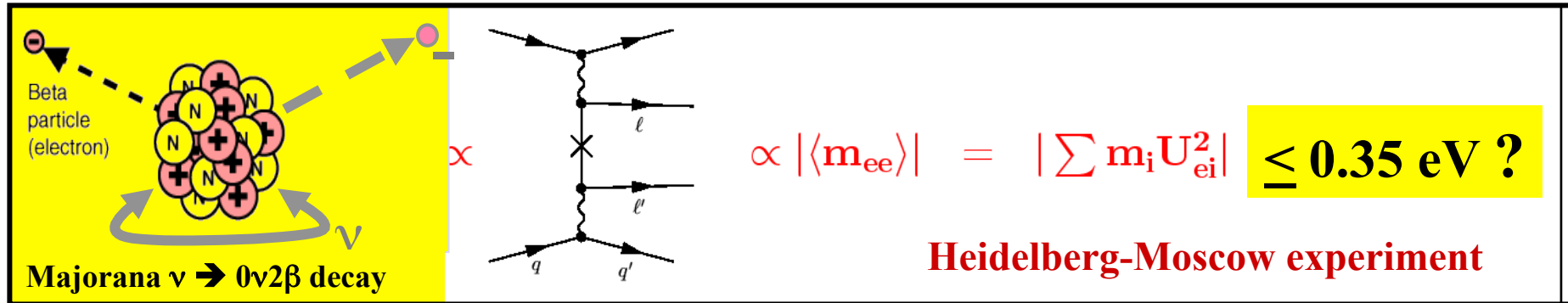


Figure 1: Expected $\beta\beta(0\nu)$ half lives for 50 meV effective neutrino mass and different NME calculations: IBM2 [17], YI09 [18], TU08 [19] and SM08 [20].

Effective Majorana Mass



The diagram shows a nucleus with neutrons (N) and protons (p) undergoing a $0\nu 2\beta$ decay. A Majorana neutrino $\bar{\nu}$ is exchanged between two neutrons, resulting in two protons and an electron (beta particle). The Feynman diagram shows a Majorana neutrino propagator (wavy line with X) connecting two vertices. The left vertex involves a quark q and a Majorana neutrino $\bar{\nu}$. The right vertex involves a lepton ℓ and a Majorana neutrino ν .

$\propto |\langle m_{ee} \rangle| = |\sum m_i U_{ei}^2| \leq 0.35 \text{ eV ?}$

Heidelberg-Moscow experiment

$$m_{ee} = |m_{ee}^{(1)}| + |m_{ee}^{(2)}| \cdot e^{i\Phi_2} + |m_{ee}^{(3)}| \cdot e^{i\Phi_3}$$

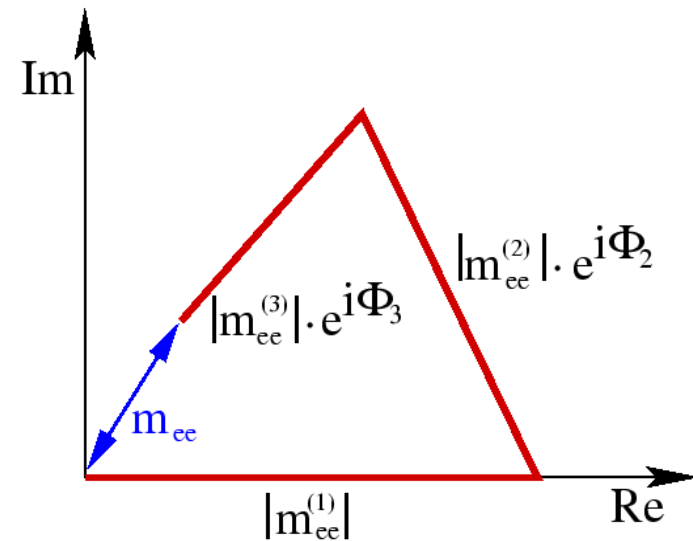
$$|m_{ee}^{(1)}| = |U_{e1}|^2 m_1$$

$$|m_{ee}^{(2)}| = |U_{e2}|^2 \sqrt{m_1^2 + \Delta m_{21}^2}$$

$$|m_{ee}^{(3)}| = |U_{e3}|^2 \sqrt{m_1^2 + \Delta m_{31}^2}$$

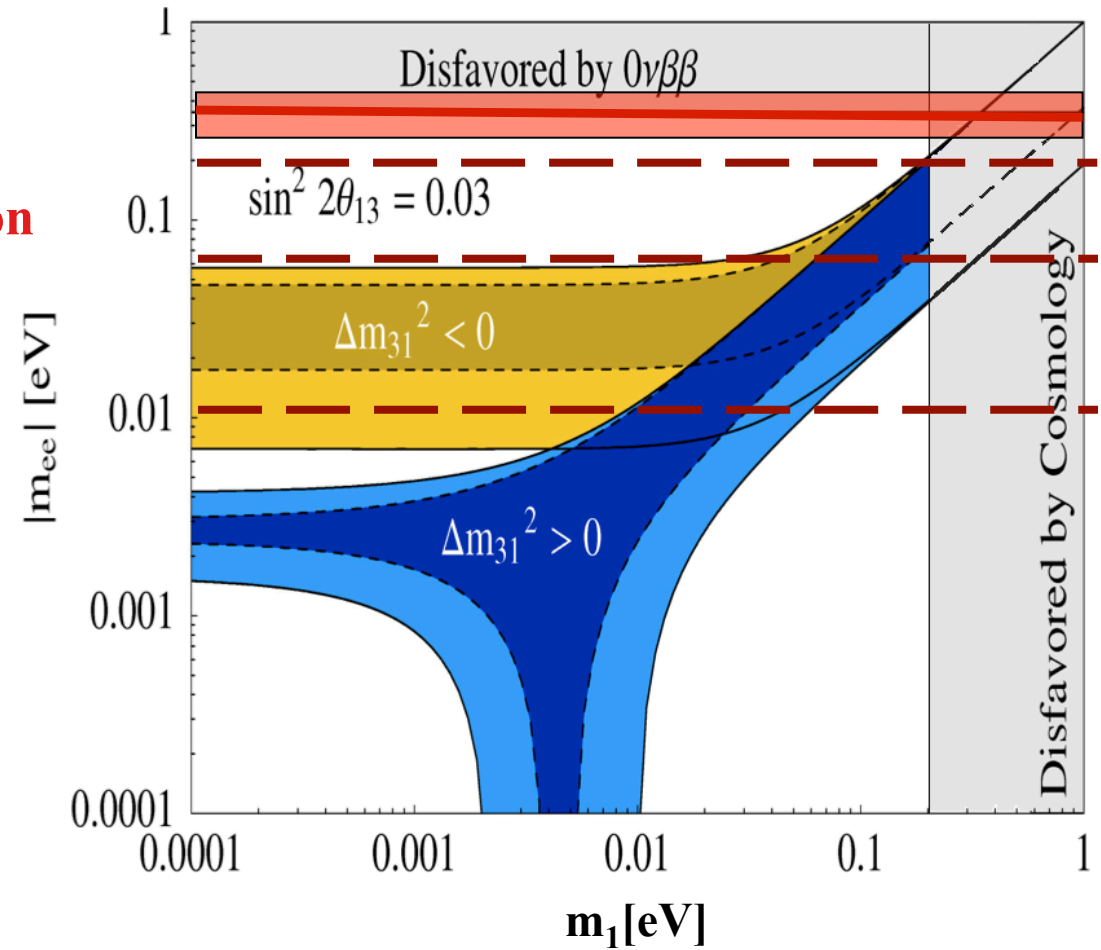
solar $\Rightarrow |U_{e1}|^2, |U_{e2}|^2, \Delta m_{21}^2$
 atmosph. $\Rightarrow |\Delta m_{31}^2|$
 CHOOZ $\Rightarrow |U_{e3}|^2 < 0.05$

\rightarrow free parameters: $m_1, \text{sign}(\Delta m_{31}^2), \text{CP-phases } \Phi_2, \Phi_3$



Claim of part of the original Heidelberg-Moscow collaboration
 \leftrightarrow cosmology \rightarrow ,tension‘

- aims of new experiments:**
- test HM claim
 - $(\Delta m_{31}^2)^{1/2} \simeq 0.05\text{eV} \pm \text{errors}$
 - \rightarrow reach 0.01eV
 - \rightarrow CUORE
 - \rightarrow GERDA phases I, II, (III)



Comments:

- cosmology: limitation by systematical errors \rightarrow ~another factor 5?
- $0\nu\beta\beta$ nuclear matrix elements ~factor 1.3-2 **theoretical** uncertainty in m_{ee}
- $\Delta m^2 > 0$ allows complete cancellation
 - \rightarrow $0\nu\beta\beta$ signal not guaranteed, but cancelation appears unlikely

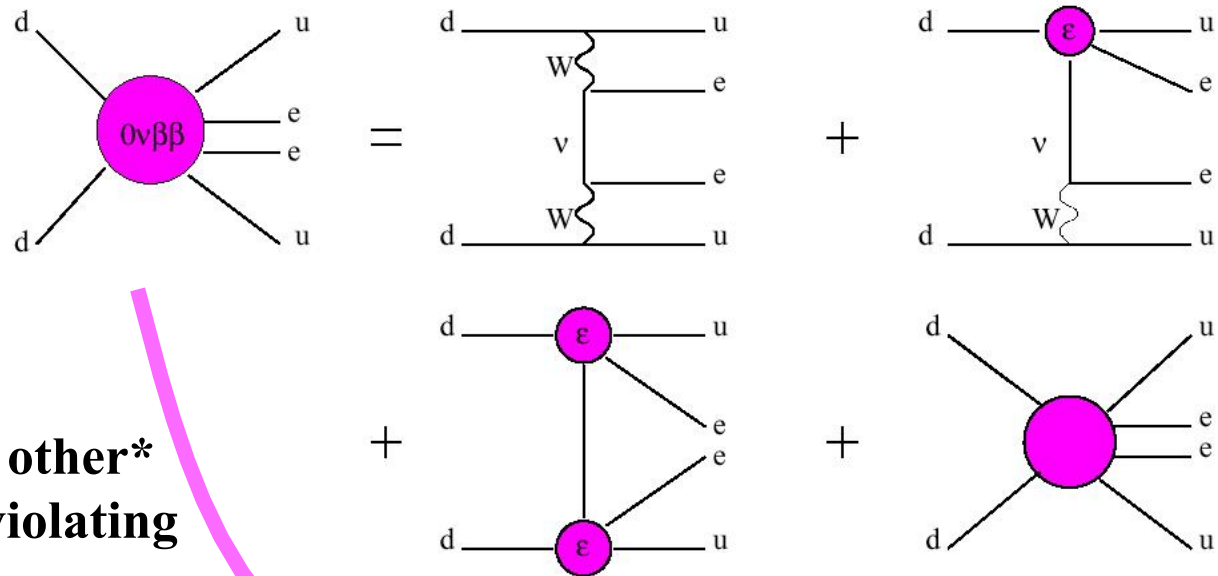
$0\nu\beta\beta$ from Alternative $\Delta L=2$ Operators

Various possibilities:

- LR symmetry
- SUSY (RPV)
- ...

→ $0\nu\beta\beta$ signal from *some other* new BSM lepton number violating operator

→ very promising interplay of neutrino mass determinations, cosmology, LHC, LNF experiments and theory

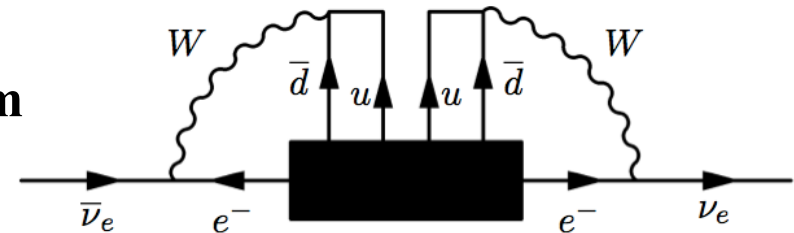


Schechter + Valle: Any $\Delta L=2$ violating operator

→ radiative generation of Majorana mass term

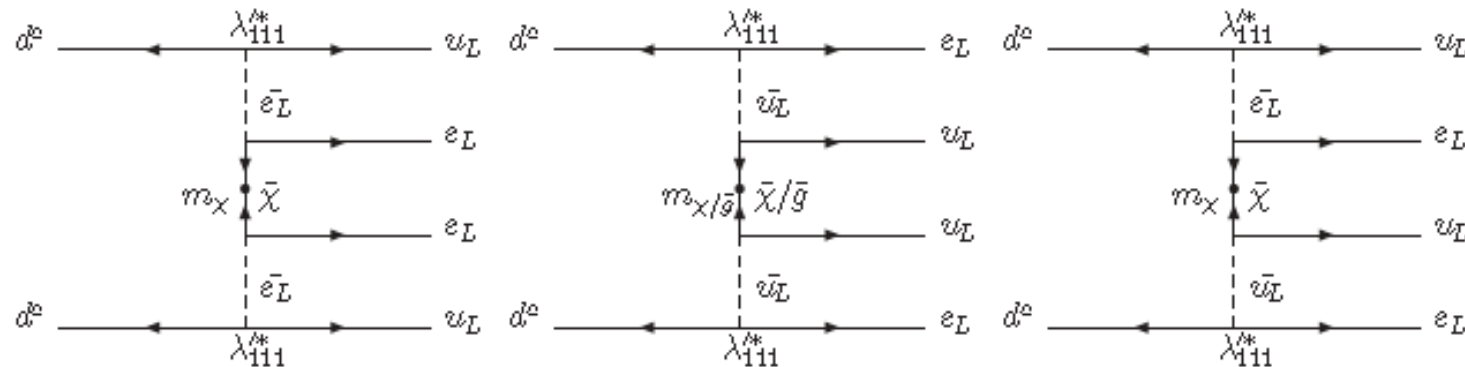
→ Majorana nature of ν 's guaranteed

→ but how big is the mass?



