

# Signals of asymmetric dark matter from galactic center pulsar implosions

VIA Lecture

Joseph Bramante  
Feb. 20, 2015

Dark matter may soon be discovered in nuclear/  
electron recoils, nuclear/electron collisions, or  
electromagnetic oscillations.

Dark matter may soon be discovered in nuclear/  
electron recoils, nuclear/electron collisions, or  
electromagnetic oscillations.

But in order to fundamentally advance  
understanding of the universe's symmetries, we  
must determine what is stabilizing the dark matter  
abundance and how it was primordially produced.

The hardest task will be to identify and categorize the interactions dark matter may have with other dark particles.

There are many possible handles on this — the collision of galaxies in bullet clusters, the inferred velocity dispersion of dwarf, spiral, and elliptic halos, and acoustic oscillations in and the normalization of the primordial power spectrum.

Asymmetric dark matter models suppose that dark matter in the halo, like ordinary matter in the halo, is composed of particles and not antiparticles charged under a continuous symmetry. This asymmetric abundance precludes  $X$ - $X$  (as opposed to  $X^*$ - $X$ ) annihilations.

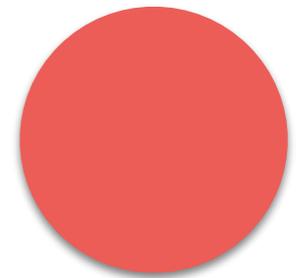
Asymmetric dark matter models suppose that dark matter in the halo, like ordinary matter in the halo, is composed of particles and not antiparticles charged under a continuous symmetry. This asymmetric abundance precludes  $X$ - $X$  (as opposed to  $X^*$ - $X$ ) annihilations.

In a very futuristic setting, the interactions (C,CP violating) that lead to a dark asymmetry might be revealed through something like a pulsed collider + kamioka program.

This talk suggests that on a much shorter timescale, old pulsars in the galactic center imploding into black holes could be a signal of asymmetric dark matter. This signal relies on the fact that annihilating dark matter cannot copiously collect in pulsars to form black holes.

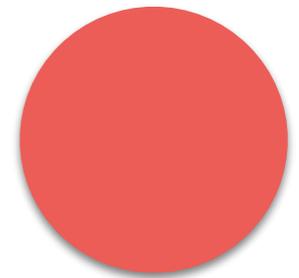
A neutron star is a ball of fermions formed from the supernova of a 10 solar mass progenitor star.

$$N_f = \frac{m_{pl}^3}{m_X^3}$$



A neutron star is a ball of fermions formed from the supernova of a 10 solar mass progenitor star.

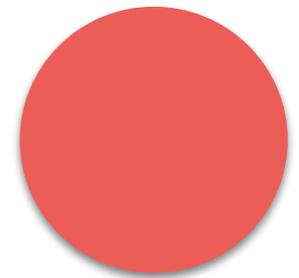
$$N_f = \frac{m_{pl}^3}{m_X^3}$$



The bound above can be compared to the same limit on a ball of bosons.

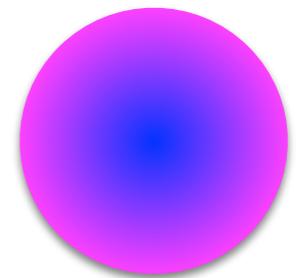
A neutron star is a ball of fermions formed from the supernova of a 10 solar mass progenitor star.

$$N_f = \frac{m_{pl}^3}{m_X^3}$$



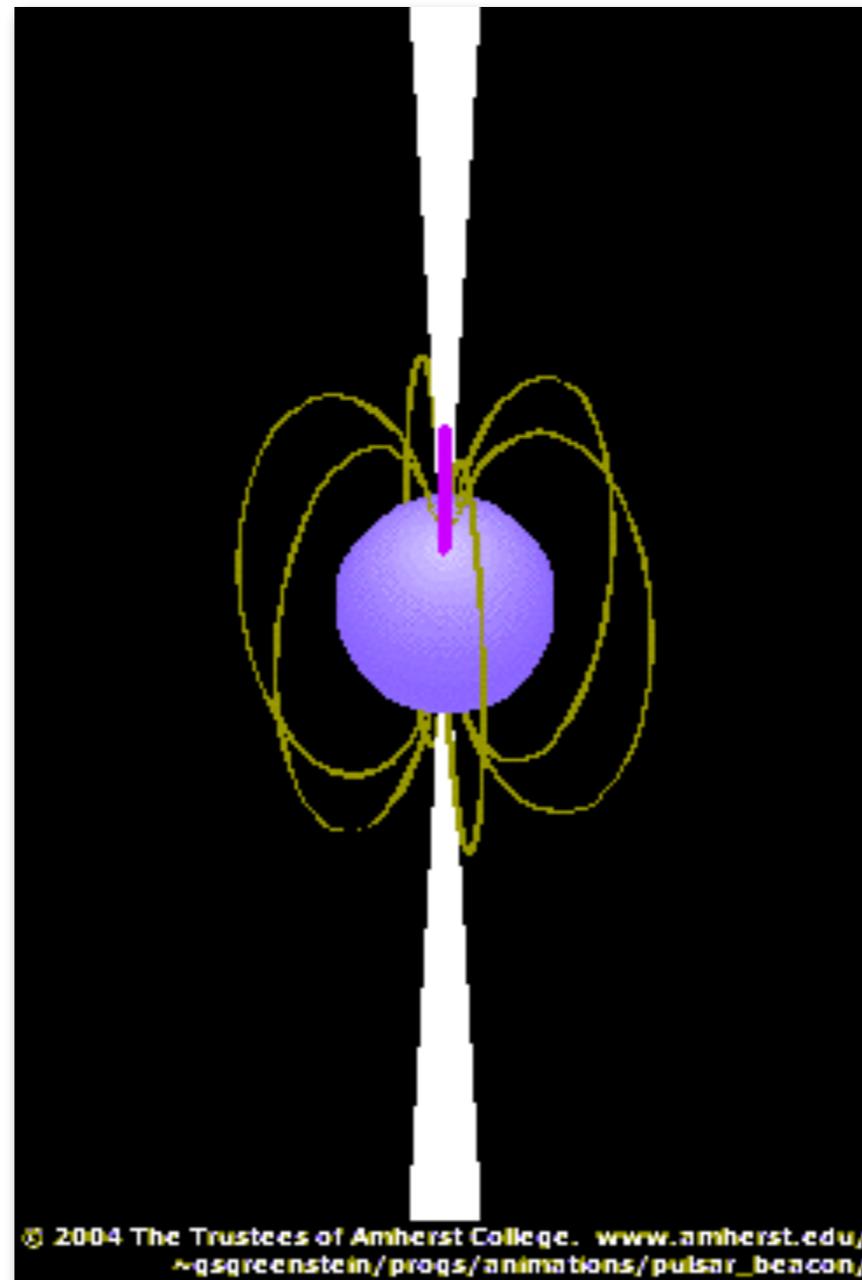
The bound above can be compared to the same limit on a ball of bosons.

$$N_b = \frac{m_{pl}^2}{m_X^2}$$



Of course, these collapse bounds will apply to both visible and dark matter.

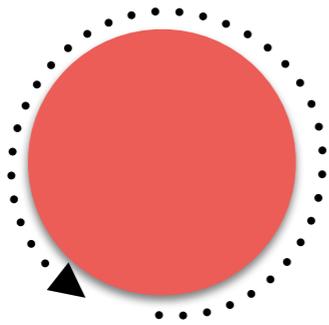
A pulsar is a spinning neutron star, observed as a regularly radiating revolving magnetic dipole.



© 2004 The Trustees of Amherst College. [www.amherst.edu/~gsqgreenstein/progs/animations/pulsar\\_beacon/](http://www.amherst.edu/~gsqgreenstein/progs/animations/pulsar_beacon/)

A pulsar is a spinning neutron star, observed as a regularly radiating revolving magnetic dipole.

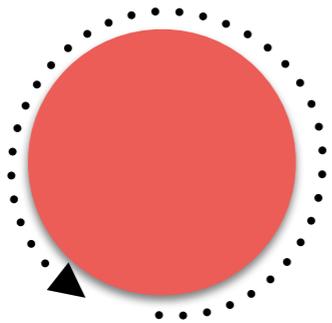
How old is it?



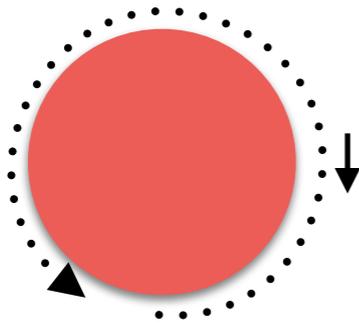
P

A pulsar is a spinning neutron star, observed as a regularly radiating revolving magnetic dipole.

How old is it?



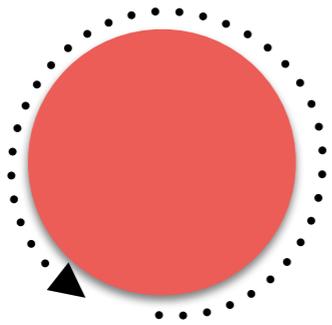
$P$



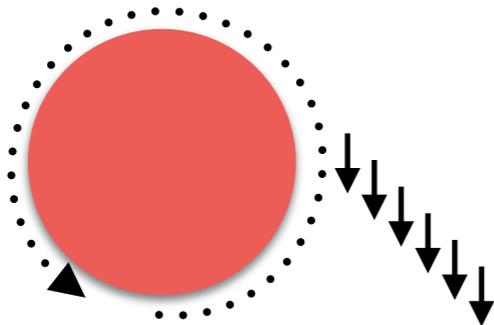
$\dot{P}$

A pulsar is a spinning neutron star, observed as a regularly radiating revolving magnetic dipole.

How old is it?



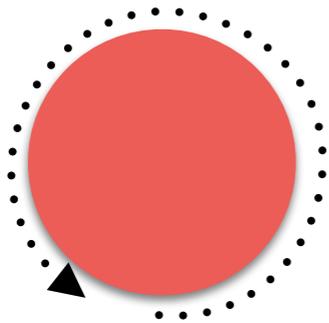
$P$



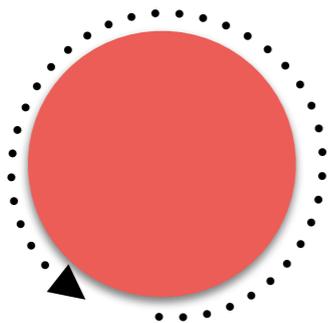
$\dot{P}$

A pulsar is a spinning neutron star, observed as a regularly radiating revolving magnetic dipole.

How old is it?



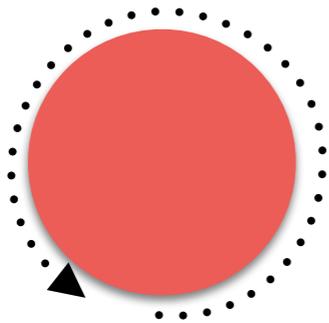
$P$



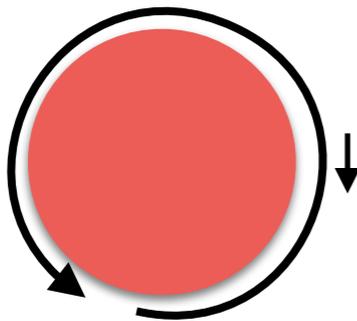
$\dot{P}$

A pulsar is a spinning neutron star, observed as a regularly radiating revolving magnetic dipole.

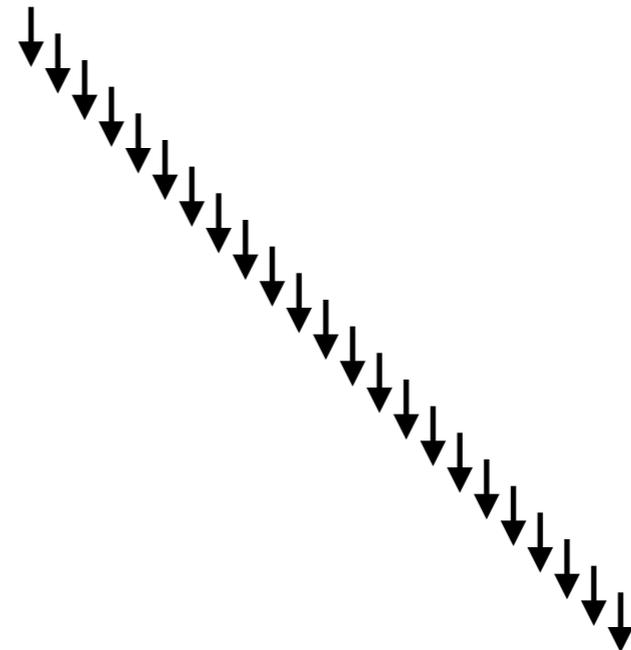
How old is it?



$P$

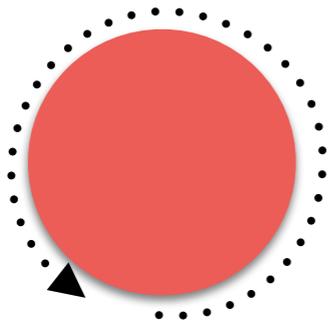


$\dot{P}$

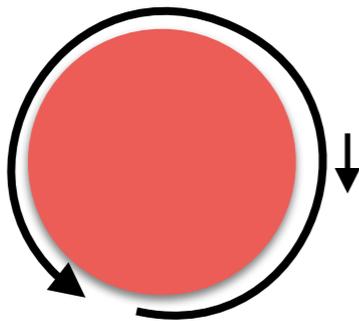


A pulsar is a spinning neutron star, observed as a regularly radiating revolving magnetic dipole.