SUSY Example

Direct, TeV scale short range mediation w/o intermediate light ν , e.g.

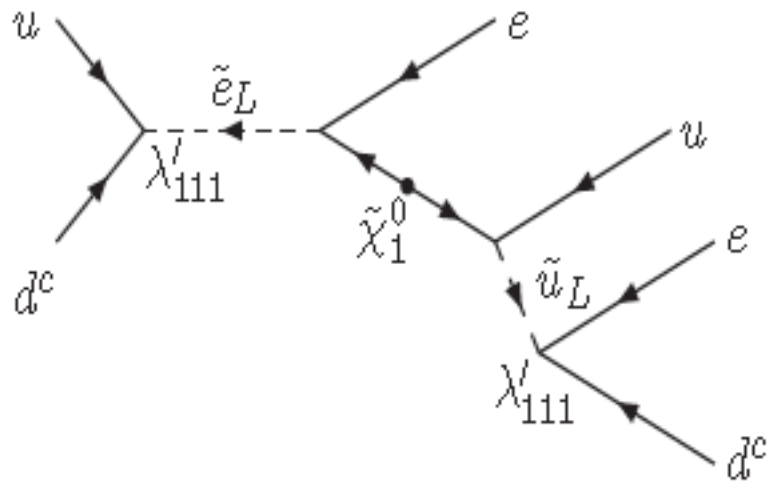


$$\mathcal{L}_{\lambda_{111}^{\prime} \lambda_{111}^*}^{eff, \Delta L_e=2}(x) = \frac{G_F^2}{2} m_p^{-1} [\bar{e}(1 + \gamma_5)e^c] \times \left[(\epsilon_{\tilde{g}} + \epsilon_{\chi})(J_{PS} J_{PS} - \frac{1}{4} J_T^{\mu\nu} J_{T\mu\nu}) + (\epsilon_{\chi\tilde{e}} + \epsilon'_{\tilde{g}} + \epsilon_{\chi\tilde{f}}) J_{PS} J_{PS} \right]$$

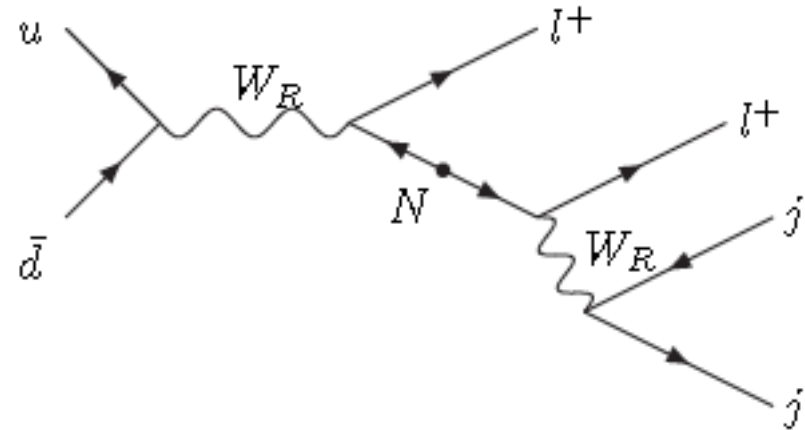
$$\epsilon_i \sim \pi \alpha_{(Strong, EW)} \frac{\lambda_{111}^{\prime 2}}{G_F^2} \frac{m_p}{m_{(\tilde{g}, \tilde{\chi})}} \frac{1}{m_{(\tilde{u}, \tilde{d}, \tilde{e})}^4}.$$

$\Delta L=2$ Operators and TeV Scale Physics

SUSY: direct test of λ'_{111}

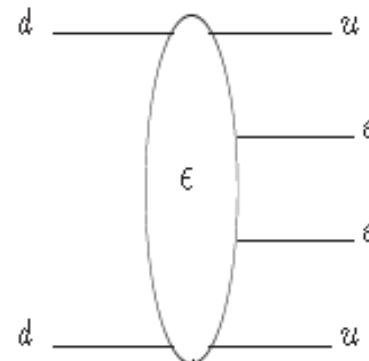
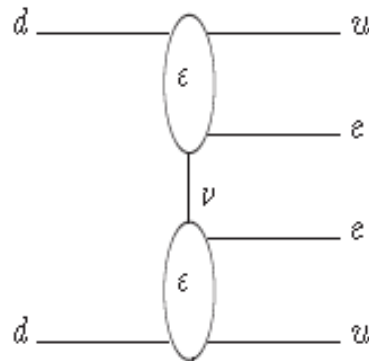


L-R symmetry: heavy N's



Relative strength of 'light' and 'heavy' $0\nu\beta\beta$ amplitudes:

$$M_{\text{light}} \sim G_F^2 \frac{m_{\beta\beta}}{\langle k^2 \rangle}$$



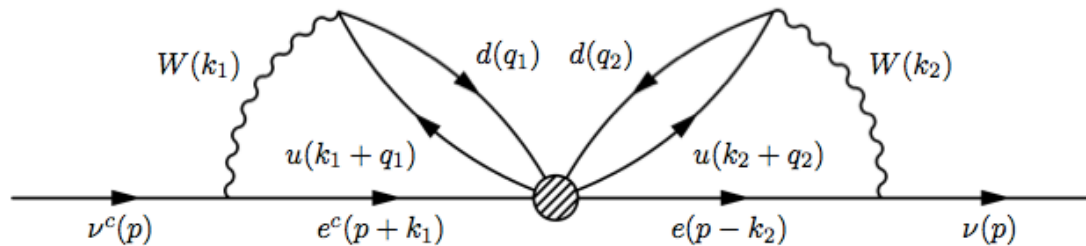
$$M_{\text{heavy}} \sim G_F^2 \left(\frac{\lambda}{g_2} \right)^4 \frac{M_W^4}{\Lambda^5}$$

SV-induced Neutrino Masses

General Lorentz-invariant Lagrangian for $0\nu\beta\beta$ (point operator)

$$\mathcal{L} = \frac{G_F^2}{2} m_p^{-1} (\epsilon_1 J J j + \epsilon_2 J^{\mu\nu} J_{\mu\nu} j + \epsilon_3 J^\mu J_\mu j + \epsilon_4 J^\mu J_{\mu\nu} j^\nu + \epsilon_5 J^\mu J j_\mu)$$

$$J = \bar{u} (1 \pm \gamma_5) d, \quad J^\mu = \bar{u} \gamma^\mu (1 \pm \gamma_5) d \text{ etc.}$$



Outcome:

M. Dürr, ML, A. Merle, arXiv:1105.0901

If other $\Delta L=2$ physics drives $0\nu\beta\beta \rightarrow$ SV gives $\delta m_\nu = 10^{-24}$ eV

\rightarrow mass correction too small to explain observed masses and splittings

\rightarrow explicit neutrino mass operators required

Dirac: $0\nu\beta\beta$ essentially unrelated to neutrino masses \leftrightarrow other BSM

Majorana: dominates over SV contribution

$0\nu\beta\beta$ may be a mixture of Majorana mass and other $\Delta L=2$ physics

\rightarrow mimics higher Majorana neutrino mass

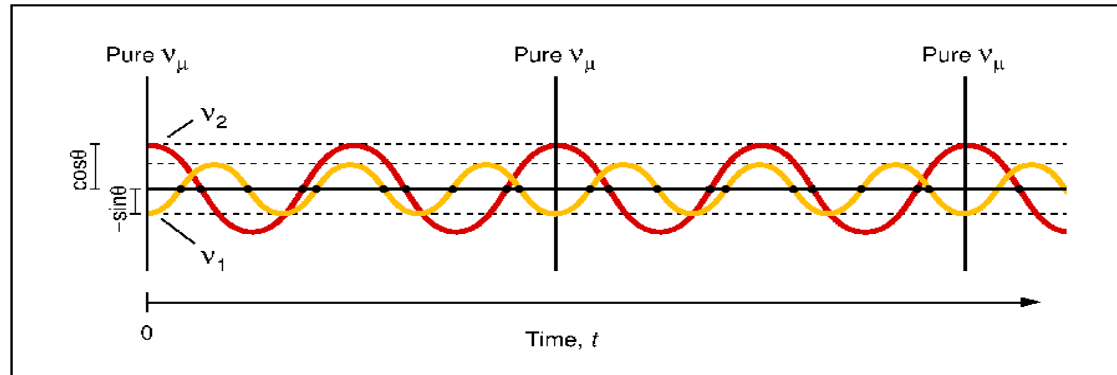
Four Methods of Mass Determination

- **kinematical**
- **lepton number violation**
 \leftrightarrow **Majorana nature**
- **oscillations**
- **astrophysics & cosmology**

Two Neutrino Oscillations

2 Neutrinos: ν_e, ν_μ

$$\begin{aligned} |\nu_e(0)\rangle &= \cos\theta |\nu_1\rangle + \sin\theta |\nu_2\rangle \\ |\nu_\mu(0)\rangle &= -\sin\theta |\nu_1\rangle + \cos\theta |\nu_2\rangle \end{aligned}$$



$$|\nu_\mu(t)\rangle = -\sin\theta \exp\left[-\frac{iE_1 t}{\hbar}\right] |\nu_1\rangle + \cos\theta \exp\left[-\frac{iE_2 t}{\hbar}\right] |\nu_2\rangle$$

$$E_i = \sqrt{p_i^2 + m_i^2} \quad p_i = p \gg m_i \quad \simeq p + \frac{m_i^2}{2p} \quad \simeq p + \frac{m_i^2}{2E}$$

$$L = c \cdot t \quad \Delta m^2 = m_2^2 - m_1^2 \Rightarrow E_2 - E_1 = \frac{\Delta m^2}{2E}$$

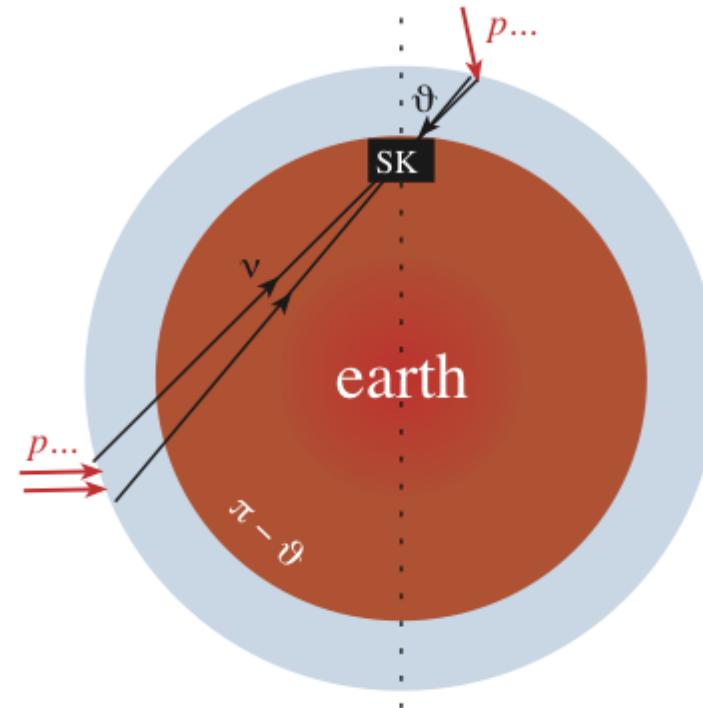
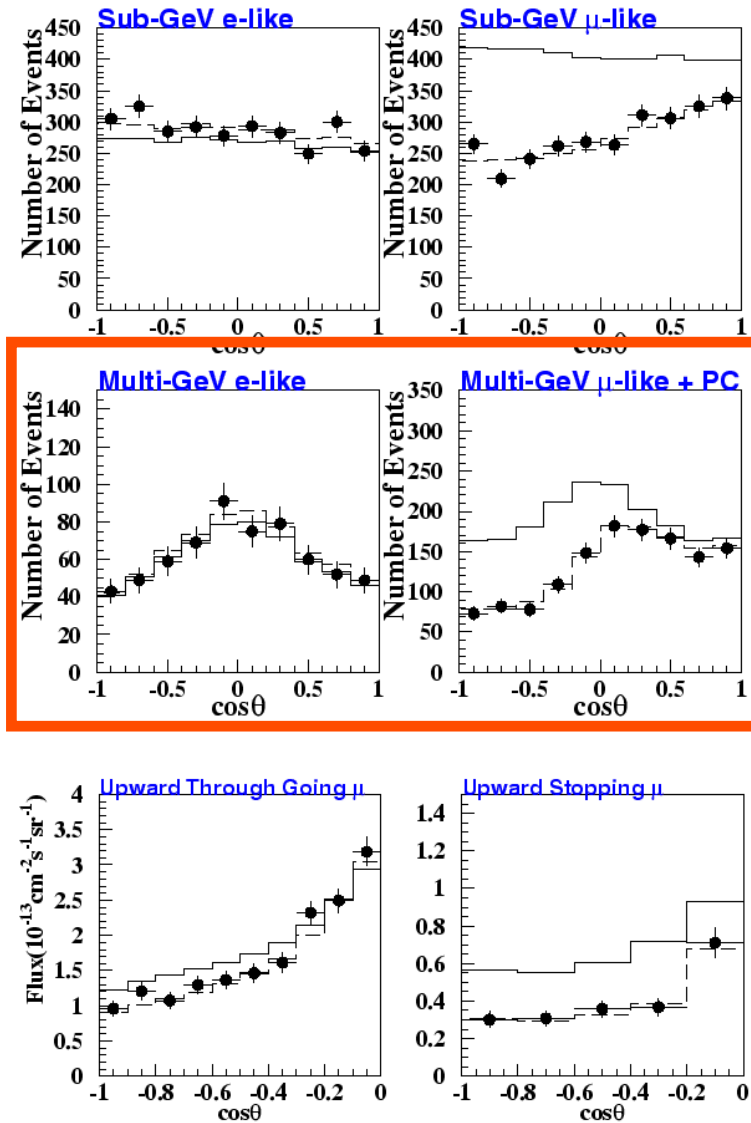
**2ν-transition-
probability:**

$$P(\nu_\mu \rightarrow \nu_e) = |\langle \nu_\mu(t) | \nu_e(0) \rangle|^2 = \sin^2 2\theta \cdot \sin^2 \left(\frac{\Delta m^2 L}{4E} \right)$$