How old is it?



$P$

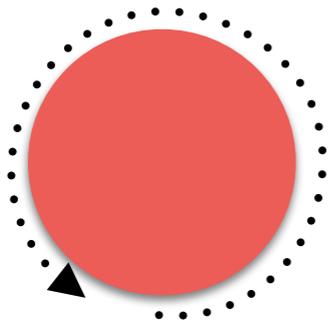


$\dot{P}$

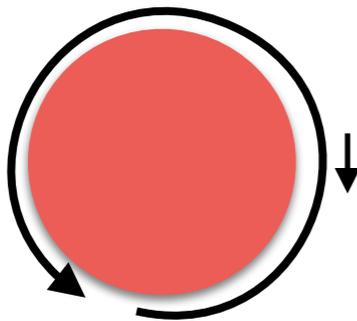
divided by

A pulsar is a spinning neutron star, observed as a regularly radiating revolving magnetic dipole.

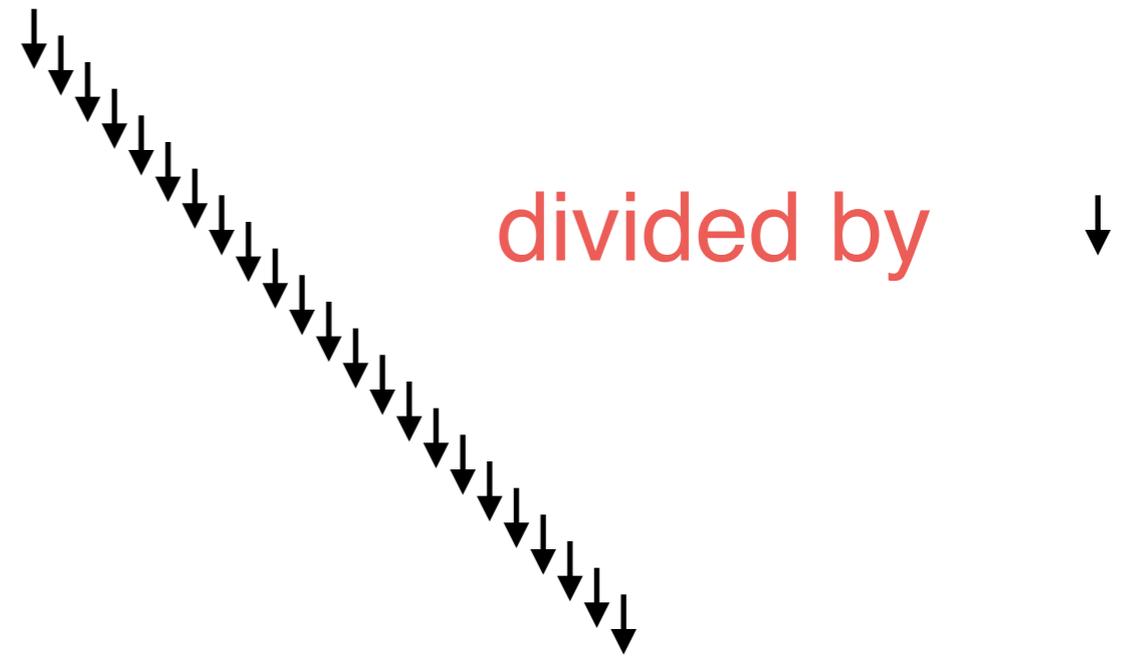
How old is it?



$P$

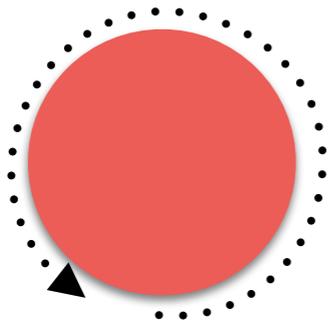


$\dot{P}$

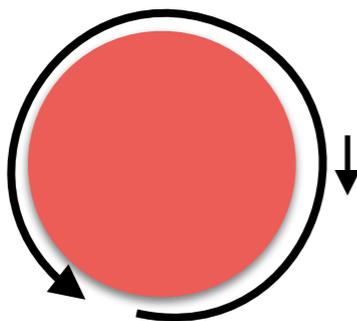


A pulsar is a spinning neutron star, observed as a regularly radiating revolving magnetic dipole.

How old is it?



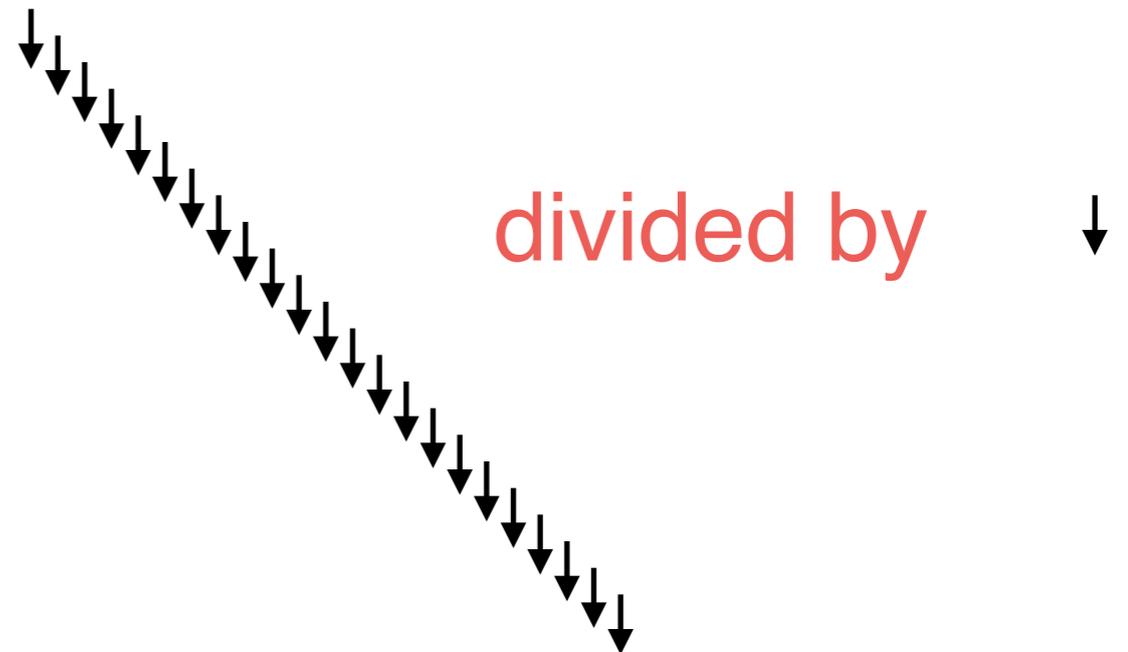
$P$



$\dot{P}$

$$t_{NS} = \frac{P}{2\dot{P}}$$

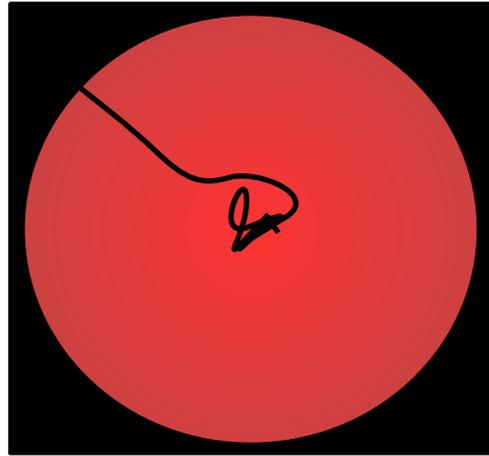
divided by



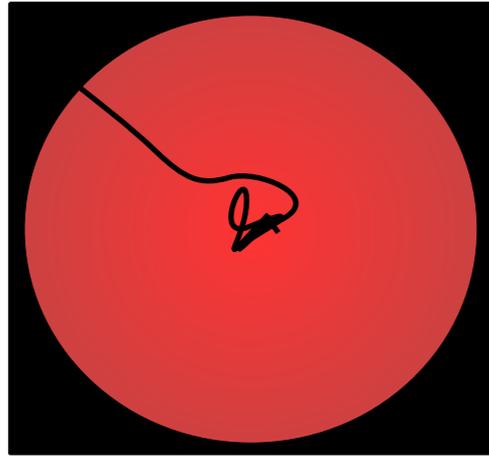
Now we know how old pulsars are, and that either Heisenberg or Fermi pressure stabilizes compact matter.

We can now discuss pulsar-destroying black holes formed from dark matter collected in pulsars.