$\nu_e, \nu_\mu, \nu_\tau \rightarrow 9$ oscillation channels for neutrinos

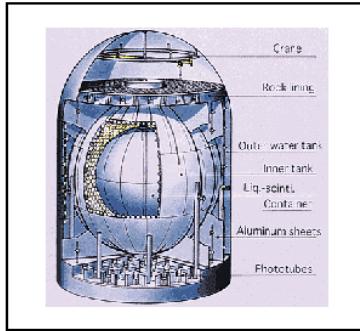
$\bar{\nu}_e, \bar{\nu}_\mu, \bar{\nu}_\tau \rightarrow 9$ channels for anti-neutrinos (assuming 3ν !)

Atmospheric Oscillations @SuperK



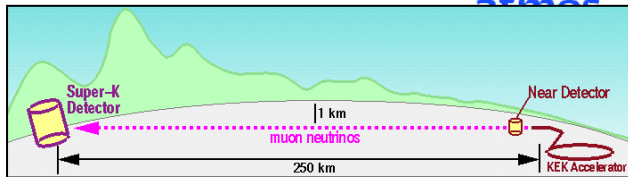
- 8σ signal for ν_μ disappearance
- NOT $\nu_\mu \rightarrow \nu_e$ (consistent with Chooz)
- $\Rightarrow \nu_\mu \rightarrow \nu_\tau$ (some τ 's seen)
- NOT $\nu_\mu \rightarrow \nu_s$ from NC/CC comparison
- L/E confirmed by K2K ($\simeq 2\sigma$)
- sensitivity for L/E-dependence of oscillations

Status of Neutrino Oscillations



Reactors: KAMLAND

**Beams: K2K → MINOS
→ OPERA**



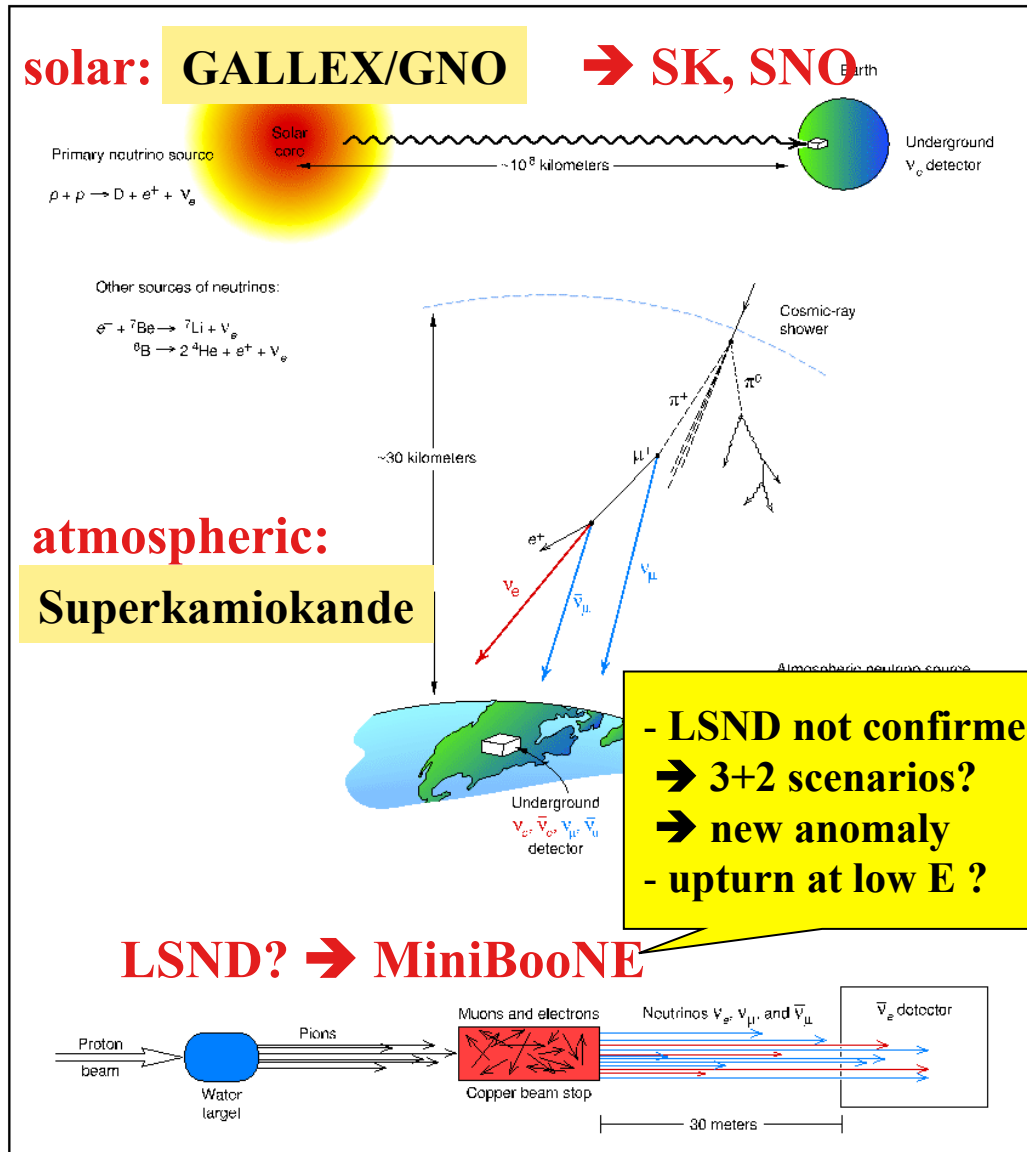
$$\Delta m_{21}^2 = (7.9 \pm 0.3) * 10^{-5} \text{ eV}^2$$

$$\tan^2 \theta_{12} = 0.39 \pm 0.05$$

$$\Delta m_{31}^2 = (2.4 \pm 0.3) * 10^{-3} \text{ eV}^2$$

$$\tan^2 \theta_{23} = 1.0 \pm 0.3$$

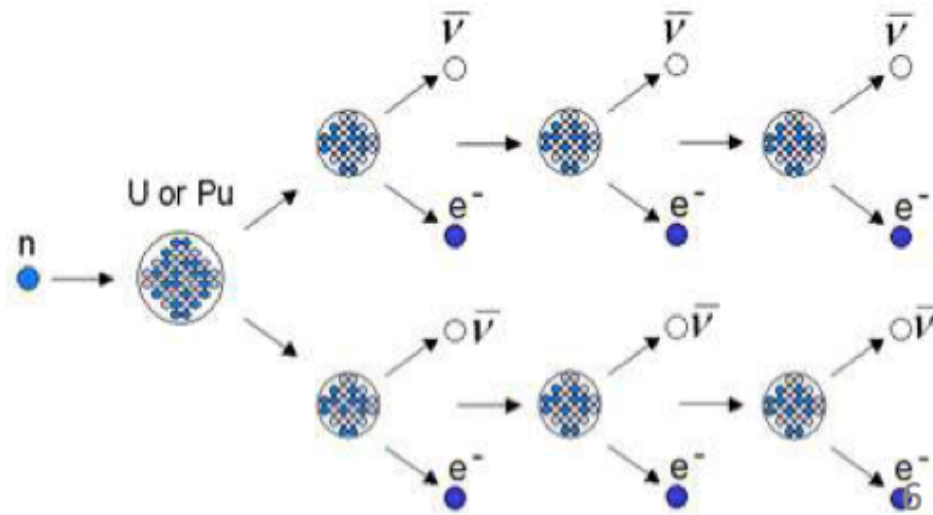
$$\sin^2 2\theta_{13} = 0.09 \pm \sim 0.02$$



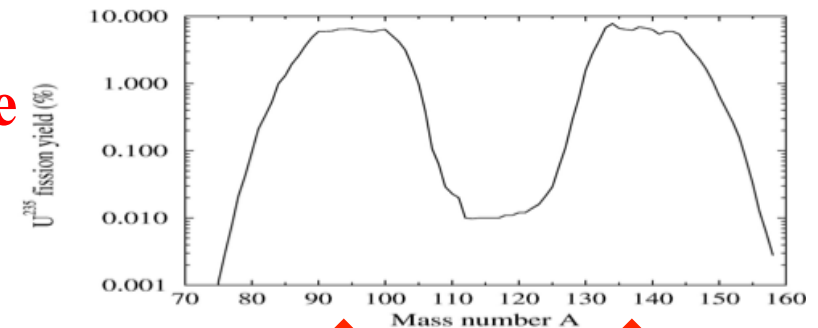
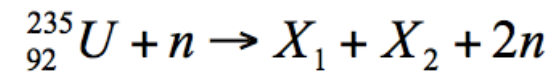
- LSND not confirmed!
 → 3+2 scenarios?
 → new anomaly
 - upturn at low E?

Nuclear Reactors as Antineutrino Source

- Reactors like Chooz $A+B \rightarrow 8.5 \text{ GW}_{\text{th}}$
- Few percent of the released energy
 - \rightarrow escapes with anti-neutrinos
 - $\rightarrow 2 \cdot 10^{21} \bar{\nu}/\text{s} \leftrightarrow O(1 \text{ kW}/\text{m}^2) @\text{fence}$



example: fission of U^{235}



most \uparrow likely A

\rightarrow on average:

- measured e^- spectrum of U^{235} , Pu^{239} , Pu^{241}
 - \rightarrow calculate $\bar{\nu}_e$ - spectrum \rightarrow precisions?

- 6 neutrons β -decay to 6 protons to reach stable matter
- $1.5 \nu_e$ emitted with $E > 1.8 \text{ MeV}$

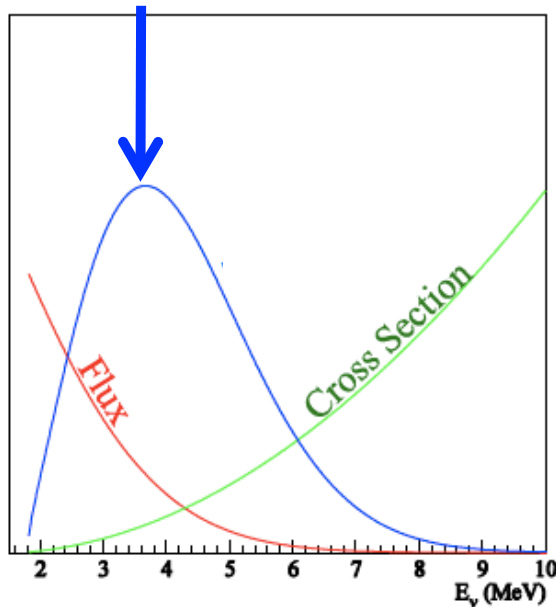
Anti-Neutrino Detection

Oscillations:

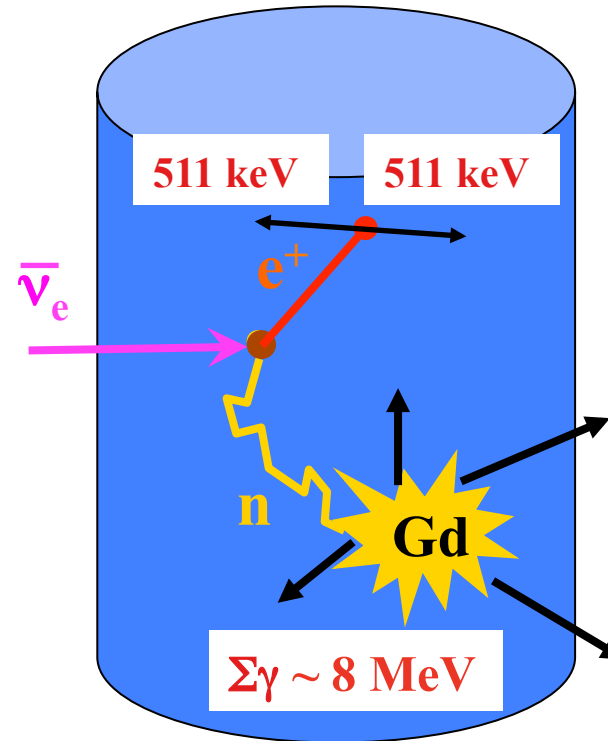
- affect rate & shape

Earlier reactor experiments:

- calculated spectrum
- rate normalized by P_{thermal}
- event rate = flux * x-section



- uncertainties in x-sections?



prompt e^+ signal

delayed n capture
 \rightarrow Gd doping
 \rightarrow delayed γ (30 μs)

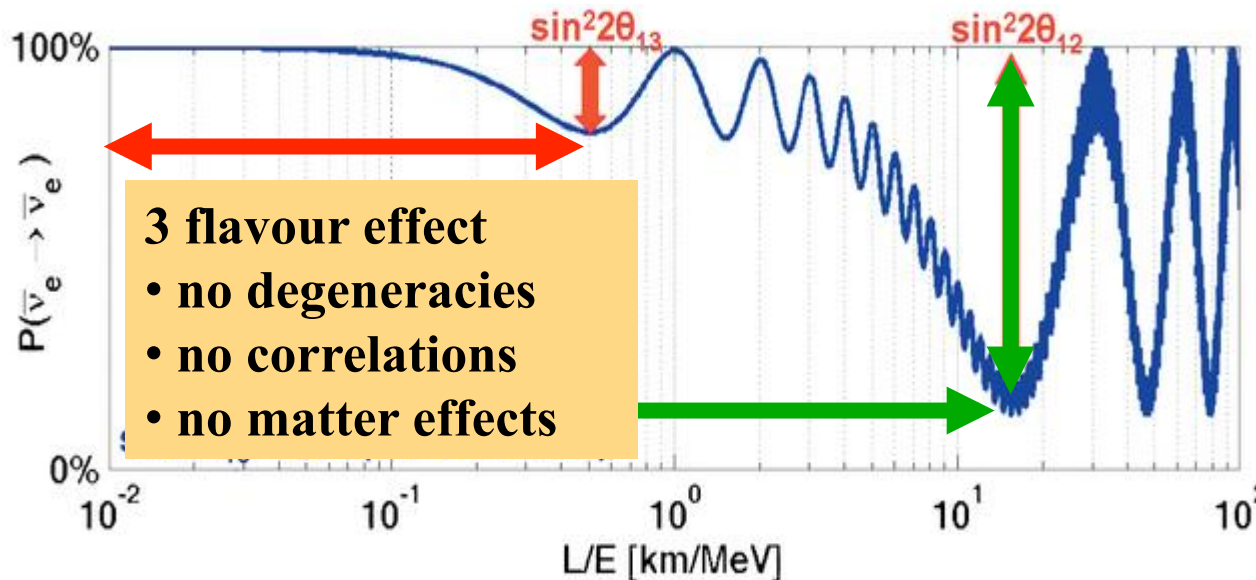
- position & time correlation
 - delayed energy information
 \rightarrow background reduction!
- Gd loaded liquid scintillator
 \rightarrow stability, transparency, WLS, ...

Precision with Reactor Experiments



identical detectors → many errors cancel

$$P_{ee} \approx 1 - \sin^2 2\theta_{13} \sin^2 \frac{\Delta m_{31}^2 L}{4E_\nu} - \left(\frac{\Delta m_{21}^2 L}{4E_\nu} \right)^2 \cos^4 \theta_{13} \sin^2 2\theta_{12}$$



clean & precise
 θ_{13} measurements
 ←→ beams

→ Double Chooz
 → Daya Bay
 → Reno

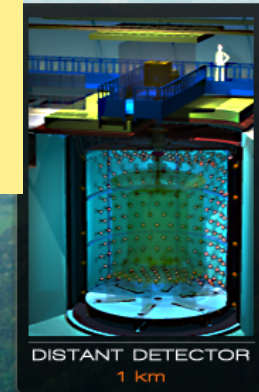
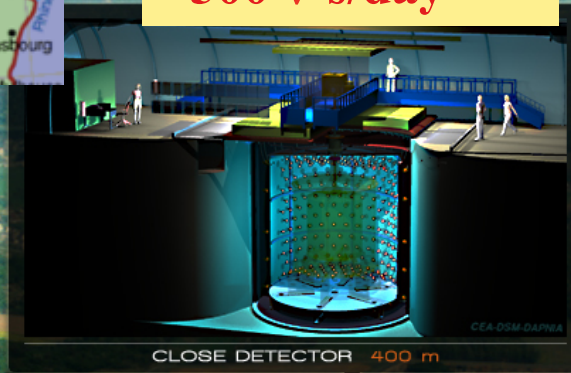
E=4MeV → 1km 180km

Double Chooz



Near lab:
410m, 120 mwe
~500 ν 's/day

Far lab:
1050m, 300mwe
~70 ν 's/day

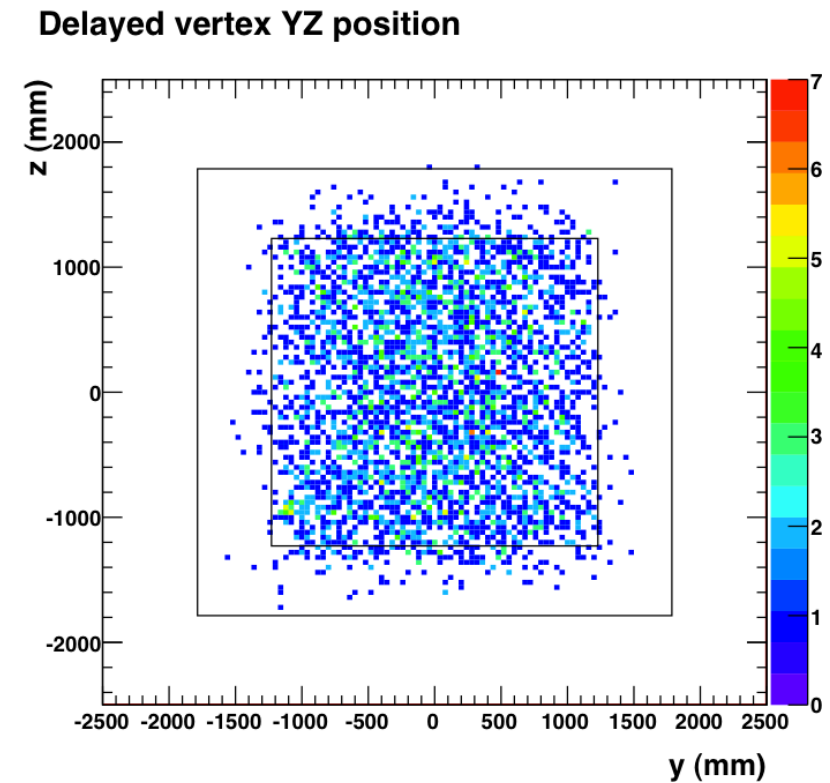
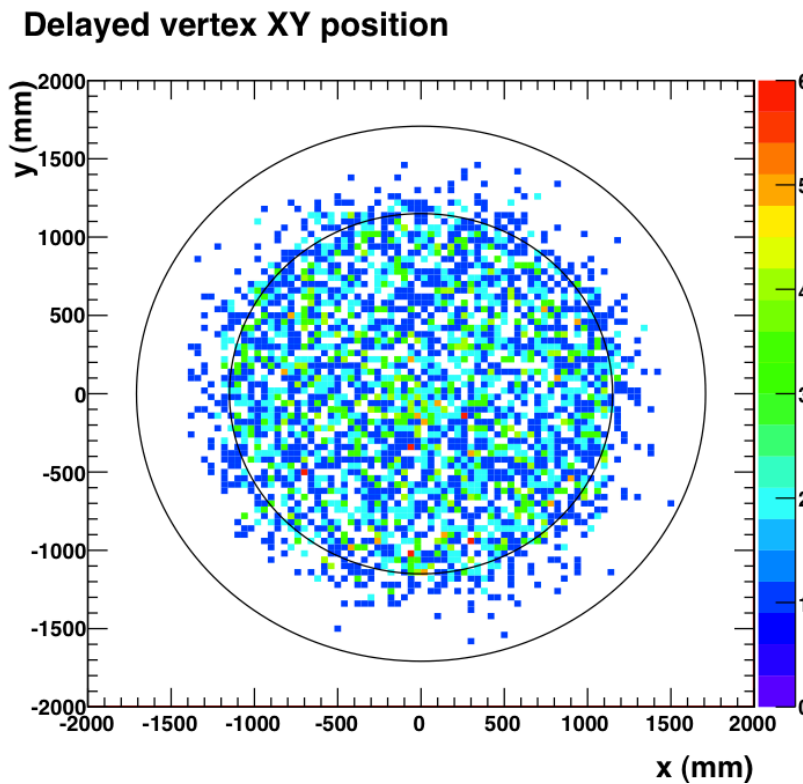


Chooz reactors A+B:
➤ 8.5 GW_{thermal}

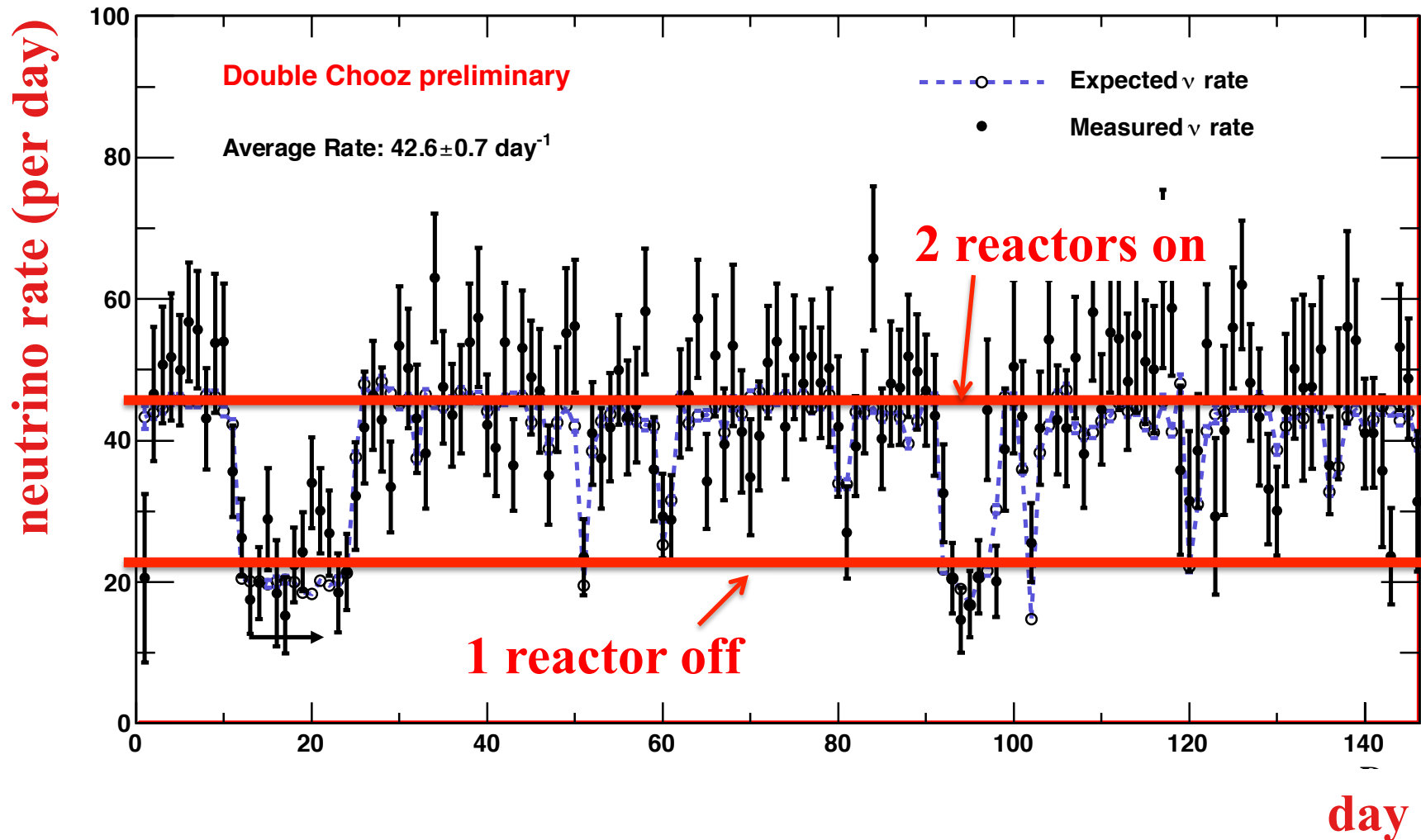


Neutrino Candidates n-Vertex Distribution

- Neutrino + Background selection
 - Prompt Energy Deposition
 - Delayed Energy Deposition
 - Time Correlation
 - No Muons