1] DM captured

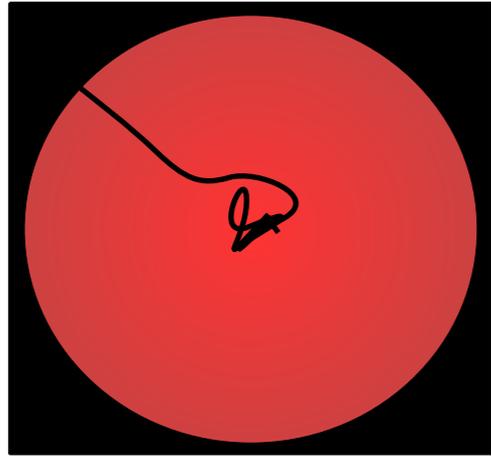


1] DM captured

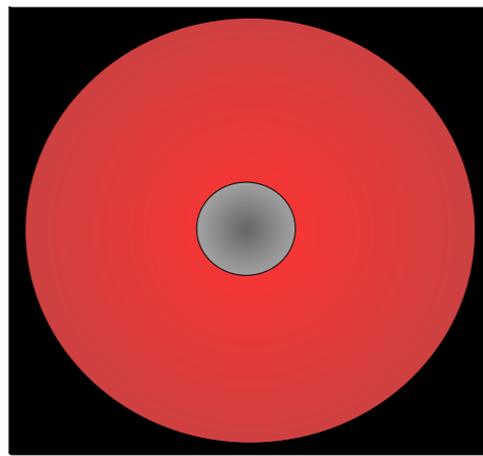


$$C_X \propto \frac{\rho_X}{\bar{v}} \sigma_{nX}$$

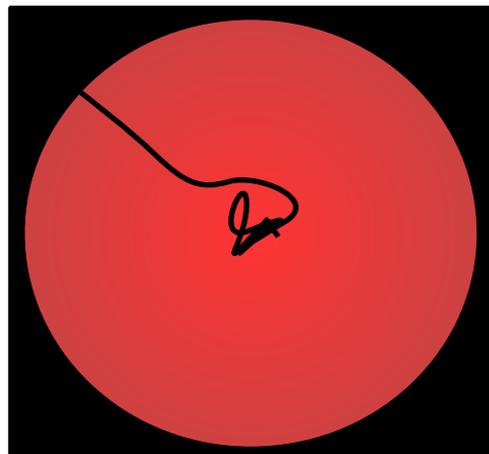
1] DM captured



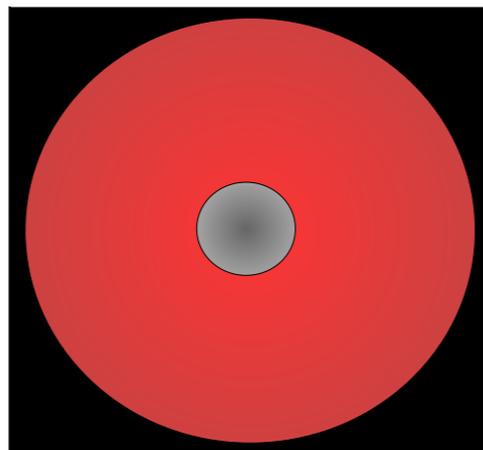
2] DM thermalizes



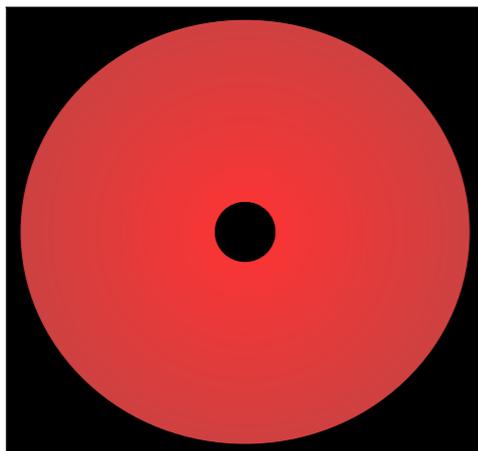
1] DM captured



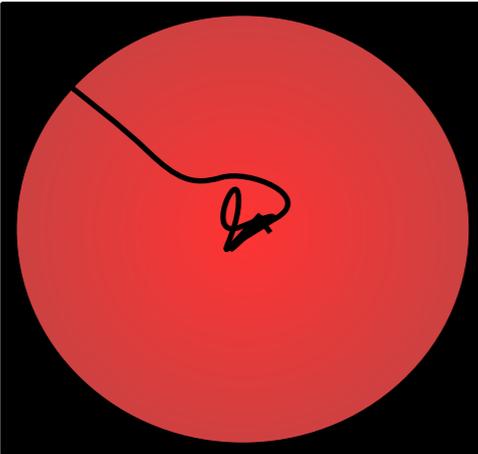
2] DM thermalizes



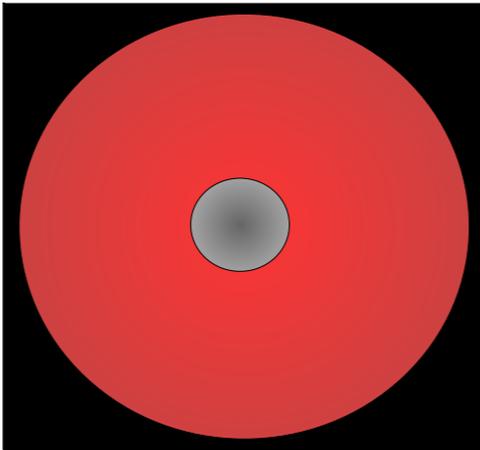
3] DM collapses



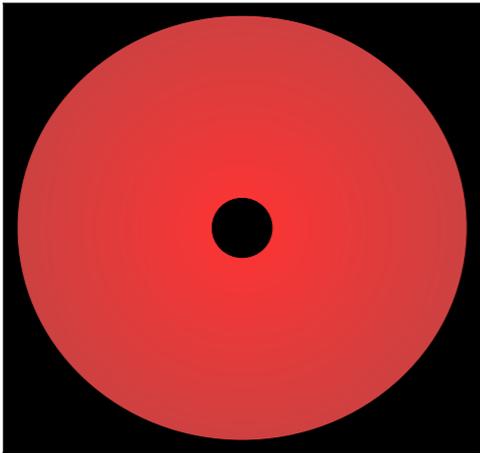
1] DM captured



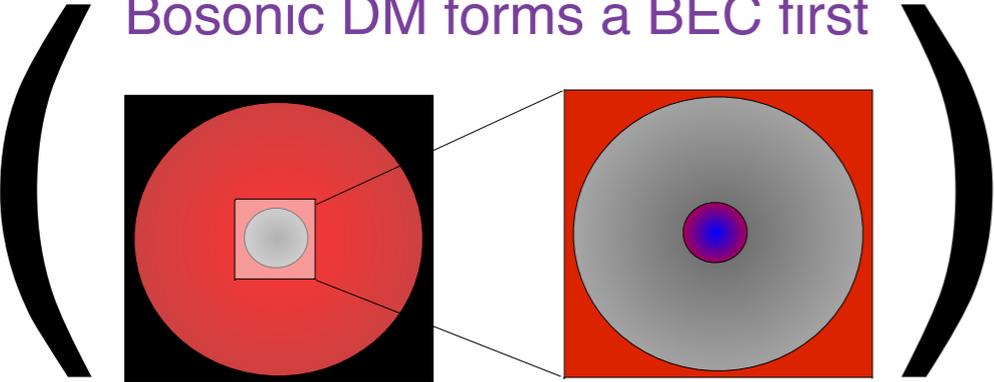
2] DM thermalizes



3] DM collapses



Bosonic DM forms a BEC first



A ball of fermions has fermi degeneracy pressure, but attractive self-interactions can counteract this.

$$N_f = \frac{m_{pl}^3}{m_X^3}$$

A ball of fermions has fermi degeneracy pressure, but attractive self-interactions can counteract this.

$$N_f < \frac{m_{pl}^3}{m_X^3}$$

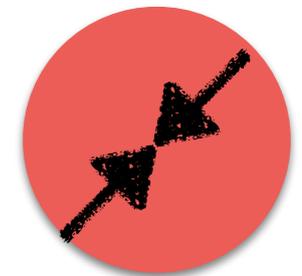
A ball of fermions has fermi degeneracy pressure,  
but attractive self-interactions can counteract this.

$$N_f < \frac{m_{\text{pl}}^3}{m_{\text{X}}^3} \quad \text{if } \mathcal{L} \supset \alpha \phi \bar{\psi} \psi$$

A ball of fermions has fermi degeneracy pressure,  
but attractive self-interactions can counteract this.

$$N_f < \frac{m_{\text{pl}}^3}{m_{\text{X}}^3} \quad \text{if } \mathcal{L} \supset \alpha \phi \bar{\psi} \psi$$

attractive



A ball of fermions has fermi degeneracy pressure, but attractive self-interactions can counteract this.

$$N_f < \frac{m_{pl}^3}{m_X^3} \quad \text{if } \mathcal{L} \supset \alpha \phi \bar{\psi} \psi$$

attractive



A ball of thermalized bosons will be repulsively stabilized by even a tiny quartic coupling.

$$N_b = \frac{m_{pl}^2}{m_X^2}$$

A ball of fermions has fermi degeneracy pressure,  
but attractive self-interactions can counteract this.

$$N_f < \frac{m_{pl}^3}{m_X^3} \quad \text{if } \mathcal{L} \supset \alpha \phi \bar{\psi} \psi$$

attractive



A ball of thermalized bosons will be repulsively  
stabilized by even a tiny quartic coupling.

$$N_b > \frac{m_{pl}^2}{m_X^2}$$

A ball of fermions has fermi degeneracy pressure, but attractive self-interactions can counteract this.

$$N_f < \frac{m_{pl}^3}{m_X^3} \quad \text{if } \mathcal{L} \supset \alpha \phi \bar{\psi} \psi$$

attractive



A ball of thermalized bosons will be repulsively stabilized by even a tiny quartic coupling.

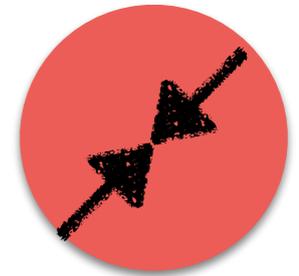
$$N_b > \frac{m_{pl}^2}{m_X^2} \quad \text{if } \mathcal{L} \supset \lambda \phi^4$$

and  $\lambda > 10^{-30}$

A ball of fermions has fermi degeneracy pressure, but attractive self-interactions can counteract this.

$$N_f < \frac{m_{\text{pl}}^3}{m_X^3} \quad \text{if } \mathcal{L} \supset \alpha \phi \bar{\psi} \psi$$

attractive



A ball of thermalized bosons will be repulsively stabilized by even a tiny quartic coupling.

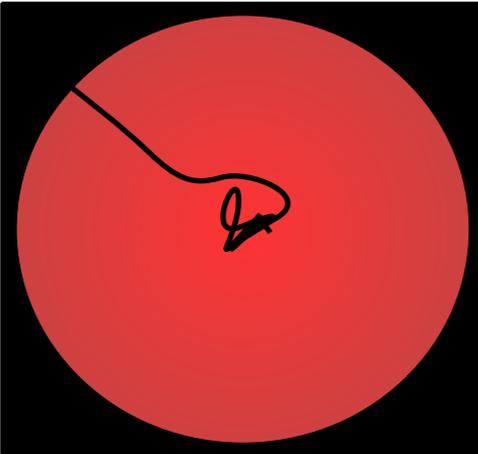
$$N_b > \frac{m_{\text{pl}}^2}{m_X^2} \quad \text{if } \mathcal{L} \supset \lambda \phi^4$$

and  $\lambda > 10^{-30}$

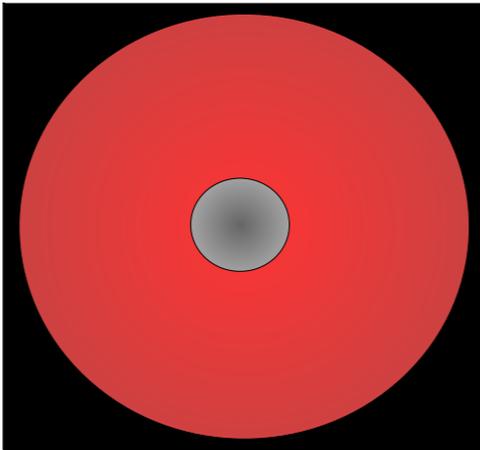
repulsive



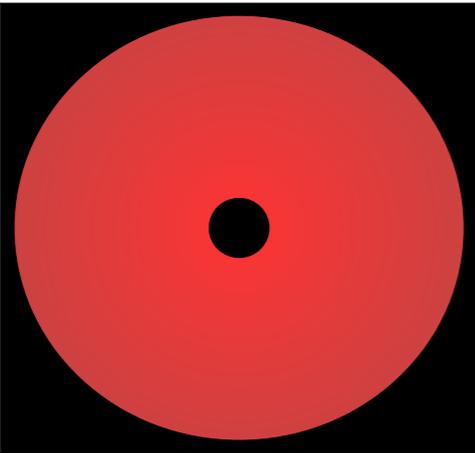
1] DM captured



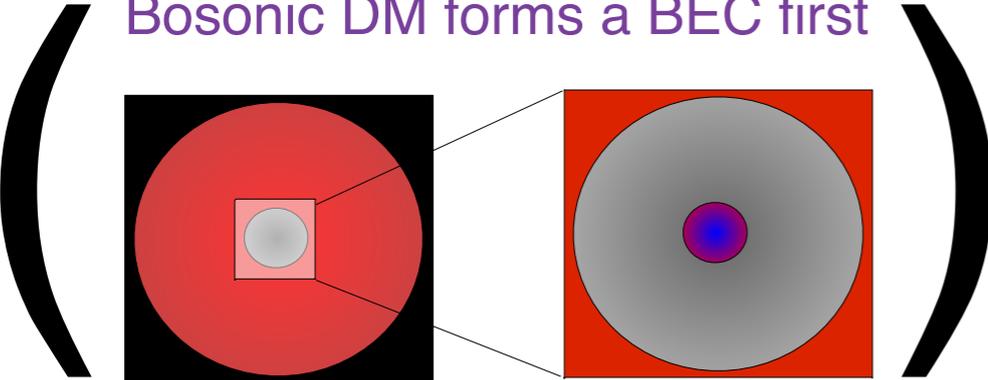
2] DM thermalizes



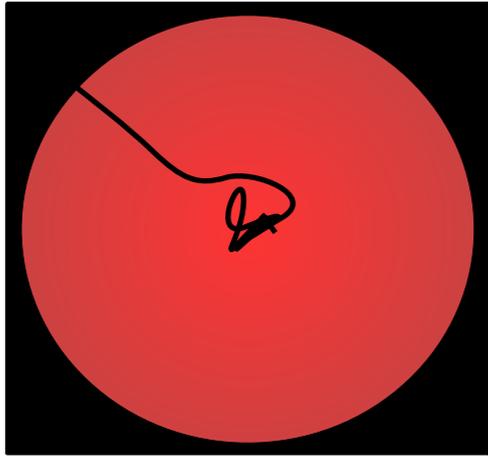
3] DM collapses



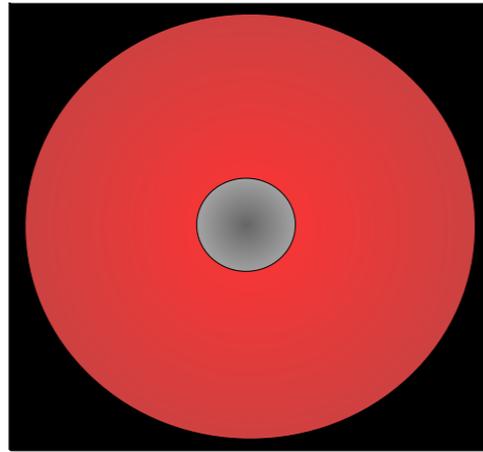
Bosonic DM forms a BEC first



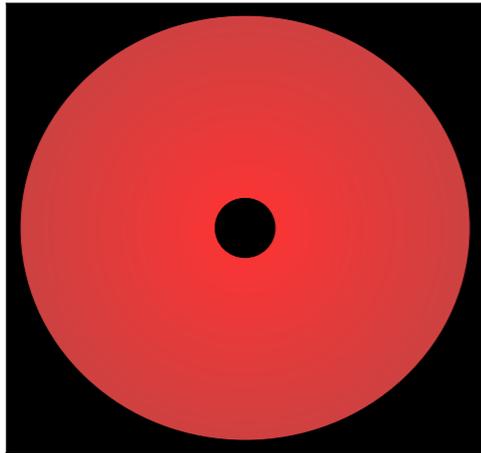
1] DM captured



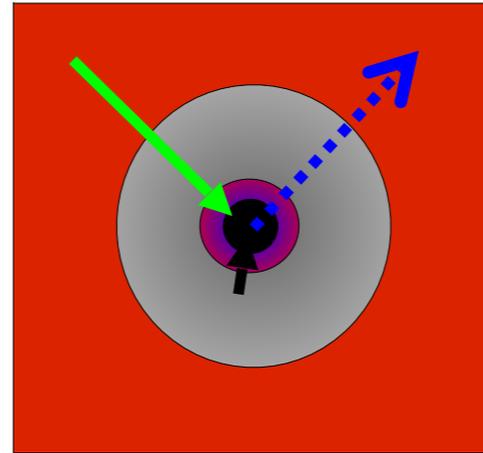
2] DM thermalizes



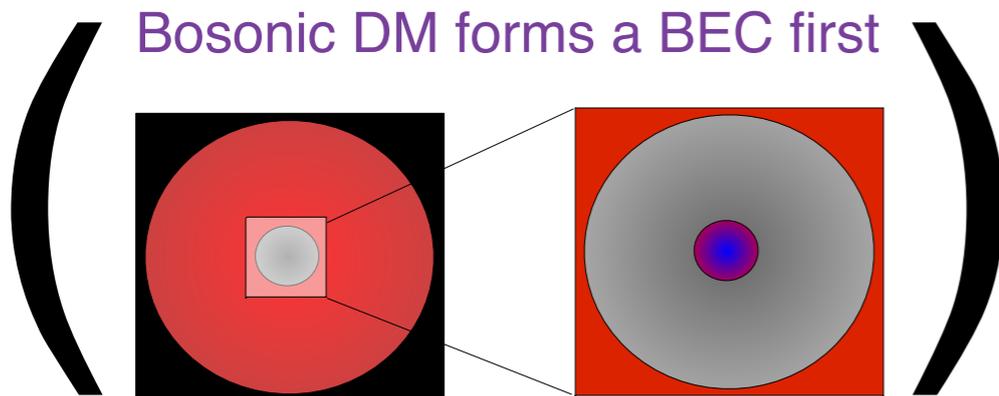
3] DM collapses



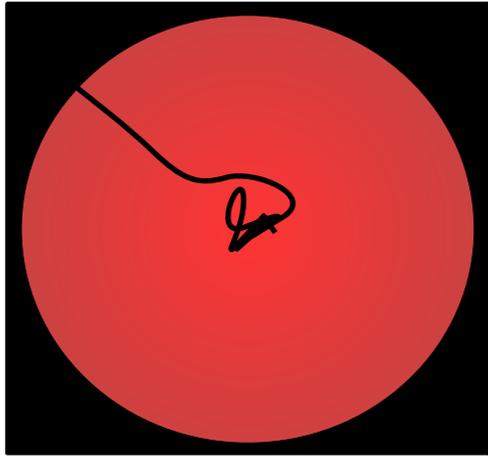
4] BH accretes, radiates



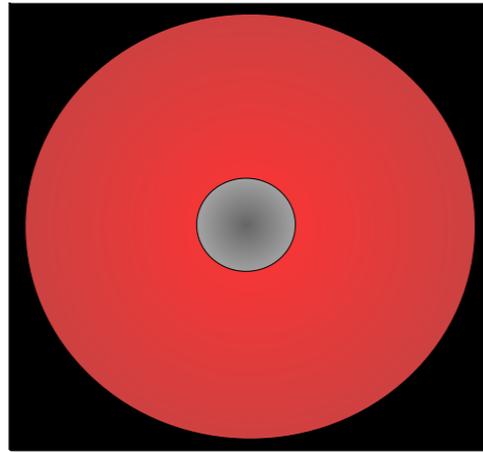
if it grows rapidly, then



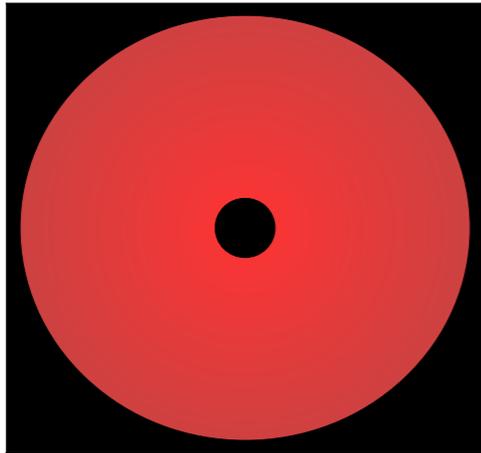
1] DM captured



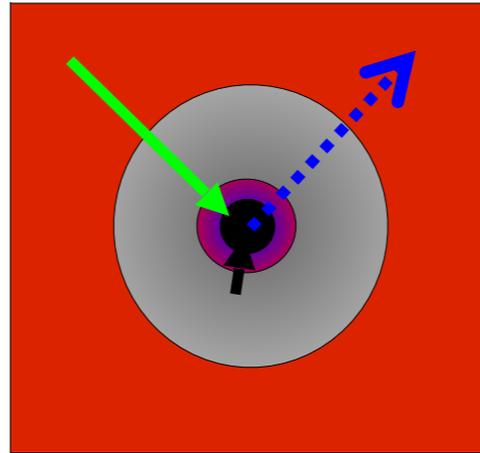
2] DM thermalizes



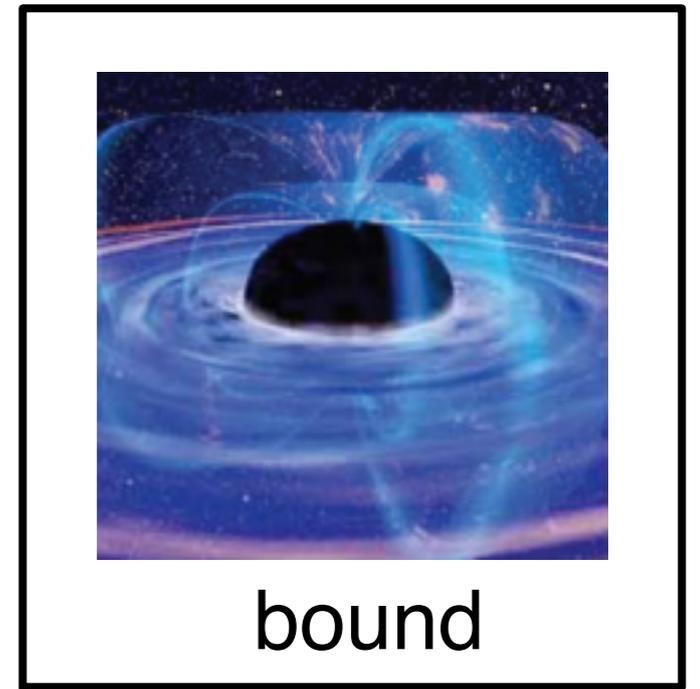
3] DM collapses



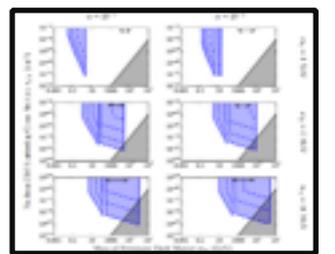
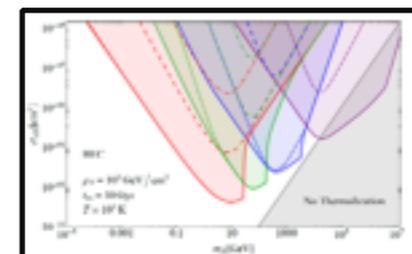
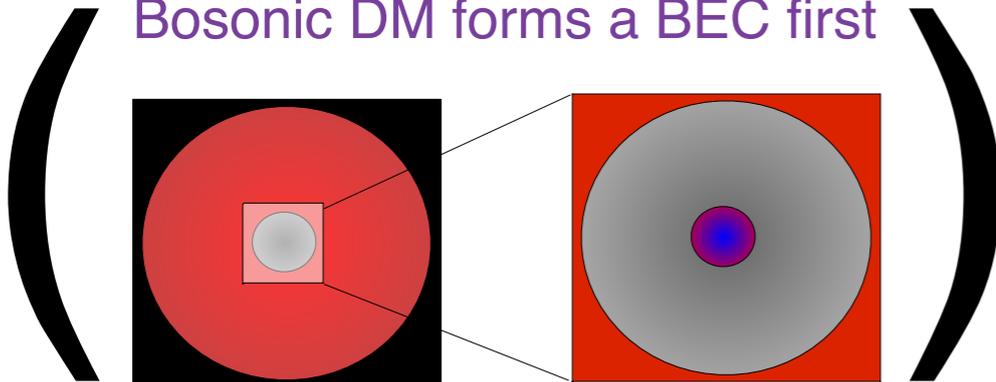
4] BH accretes, radiates



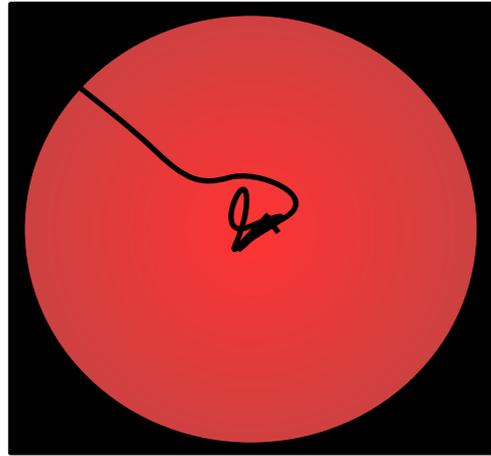
if it grows rapidly, then



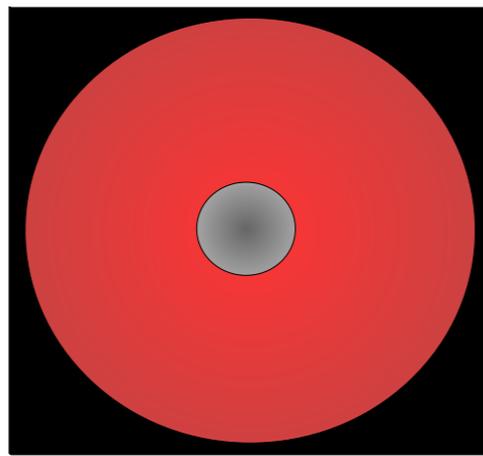
Bosonic DM forms a BEC first



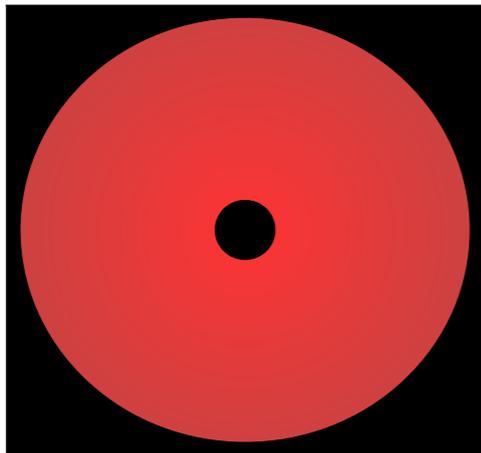
1] DM captured



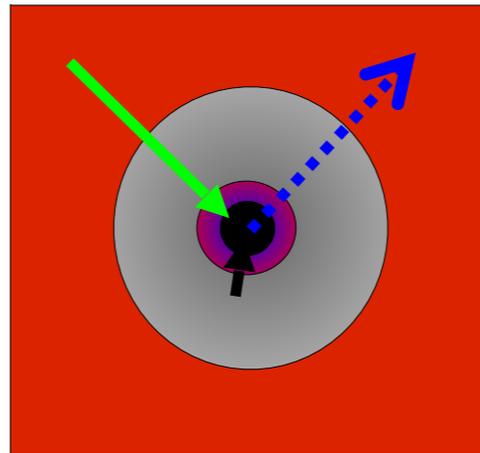
2] DM thermalizes



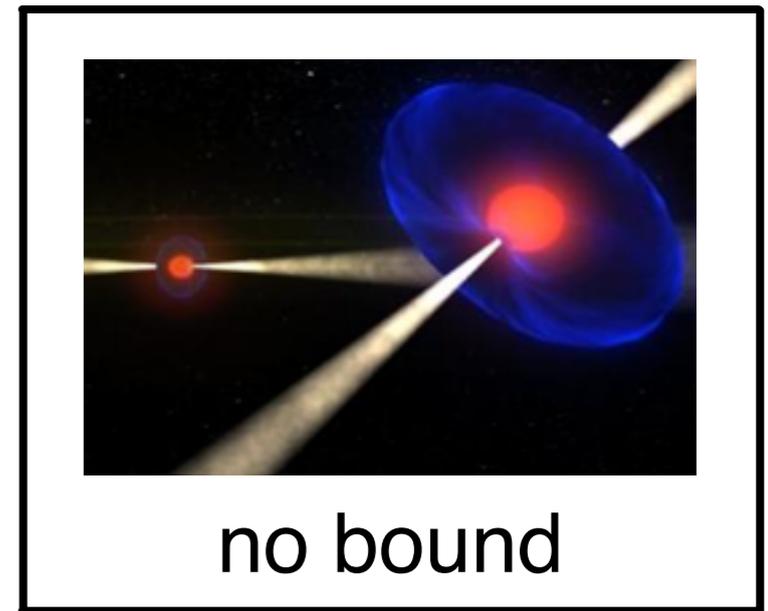
3] DM collapses



4] BH accretes, radiates

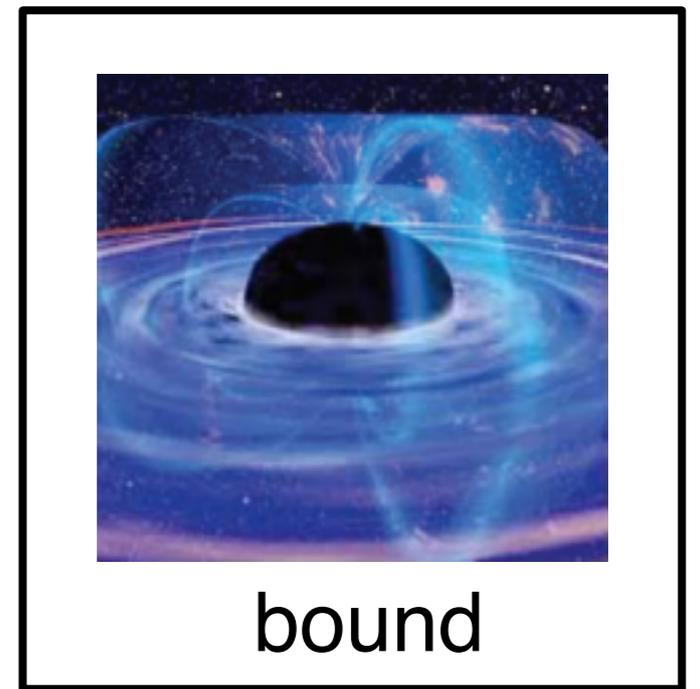


if it shrinks, (Hawking)