Measured Neutrino Rate & Reactors



Both reactors off for about 24 hours → pure background measurement

First Results

Mid April - mid September 2011

- 4121 neutrino candidates
- 328 background events
- 4041 neutrinos expected without oscillations

➤ 3793 neutrinos seen (94 %)

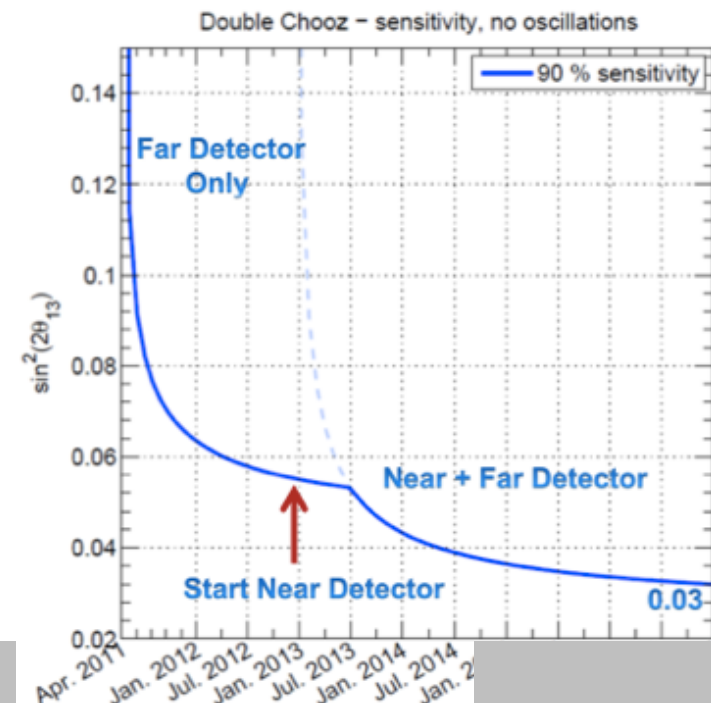
➔ $\sin^2(2\theta_{13}) = 0.085 \pm 0.051$
➔ >90% probability for $\theta_{13} > 0$

➔ More data until next spring
(5 months → 1 year)

➔ Further improved precision
with 2 detectors



H.de Kerret, Seoul, 9.November '11



Further Results

Electron neutrinos from a ν_μ -beam

T2K: $p \rightarrow$ target @ J-PARC

Detector: Superkamiokande (50 kt)

MIONS: Beam from FNAL to Soudan

Daya Bay: Reactor experiment

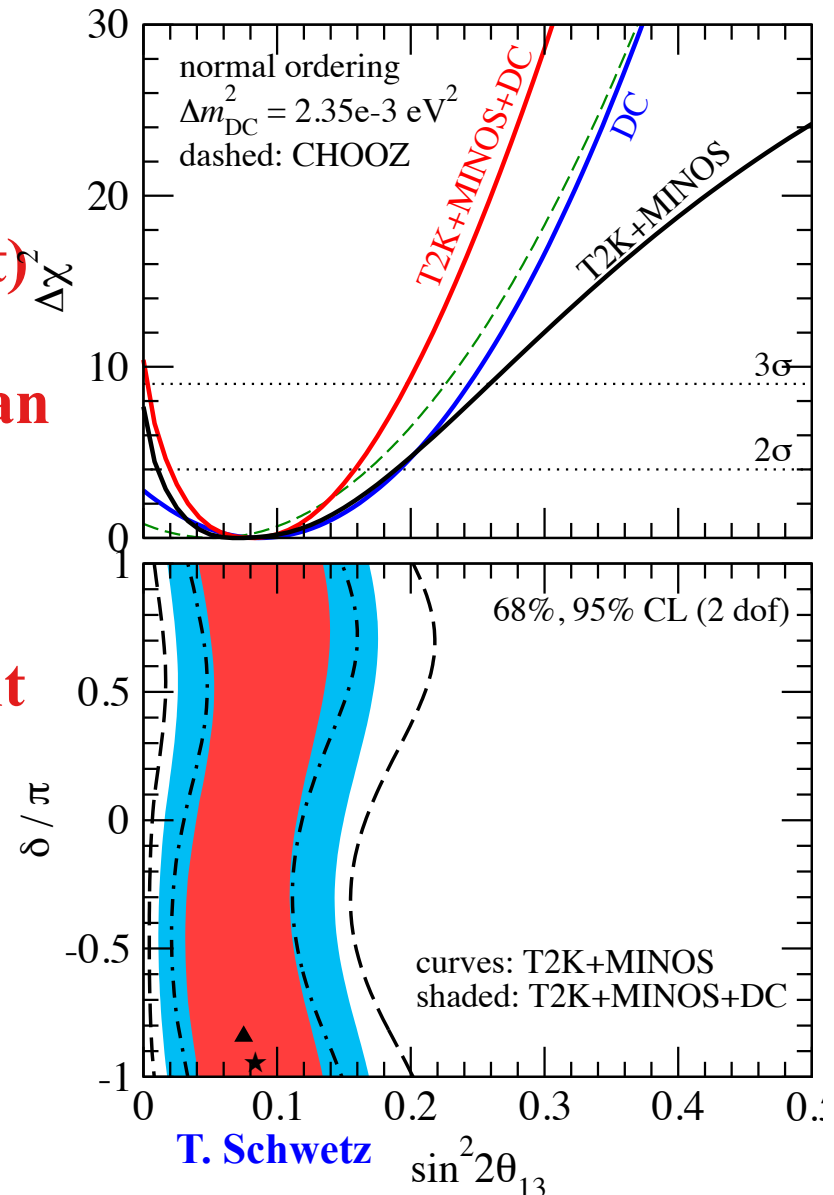
RENO: Reactor neutrino experiment

Double Chooz:

$$\sin^2(2\theta_{13}) = 0.085 \pm 0.051$$

will further improve very soon

All experiments: many sigma, CP...



Implications of sizable $\sin^2(2\theta_{13})$

... in the standard 3 neutrino picture:

- **Experiment:**

- good chances to measure θ_{13} soon with good precision
- leptonic CP violation (Dirac phase δ) accessible

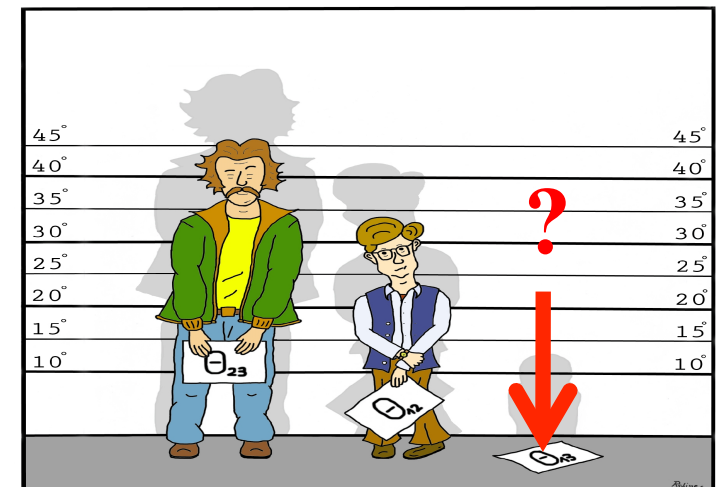
$$U_{\text{PMNS}} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix} \times \begin{pmatrix} c_{13} & 0 & s_{13}e^{-i\delta} \\ 0 & 1 & 0 \\ -s_{13}e^{+i\delta} & 0 & c_{13} \end{pmatrix} \times \begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix} \times \begin{pmatrix} 1 & 0 & 0 \\ 0 & e^{i\alpha_1} & 0 \\ 0 & 0 & e^{i\alpha_2} \end{pmatrix}$$

- standard neutrino factory out – matter uncertainties too big
- low energy NF or beta beam or conventional beams

- **Theory:**

- no apparent necessity for a special θ_{13} suppression (symmetry, ...)
- TBM ruled out \rightarrow sizable corrections, other starting points, ...

- ...maybe not the complete story



Potential Neutrinos Oscillation Surprises

... various assumptions: Majorana, 3 ν 's, $\Delta L=2$ ops mass mechanism

→ examples:

- extra light sterile neutrinos → C. Rubbia's lectures
- how NSI's can fool precision experiments

GLOBES simulation:

Source	⊗	Oscillation	⊗	Detector
<ul style="list-style-type: none">- neutrino energy E- flux and spectrum- flavour composition- contamination- symmetric $\nu/\bar{\nu}$ operation		<ul style="list-style-type: none">- oscillation channels- realistic baselines- MSW matter profile- degeneracies- correlations		<ul style="list-style-type: none">- effective mass, material- threshold, resolution- particle ID (flavour, charge, event reconstruction, ...)- backgrounds- x-sections (at low E)

new effects beyond oscillations?!

NSI Operators

- **Good reasons for physics beyond the SM+ (with ν 's)**
 - expect effects beyond 3 flavours in many models
 - effective 4f interactions

$$\mathcal{L}_{NSI} \simeq \epsilon_{\alpha\beta} 2\sqrt{2}G_F (\bar{\nu}_{L\beta} \gamma^\rho \nu_{L\alpha}) (\bar{f}_L \gamma_\rho f_L)$$

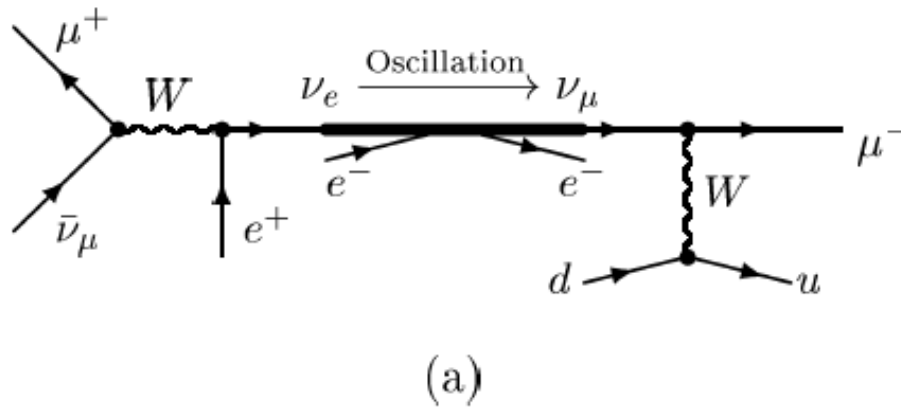
- **integrating out heavy physics (c.f. $G_F \leftrightarrow M_W$)**

$$|\epsilon| \simeq \frac{M_W^2}{M_{NSI}^2}$$

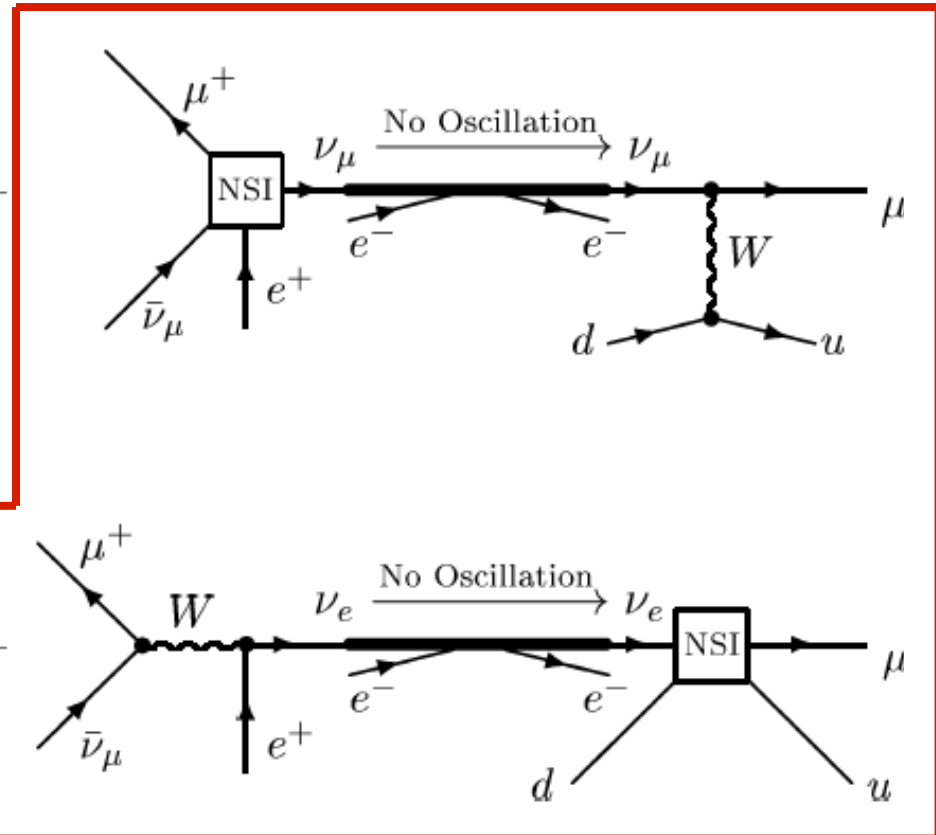
Grossman, Bergmann+Grossman, Ota+Sato, Honda et al., Friedland+Lunardini, Blennow+Ohlsson+Skrotzki, Huber+Valle, Huber+Schwetz+Valle, Campanelli+Romanino, Bueno et al., Kopp+ML+Ota, ...