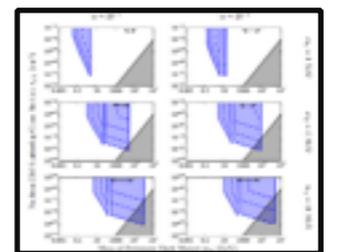
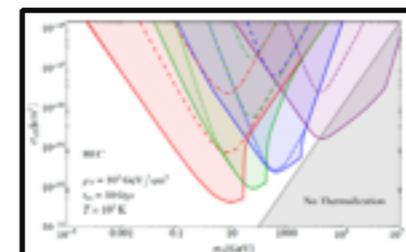
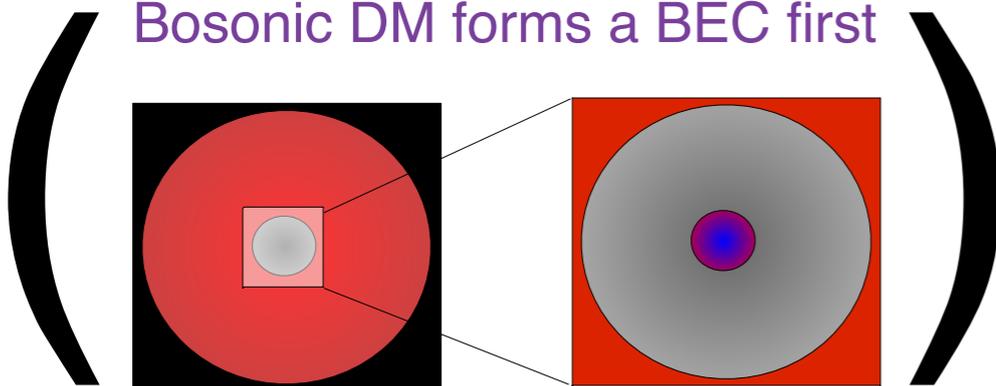
no bound

if it grows rapidly, then

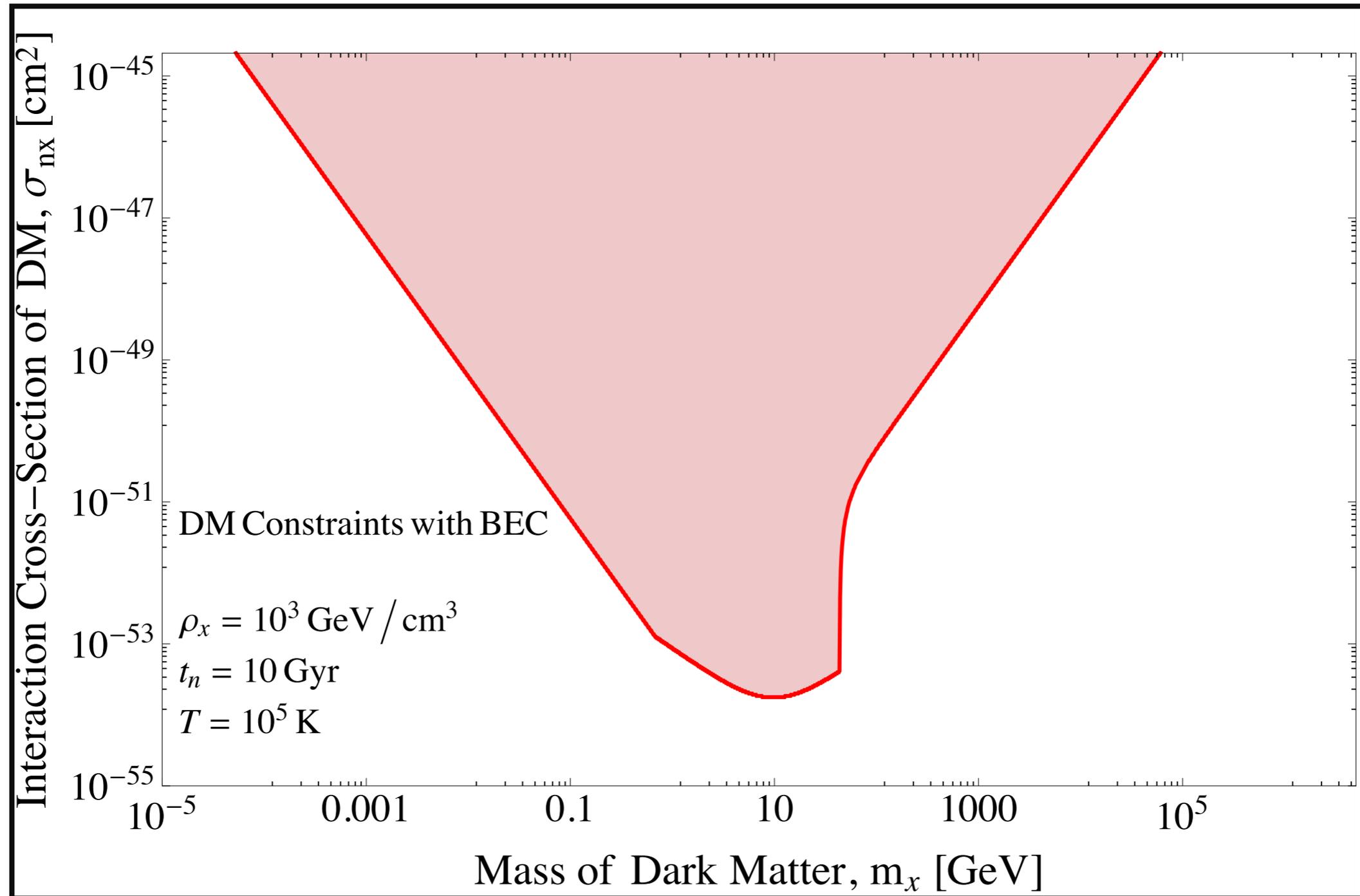


bound

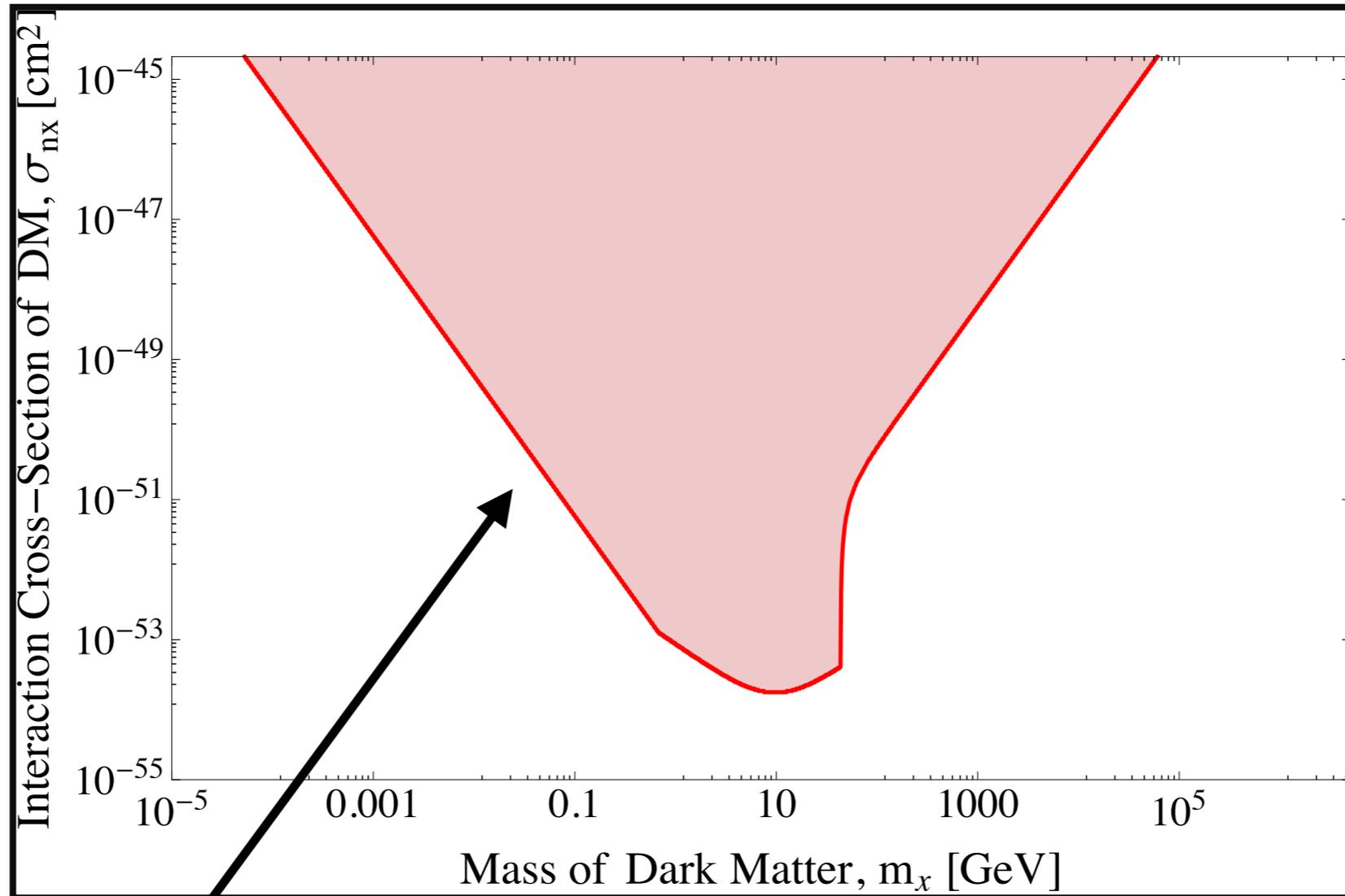
Bosonic DM forms a BEC first



# Asymmetric Boson DM Bounds

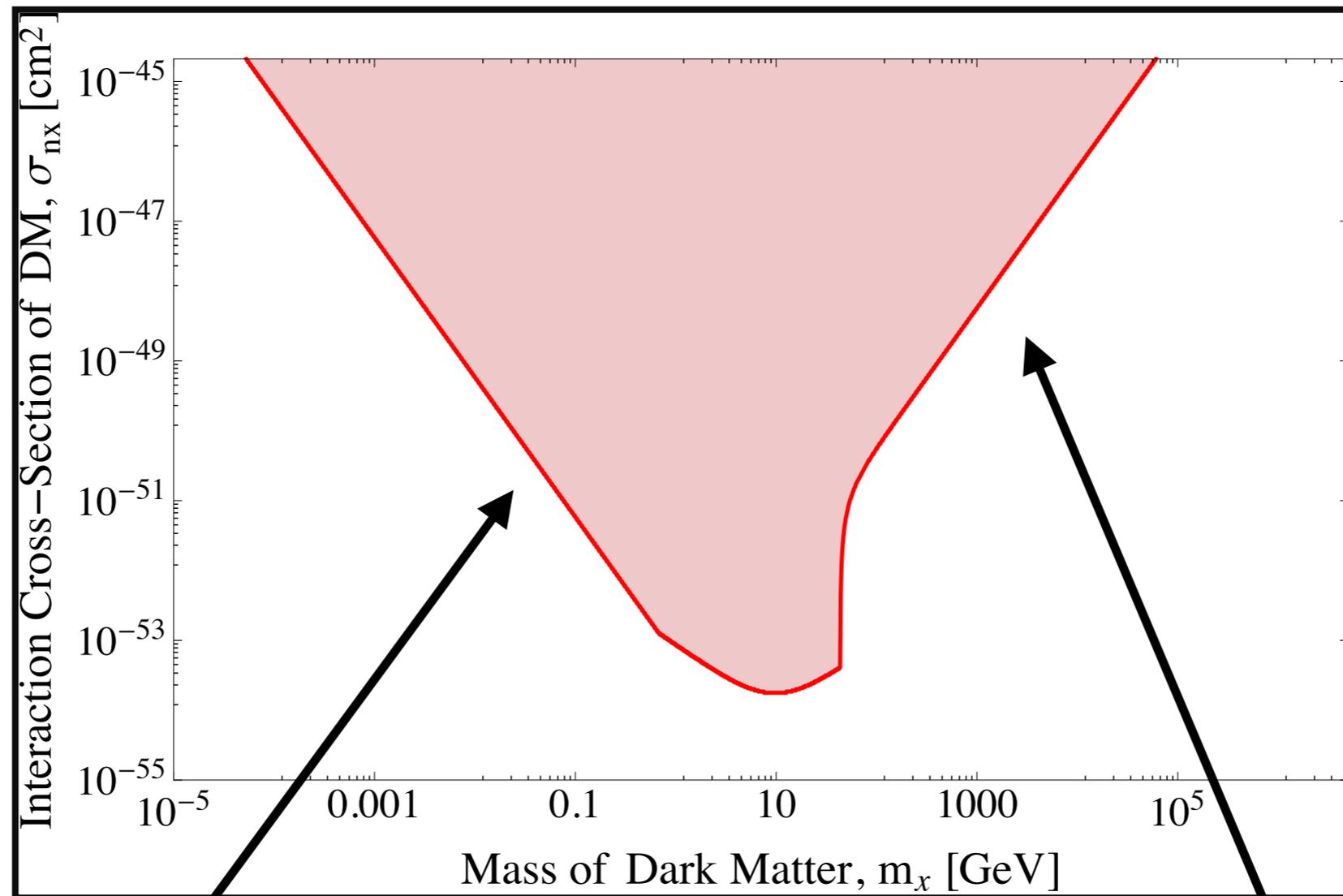


# Asymmetric Boson DM Bounds



Pauli Blocking:  $1/m_x$

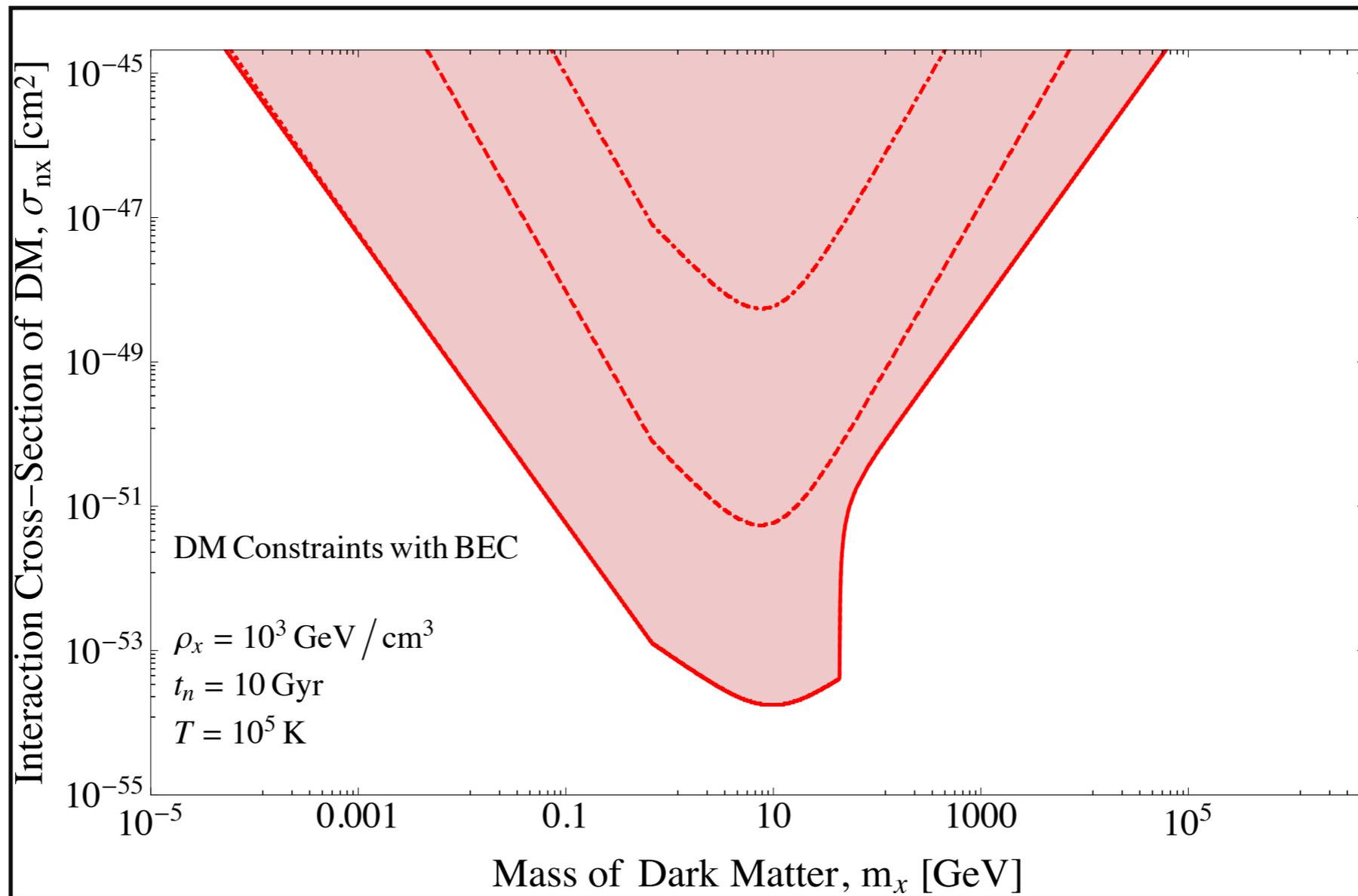
# Asymmetric Boson DM Bounds



Pauli Blocking:  $1/m_x$

Hawking vs Accretion  
of black hole

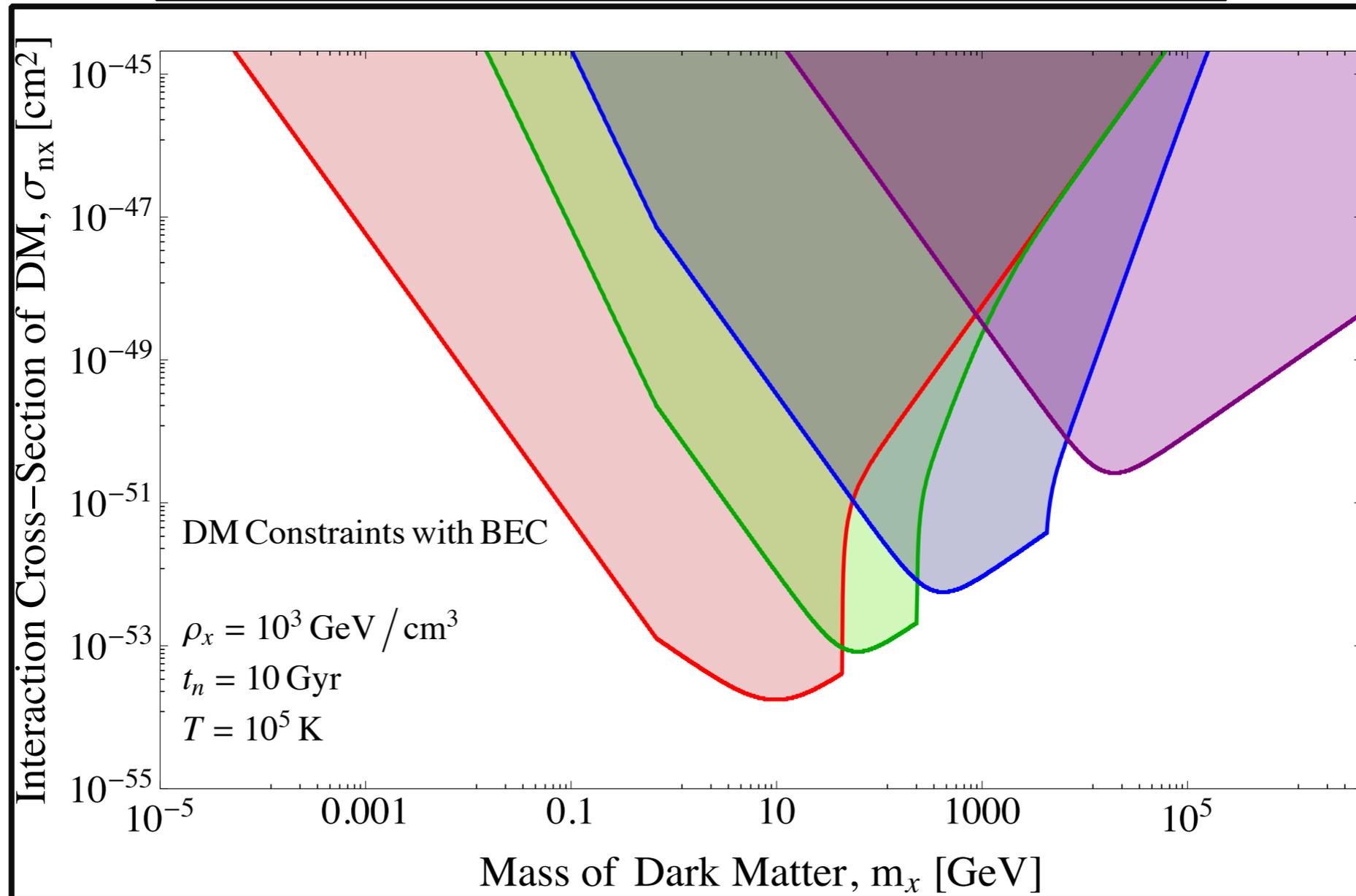
# Asymmetric Boson DM Bounds



Extremely tiny self-annihilations remove the bound.

$$\langle \sigma_a v \rangle = \boxed{0} \quad \boxed{10^{-50}} \quad \boxed{10^{-45}} \quad \boxed{10^{-42}} \text{ cm}^3/\text{s}$$

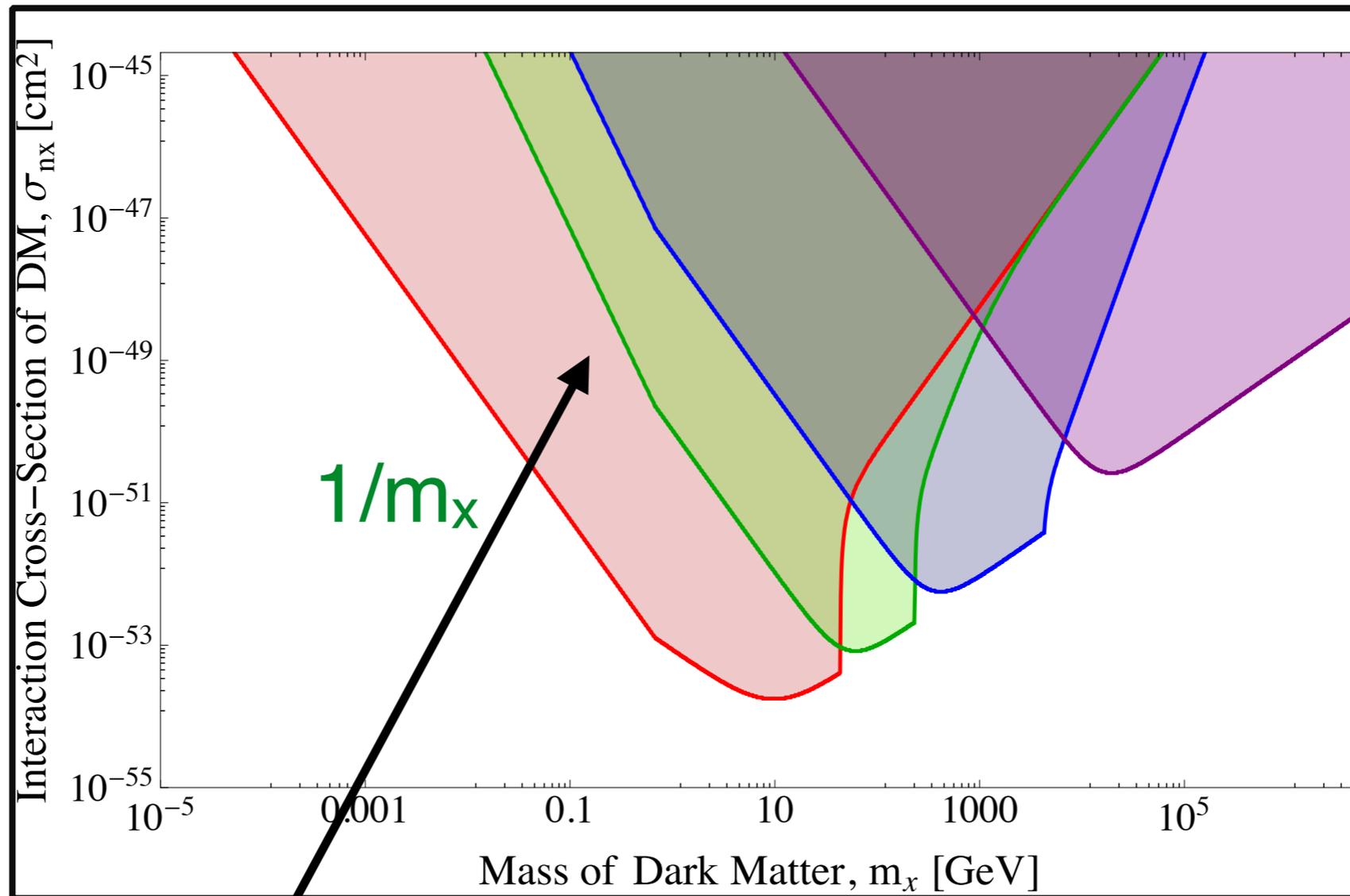
# Asymmetric Boson DM Bounds



Repulsive self-interactions push the bounds to higher masses — **lower mass DM can no longer collapse**, **higher mass DM collapses to a bigger black hole**.