NSIs interfere with Oscillations

the “golden” oscillation channel

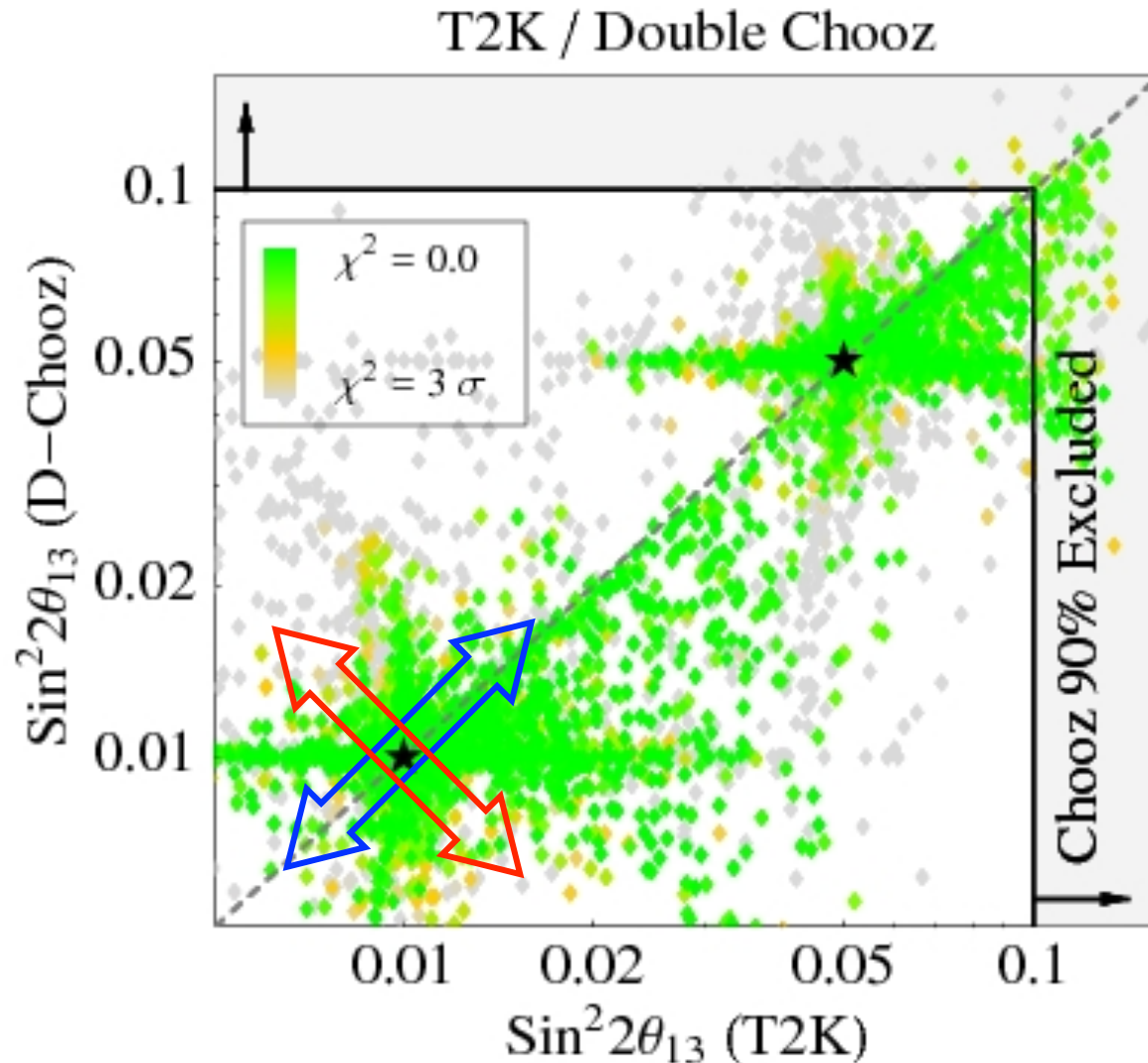


NSI contributions to the “golden” channel



note: interference in oscillations $\sim \epsilon$ \leftrightarrow FCNC effects $\sim \epsilon^2$

NSI: Offset and Mismatch in θ_{13}



Redundant measurements:

Double Chooz + T2K

***=assumed 'true' values of θ_{13}**

scatter-plot: ϵ values random

- below existing bounds

- random phases

NSIs can lead to:

- **offset**

- **mismatch**

➔ **redundancy**

➔ **interesting potential**

Four Methods of Mass Determination

- **kinematical**
- **lepton number violation**
 \leftrightarrow **Majorana nature**
- **oscillations**
- **astrophysics & cosmology**

Neutrinos & Cosmology

- Dark Matter ~ 25% & Dark Energy 70%
- mass of all neutrinos: $0.001 \leq \Omega_\nu \leq 0.02$
- baryonic matter $\Omega_B \sim 0.04$

Neutrino mass contribution
possibly as big as all baryonic matter \gg visible matter
much more COLD dark matter & dark energy
neutrinos are an important hot dark matter component

Comological impact of neutrinos:

- hot component in structure formation: $330\nu/\text{cm}^3 \times \text{mass} \rightarrow$
- Big Bang Nucleosynthesis \rightarrow
- Baryon asymmetry \rightarrow Leptogenesis $\rightarrow ?$

-...

Source: Robert Kirshner

Source: David Aguilar, Harvard-Smithsonian Center for Astrophysics

Neutrinos and Dark Matter

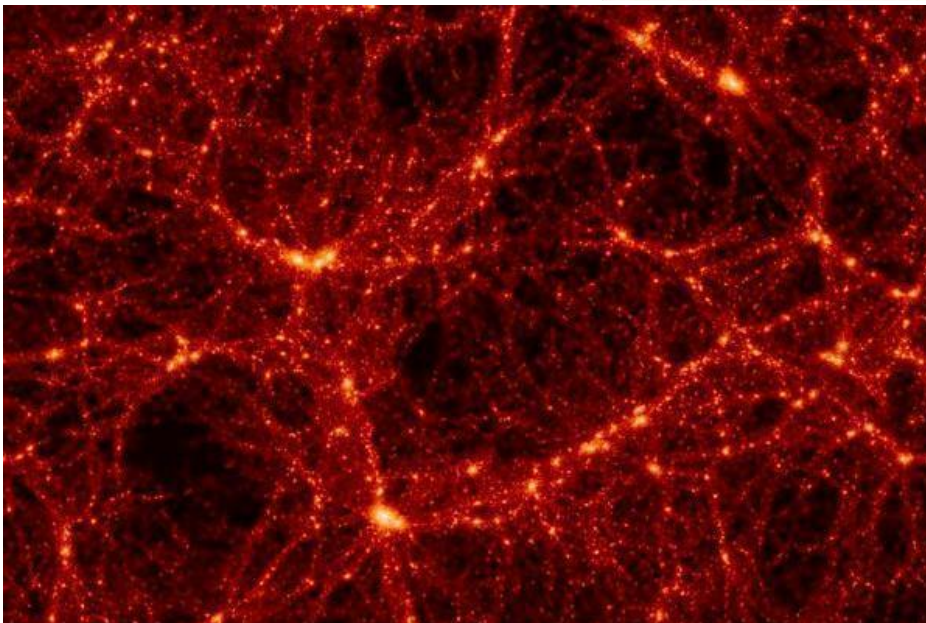
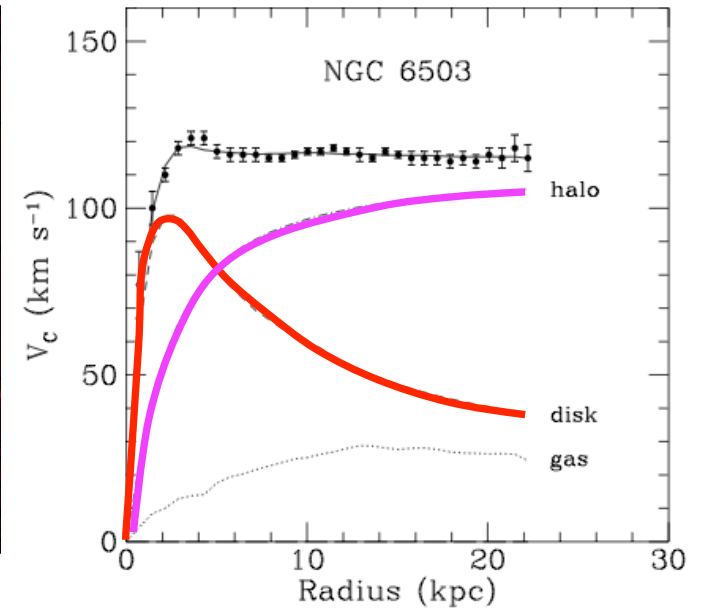
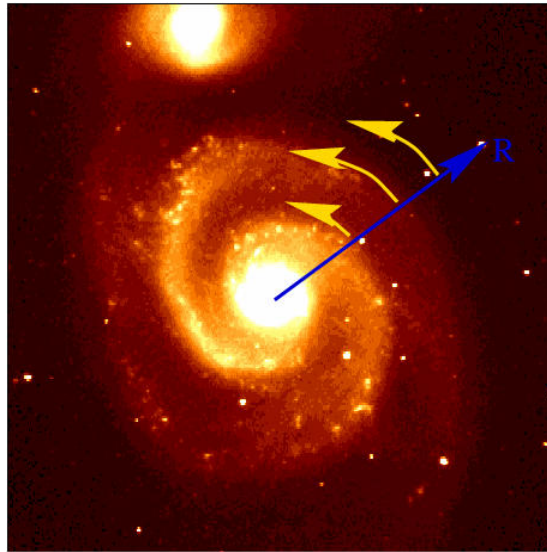
Dynamical Evidence for Dark Matter

F. Zwicky 1933

→ proplem with galaxy clusters

→ galaxies →

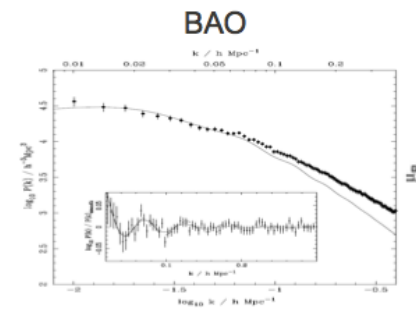
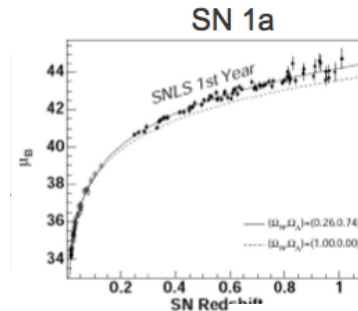
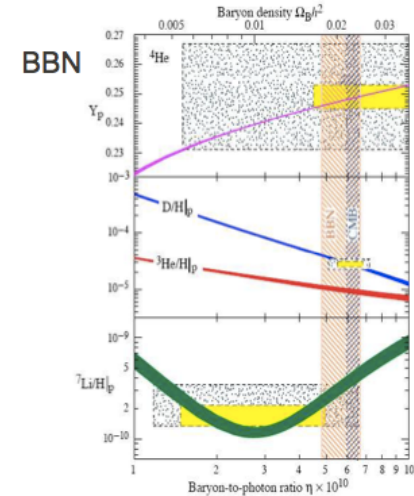
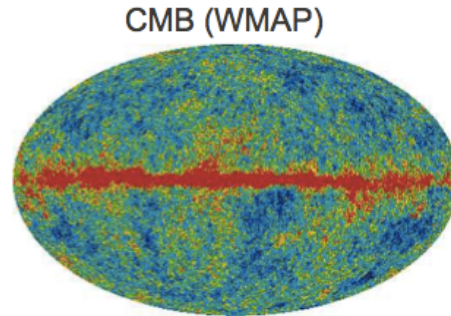
→ galaxies have a large DM halo



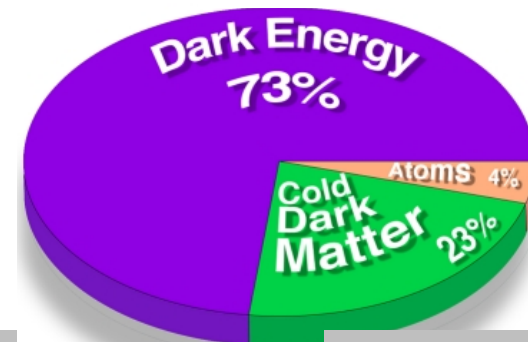
comparison of simulated and real structures
→ cosmological DM dynamics

Many consistent Evidences for Dark Matter

- + Galactic rotation curves
- + Galaxy clusters & GR lensing
- + Bullet Cluster
- + Velocity dispersions of galaxies
- + Cosmic microwave background
- + Sky Surveys and Baryon Acoustic Oscillations
- + Type Ia supernovae distance measurements
- + Big Bang Nucleosynthesis (BBN)
- + Lyman-alpha forest
- + Structure formation



→ Coherent picture of the composition of the Universe →



Sterile Neutrino Spectrum

The standard picture:

3 heavy sterile neutrinos typ. $\geq 10^{13}$ GeV

→ leptogenesis, role in GUTs, ...

some mechanism which makes
1 or 2 heavy sterile neutrinos light?

→ keV sterile neutrino

→ tiny heavy-light mixing

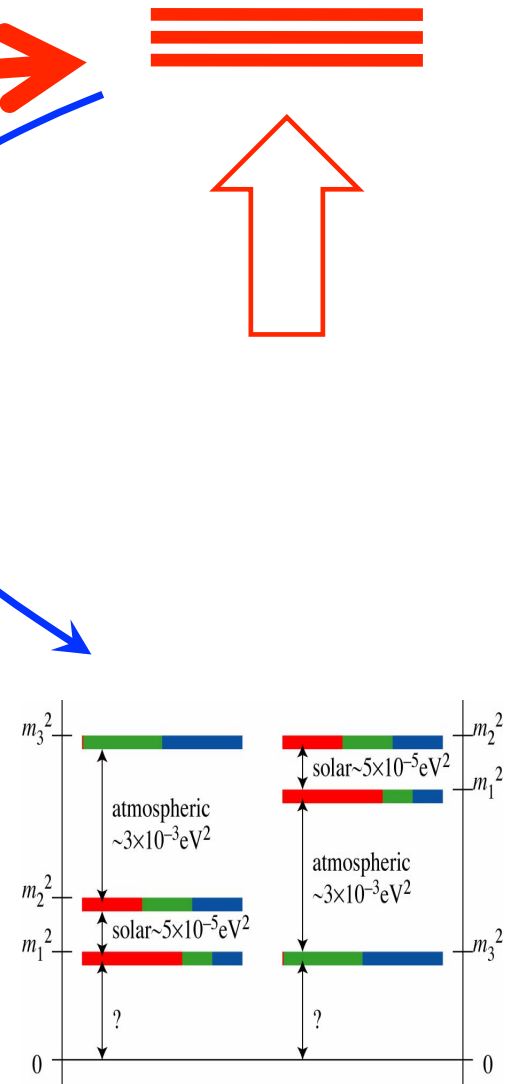
3 light active neutrinos

→ this could easily be wrong

- more than 3 N_R states, ...

- M_R may have special eigenvalues, ...

→ light sterile neutrinos ?!



Evidences for Light Sterile Neutrinos ?

Particle Physics:

Reactor anomaly, LSND, MiniBooNE, MINOS, Gallex...

- New and better data / experiments are needed to clarify the situation
- maybe something exciting around the corner?
- but eV scale and sizable mixings

CMB: extra eV-ish neutrinos J. Hamann et al. , ...

BBN: extra ν 's possible: $N_\nu \simeq 3.7 \pm 1$

E. Aver, K. Olive, E. Skillman (2010), Y. Izotov, T. Thuan(2010)

Astrophysics:

e.g. effects of keV-ish sterile neutrinos on pulsar kicks

Kusenko, Segre, Mocioiu, Pascoli, Fuller et al., Biermann & Kusenko, Stasielak et al., Loewenstein et al., Dodelson, Widrow, Dolgov, ...

Most likely not all of them are correct!

→ consequences! → assume first only 3 active neutrinos

Could Neutrinos be Dark Matter?

- Active neutrinos would be perfect Hot Dark Matter → ruled out:
 - destroys small scale structures in cosmological evolution
 - measured neutrino masses too small → maybe HDM component
- keV sterile neutrinos: Warm Dark Matter → workes very well:
 - relativistic at decoupling
 - non-relativistic at radiation to matter dominance transition
 - OK for $M_X \simeq \text{few keV}$ with very tiny mixing
 - reduced small scale structure → smoother profile, less dwarf satellites
 - scenario where one sterile neutrino is keV-ish, the others heavy
 - tiny active – sterile mixings $O(m_\nu/M_R)$
 - ↔ observational hints from astronomy
 - hints that a keV sterile particle may exist → right-handed neutrino?

Note: Right-handed neutrinos exist probably anyway – just make one light!

keV Neutrinos as WDM

The ν MSM

Asaka, Blanchet, Shaposhnikov, 2005 Asaka, Shaposhnikov, 2005

Particle content:

- Gauge fields of $SU(3)_c \times SU(2)_W \times U(1)_Y$: γ, W_{\pm}, Z, g
- Higgs doublet: $\Phi=(1,2,1)$

• Matter

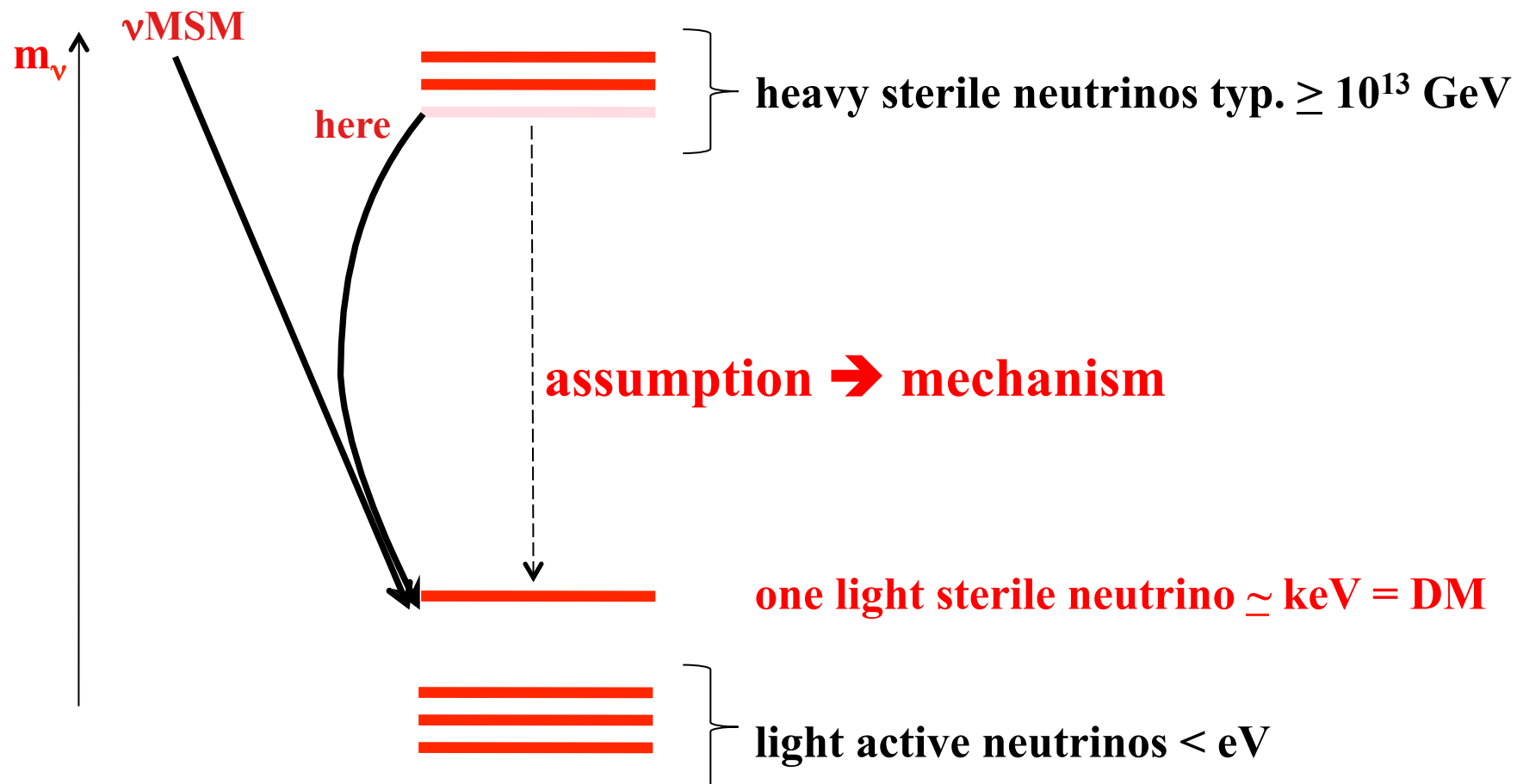
	$SU(3)_c$	$SU(2)_W$	$U(1)_Y$	$U(1)_{em}$
$\begin{pmatrix} u \\ d \end{pmatrix}_L$	3	2	+1/3	$\begin{pmatrix} +2/3 \\ -1/3 \end{pmatrix}$
u_R	3	1	+4/3	+2/3
d_R	3	1	-2/3	-1/3
$\begin{pmatrix} \nu_e \\ e \end{pmatrix}_L$	1	2	-1	$\begin{pmatrix} 0 \\ -1 \end{pmatrix}$
e_R	1	1	-2	-1
N	1	1	0	0

x3 generations

- lepton sector more symmetric to the quark sector
- Majorana masses for N
- choose for one sterile $\nu \sim \text{keV}$ mass → exceeds lifetime of Universe

Virtue and Problem of the ν MSM

- ν MSM:** Scenario with sterile ν and tiny mixing \rightarrow never enters thermal equilibrium
- \rightarrow requires **non-thermal production** from other particles (avoid over-closure)
 - \rightarrow **new physics** before the beginning of the thermal evolution sets abundance



Alternative Scenario with Thermal Abundance

An alternative scenario: Bezrukov, Hettmannsperger, ML

- Three right-handed neutrinos N_1, N_2, N_3
- Dirac and Majorana mass terms
- **N Charged under some (BSM) gauge group \rightarrow scale M (\sim sterile)**
- **Specific example: LR-symmetry $SU(3)_c \times SU(2)_L \times SU(2)_R \times U(1)_{B-L}$**

Roles played by the sterile (\sim right-handed) neutrinos:

N_1 – Warm Dark Matter

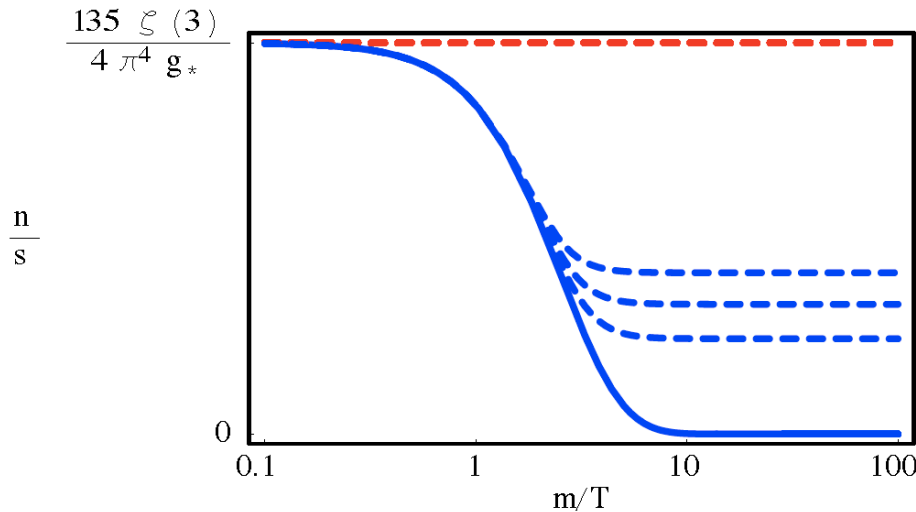
- Mass $M_1 \sim \text{keV}$
- Lifetime $\tau_1 > \tau_{\text{Universe}} \sim 10^{17} \text{ s}$

$N_{2,3}$ – dilute entropy after DM decoupling

- Mass $M_{2,3} > \text{GeV}$
- Lifetime $\tau_{2,3} \lesssim 0.1 \text{ s}$

Obtaining the correct Abundance

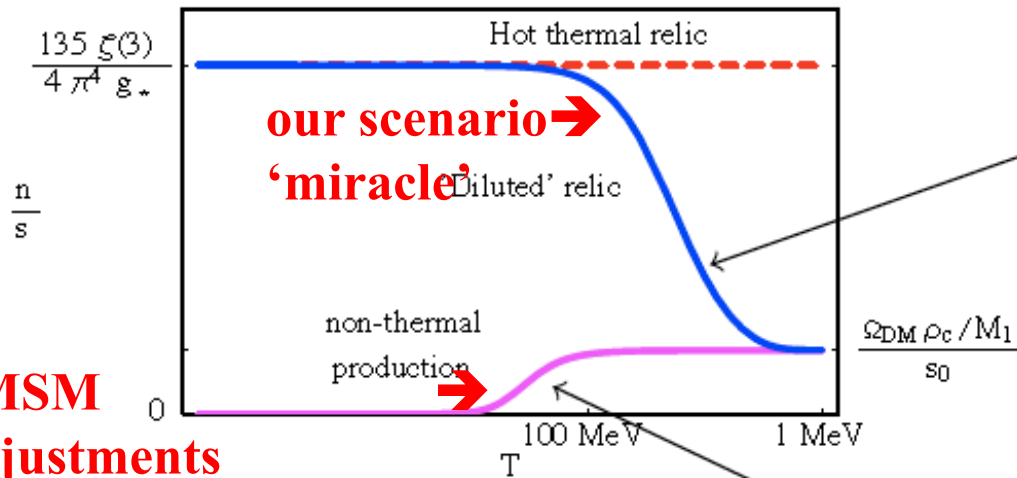
Usual thermal case:



HDM: $\frac{\Omega}{\Omega_{DM}} \simeq \left(\frac{10}{g_{*f}}\right) \left(\frac{M}{10\text{eV}}\right)$
Decoupled relativistic

CDM: $\Omega \sim \Omega_{DM}$
($M \gg \text{MeV}$)
Decoupled nonrelativistic

keV sterile neutrinos:



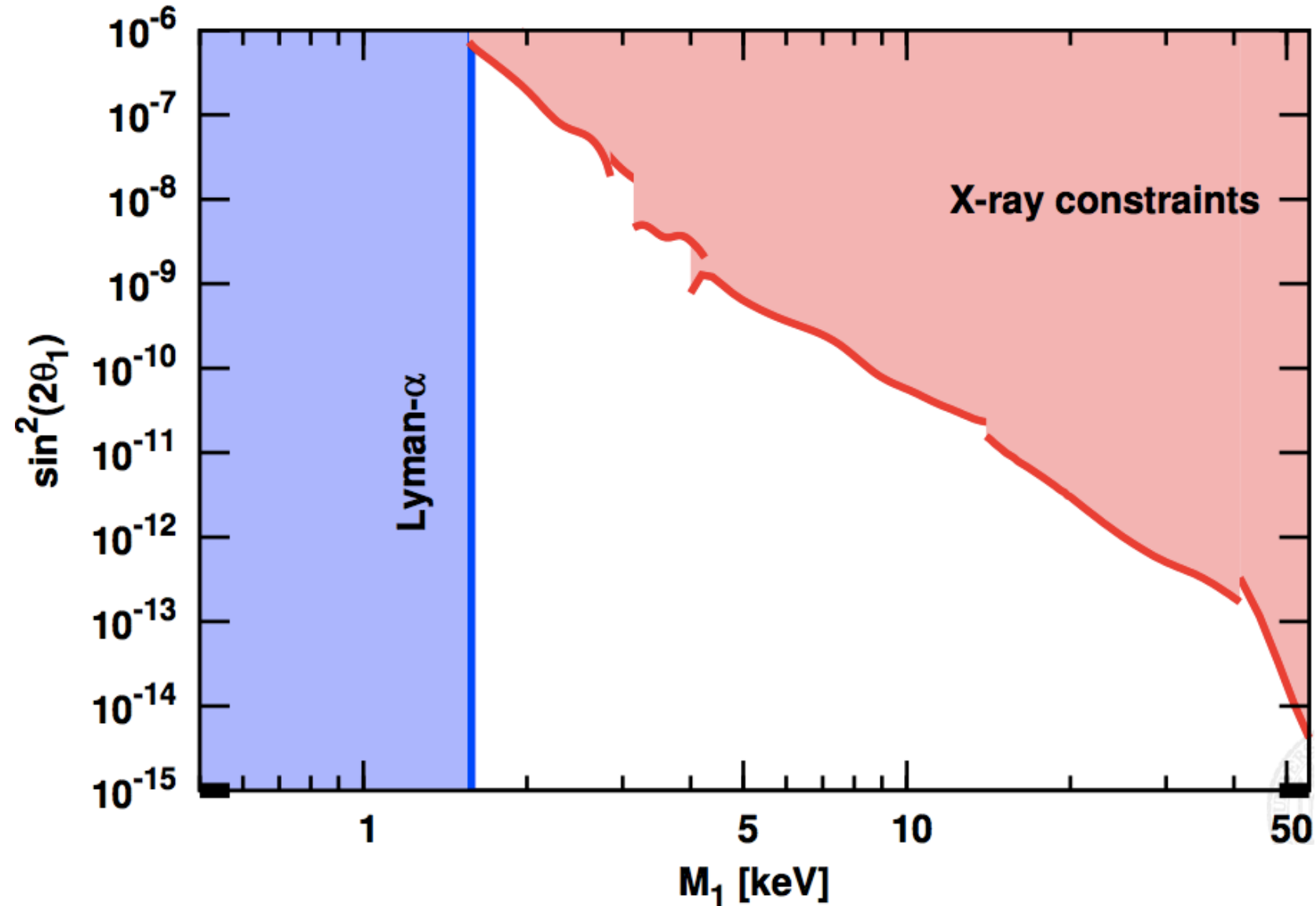
ν MSM adjustments

Diluted after decoupling
(entropy generated by other particle decay)

$$\Omega \sim \Omega_{DM}$$

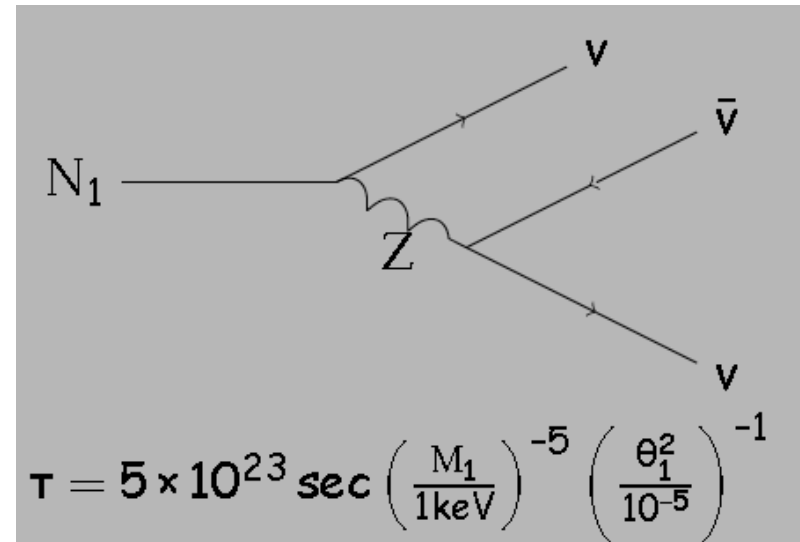
Never entered thermal equilibrium

Allowed Parameter Range

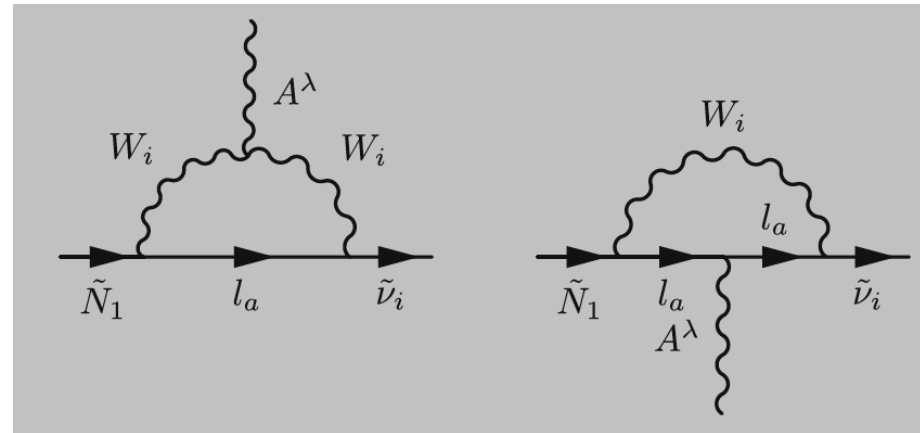


Observing keV-ish Neutrino DM

- **LHC**
 - sterile neutrino DM is not observable
 - WIMP-like particles still possible – but not DM
- **direct searches**
 - sterile ν DM extremely difficult; maybe in β -decay (MARE)
- **astrophysics/cosmology** \rightarrow at some level: keV X-rays
 - \rightarrow sterile neutrino DM is decaying into active neutrinos
 - decay $N_1 \rightarrow \nu\bar{\nu}$, $N_1 \rightarrow \nu\nu$
 - not very constraining since $\tau \gg \tau_{\text{Universe}}$



- radiative decays $N_1 \rightarrow \nu\gamma$



- so far: observational limit on active-sterile mixing angle

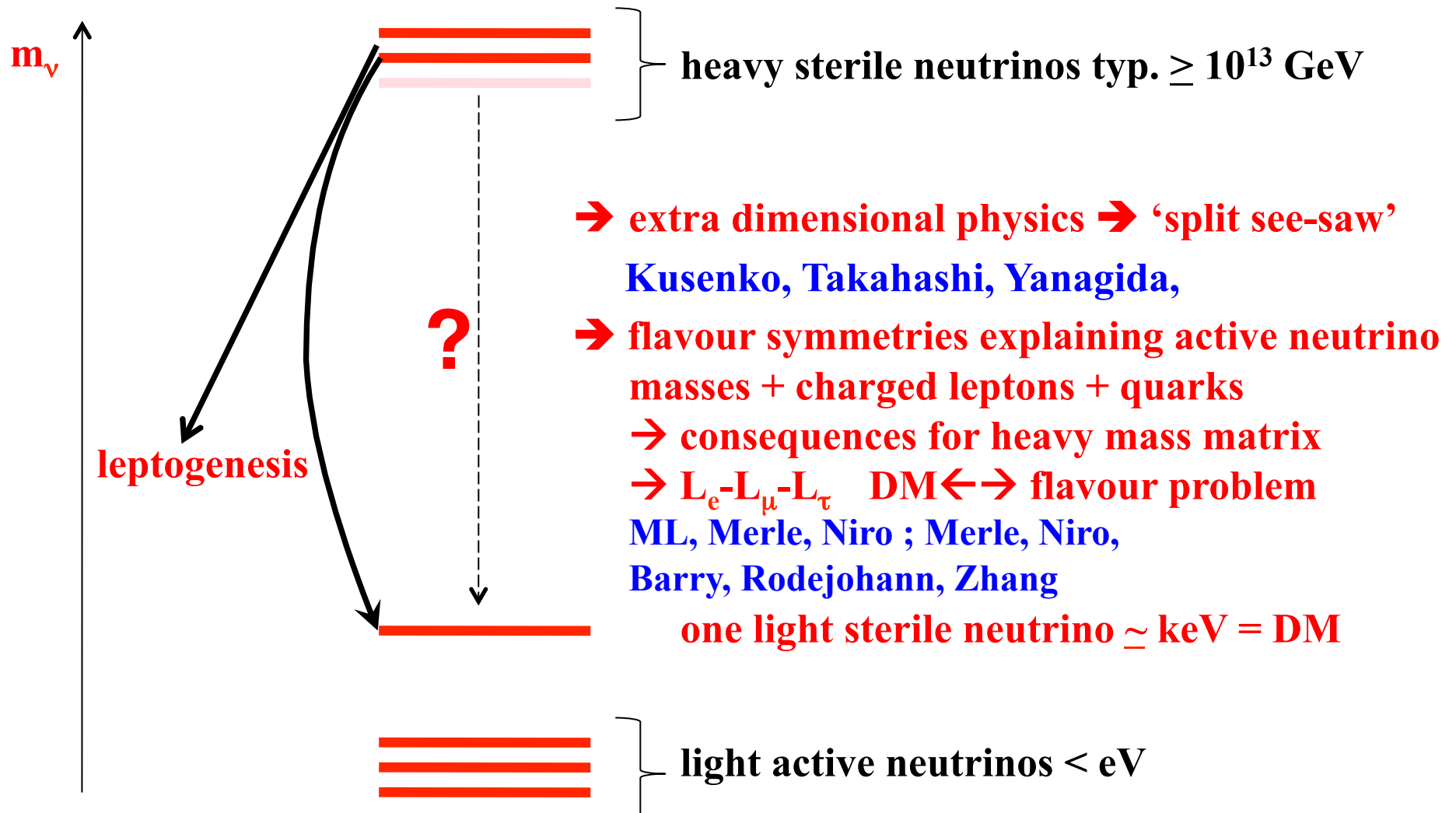
$$\Gamma_{N_1 \rightarrow \nu\gamma} \simeq 5.5 \times 10^{-22} \theta_1^2 \left(\frac{M_1}{1 \text{ keV}} \right)^5 \text{ s}^{-1}$$

$$\theta_1^2 \lesssim 1.8 \times 10^{-5} \left(\frac{1 \text{ keV}}{M_1} \right)^5$$

- mixing tiny, but naturally expected to be tiny: $O(\text{scale ratio})$

Explaining keV-ish Sterile Neutrinos

Possible scenario: See-saw + a reason why 1 sterile ν is light



Conclusions

- **Dark Matter exists**
 - **But we cannot be sure if Dark Matter is**
 - **made of particles (→ smarter modified gravity ?)**
 - **particles → prefer scenarios with a ‘miracle’ (abundance)**
 - **WIMPs (SUSY, ...)** - are very well motivated
 - **GIMPs – example: a keV-ish sterile neutrino**
very well motivated and good working **Warm Dark Matter**
related to ν -masses
- WIMP detection/limits down to 10^{-47}cm^2 will be reached**
- **either find a WIMP or very strong constraint!**
 - **combined with LHC searches**
 - **excellent discovery potential!**

Summary & Conclusions

- neutrinos provide **unique information & insight into new** physics
- three active neutrinos are still the paradigm
 - ➔ Possibility to measure **leptonic CP-violation** \leftrightarrow BAU...
- But maybe more...
 - ➔ **sterile neutrinos theoretically conceivable & quite natural**
- various **hints** which may point towards new surprises
 - ➔ requires clarification
 - ➔ **very interesting consequences**
 - neutrino mass generation
 - neutrinos as Dark Matter

Neutrinos have always been surprising
➔ **very interesting ν -times ahead!**

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