$$\lambda = \boxed{0} \quad \boxed{10^{-30}} \quad \boxed{10^{-25}} \quad \boxed{10^{-15}}$$

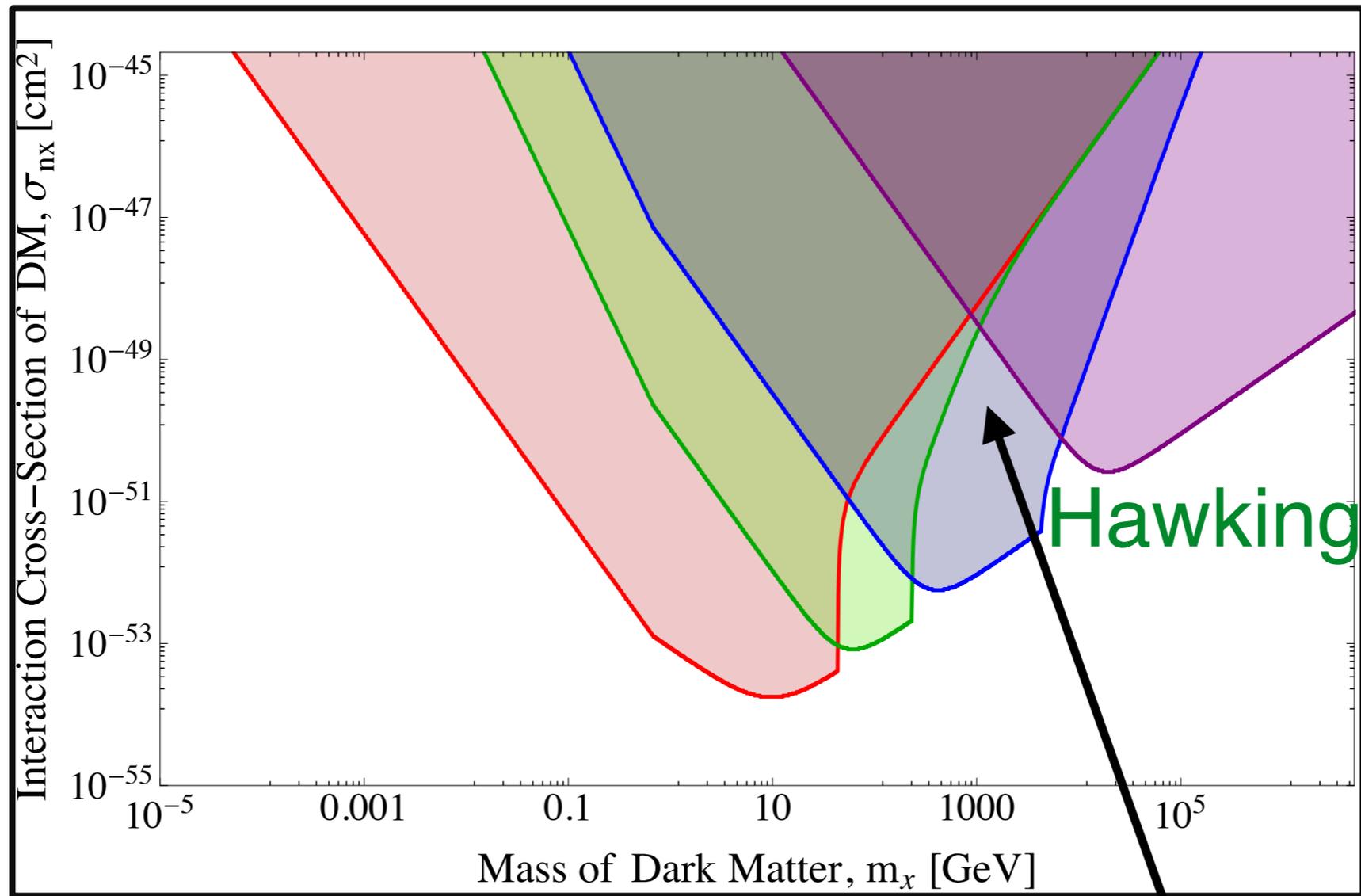
# Asymmetric Boson DM Bounds



$$\lambda = \boxed{0} \quad \boxed{10^{-30}} \quad \boxed{10^{-25}} \quad \boxed{10^{-15}}$$

Need more massive ball of DM to collapse

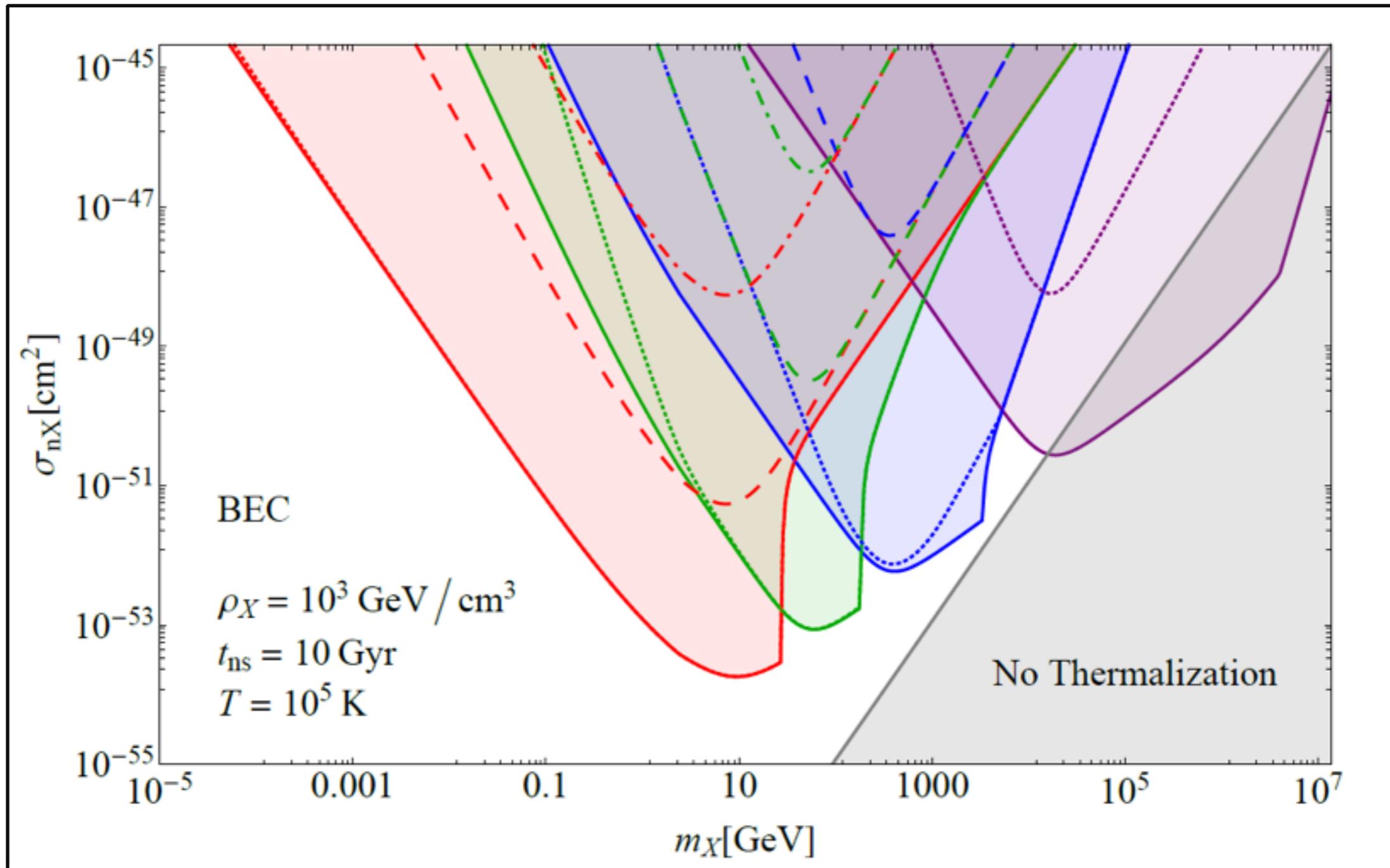
# Asymmetric Boson DM Bounds



$$\lambda = \boxed{0} \quad \boxed{10^{-30}} \quad \boxed{10^{-25}} \quad \boxed{10^{-15}}$$

Need more massive ball of DM to collapse

# Asymmetric Boson DM Bounds



$$\lambda = \boxed{0} \quad \boxed{10^{-30}} \quad \boxed{10^{-25}} \quad \boxed{10^{-15}} \quad [\sigma_{\text{XX}} \sim 0, 10^{-118}, 10^{-98}, 10^{-58} \text{ cm}^2]$$

$$\langle \sigma_a v \rangle = \boxed{0} \quad \boxed{10^{-50}} \quad \boxed{10^{-45}} \quad \boxed{10^{-42}} \text{ cm}^3/\text{s}$$

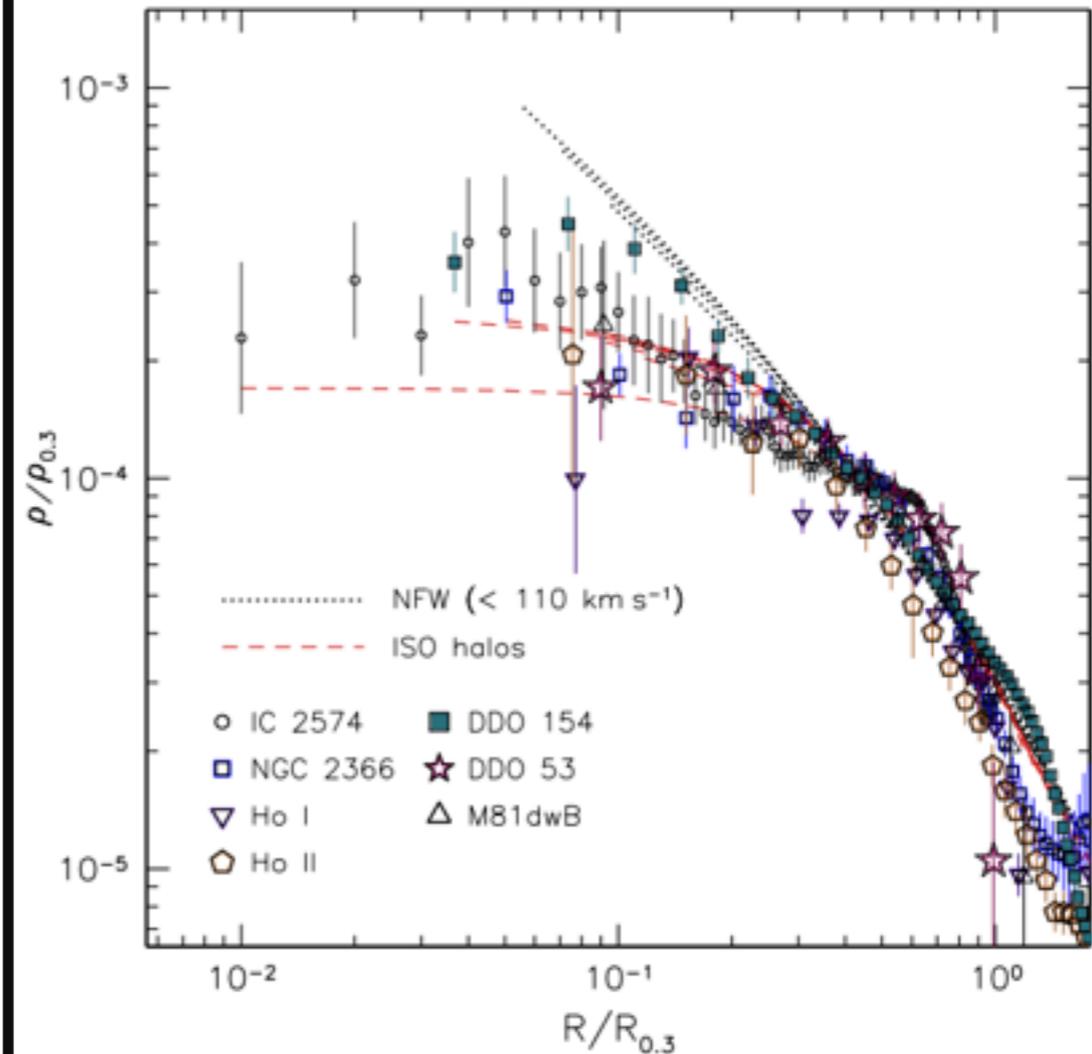
Fermionic asymmetric dark matter

# Seven dwarves with mined out DM cores

7 dwarf galaxies measured by THINGS do not show a cold, collisionless NFW profile which would **cusp** in the center. They instead have a more cored out shape.

-Caveat: **Baryonic outflow** via SN

-Counter: less luminous galaxies should **not experience outflow**, but seem to.



Heon et al.  
1011.0899

# Seven dwarves with mined out DM cores

7 dwarf galaxies measured by THINGS do not show a cold, collisionless NFW profile which would **cusp** in the center. They instead have a more cored out shape.

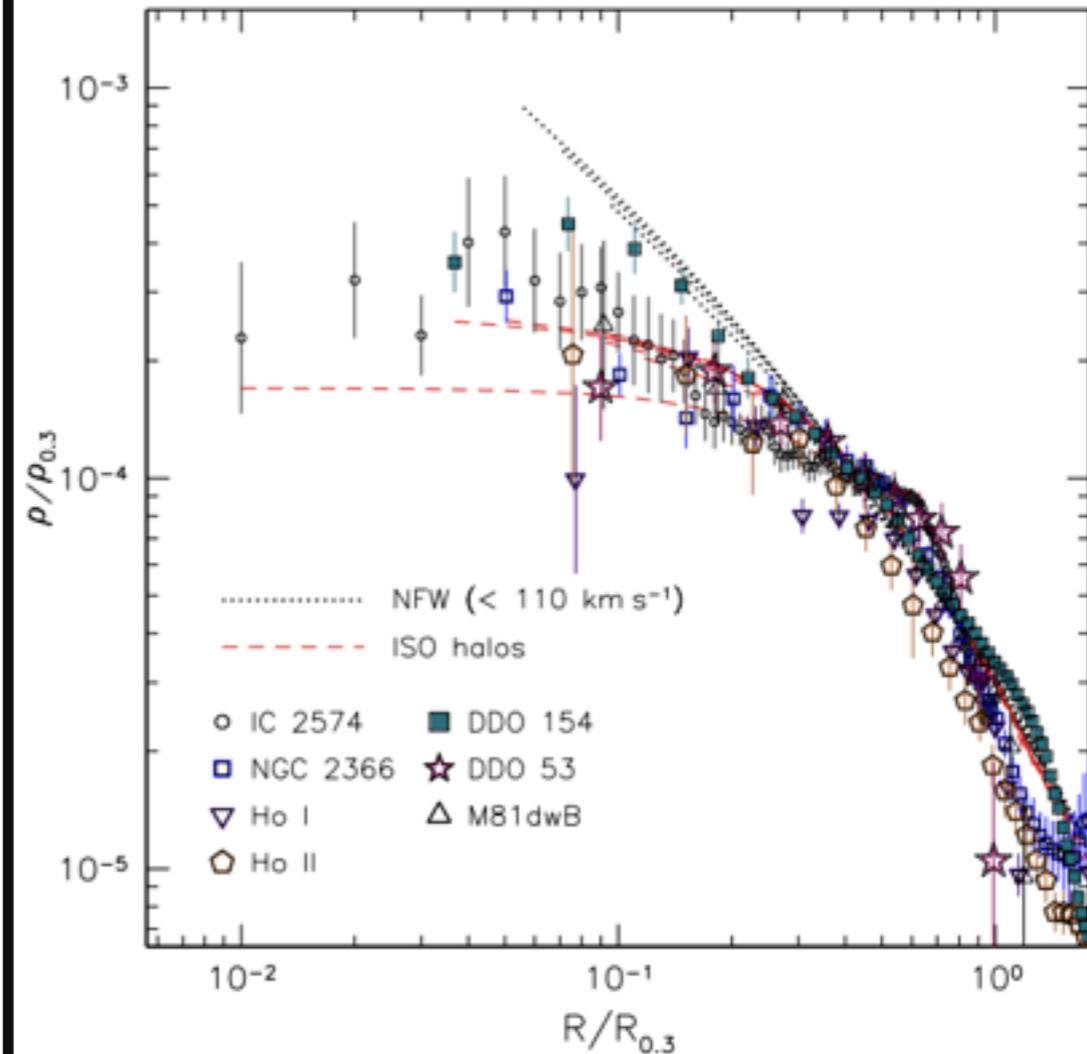
-Caveat: **Baryonic outflow** via SN

-Counter: less luminous galaxies should **not experience outflow**, but seem to.

Also, many simulations suggest that we should have  $\sim 50$  subhalos in the MW, we see only 12.

-Caveat: Different models of star formation, **subhalos too dim?**

-Counter: **“Too big to fail to form star subhalos not seen in the Milky Way.**



Heon et al.  
1011.0899

# Resonant Self-Interacting Fermionic DM

Bullet clusters (1000 km/s) and spiral galaxies (200 km/s) constrain the cross-section of dark matter with itself to

$$\sigma/m < 1 \text{ cm}^2/\text{g}, \quad \sigma/m < 1 \text{ cm}^2/\text{g}$$

# Resonant Self-Interacting Fermionic DM

Bullet clusters (1000 km/s) and spiral galaxies (200 km/s) constrain the cross-section of dark matter with itself to

$$\sigma/m < 1 \text{ cm}^2/\text{g}, \quad \sigma/m < 1 \text{ cm}^2/\text{g}$$

But the preferred cross-section to core the dwarf halo (10 km/s) is

$$\sigma/m \sim .1\text{-}10 \text{ cm}^2/\text{g}$$

Answer: velocity dependent cross-section provided by light mediator.

$$\text{if } \mathcal{L} \supset \alpha \phi \bar{\psi} \psi$$

# Resonant Self-Interacting Fermionic DM

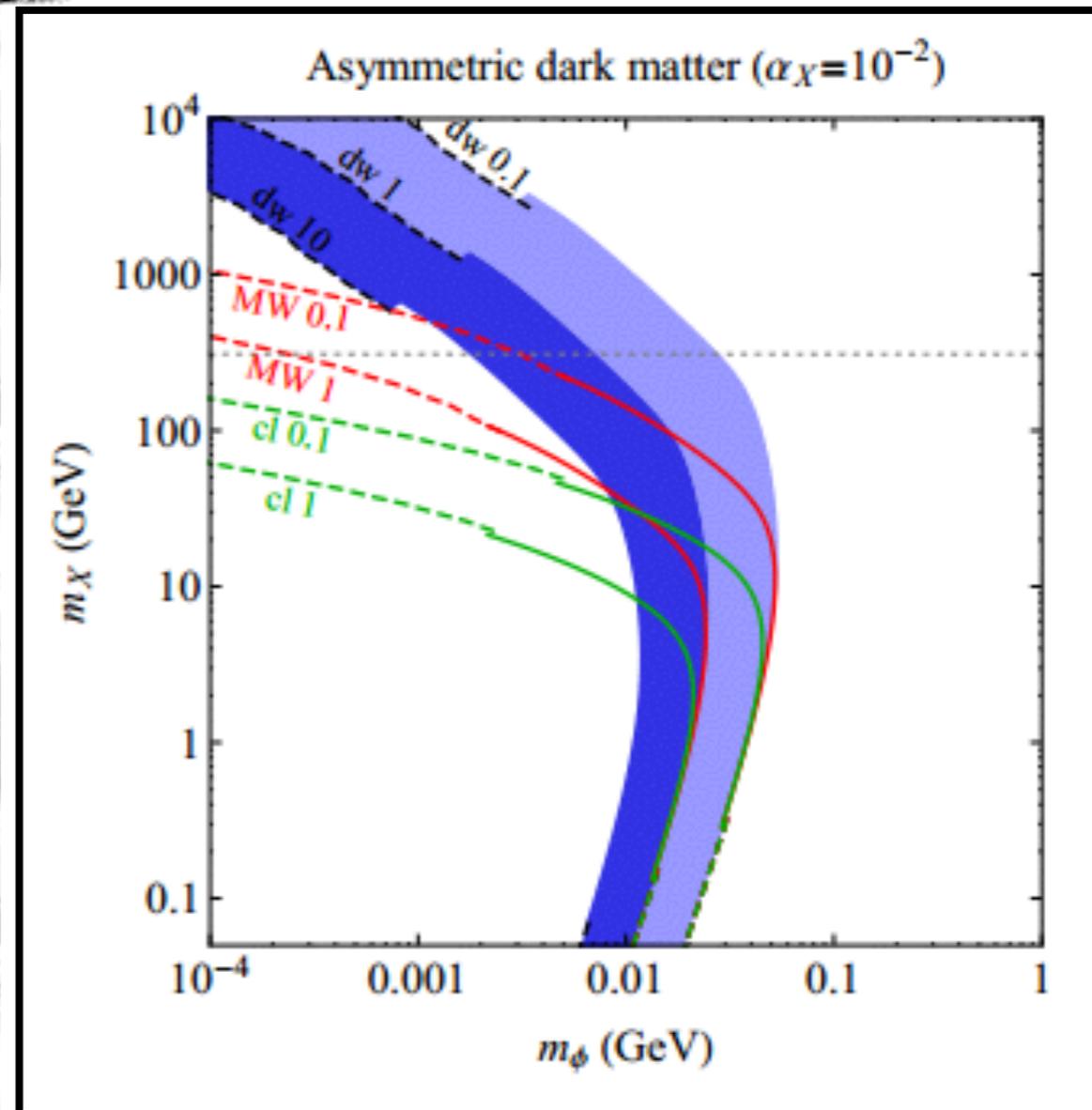
Bullet clusters (1000 km/s) and spiral galaxies (200 km/s) constrain the cross-section of dark matter with itself to

$$\sigma/m < 1 \text{ cm}^2/\text{g}, \quad \sigma/m < 1 \text{ cm}^2/\text{g}$$

But the preferred cross-section to core the dwarf halo (10 km/s) is

$$\sigma/m \sim .1\text{-}10 \text{ cm}^2/\text{g}$$

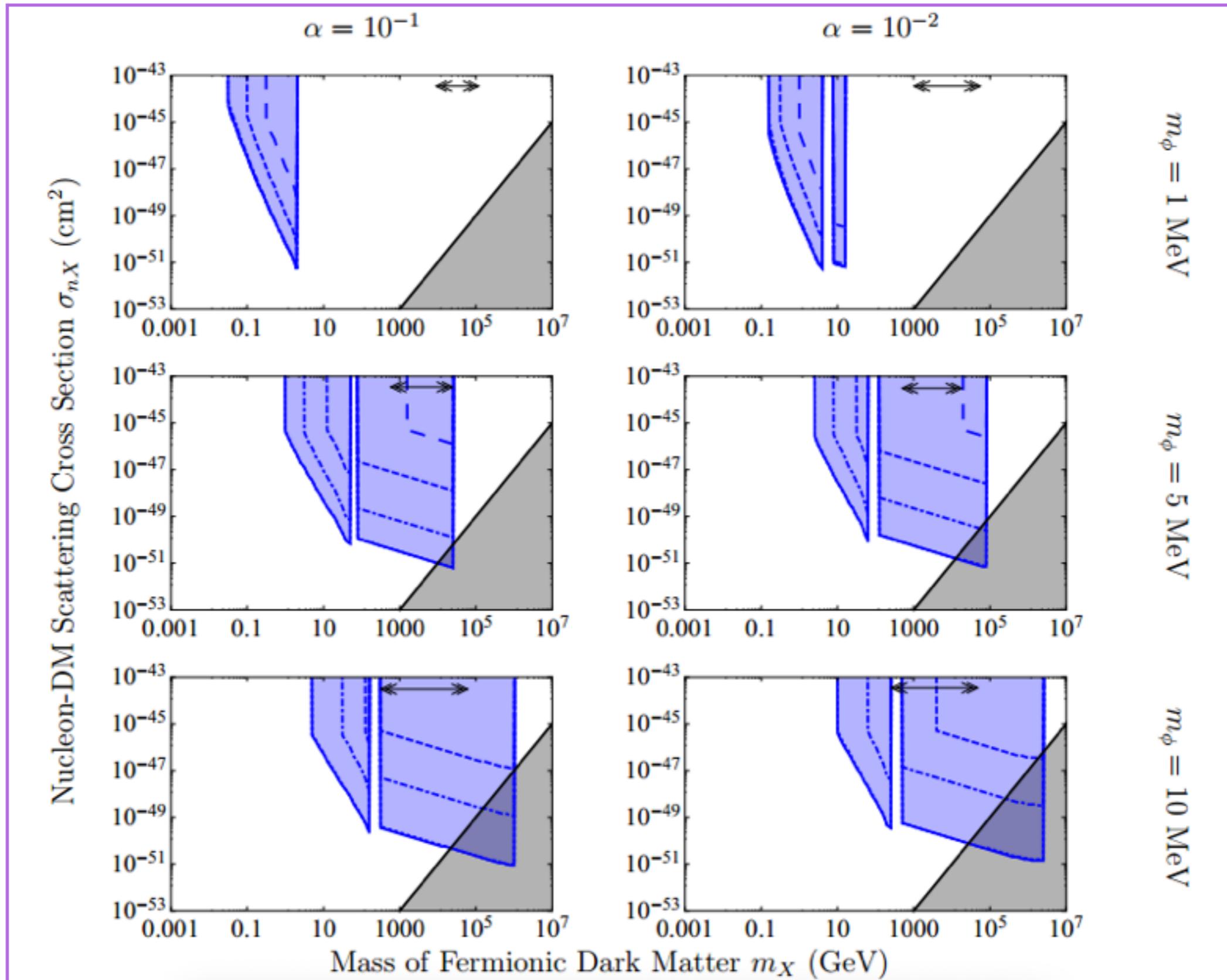
Answer: velocity dependent cross-section provided by light mediator.



Tulin, Yu, Zurek  
1210.0900

$$\text{if } \mathcal{L} \supset \alpha \phi \bar{\psi} \psi$$

# Bounds on SIDM Fermions



$$\langle \sigma_a v \rangle = \boxed{0} \boxed{10^{-47}} \boxed{10^{-45}} \boxed{10^{-43}} \text{ cm}^3/\text{s}$$

JB, Kumar, Fukushima, Stopnitzky (2013)

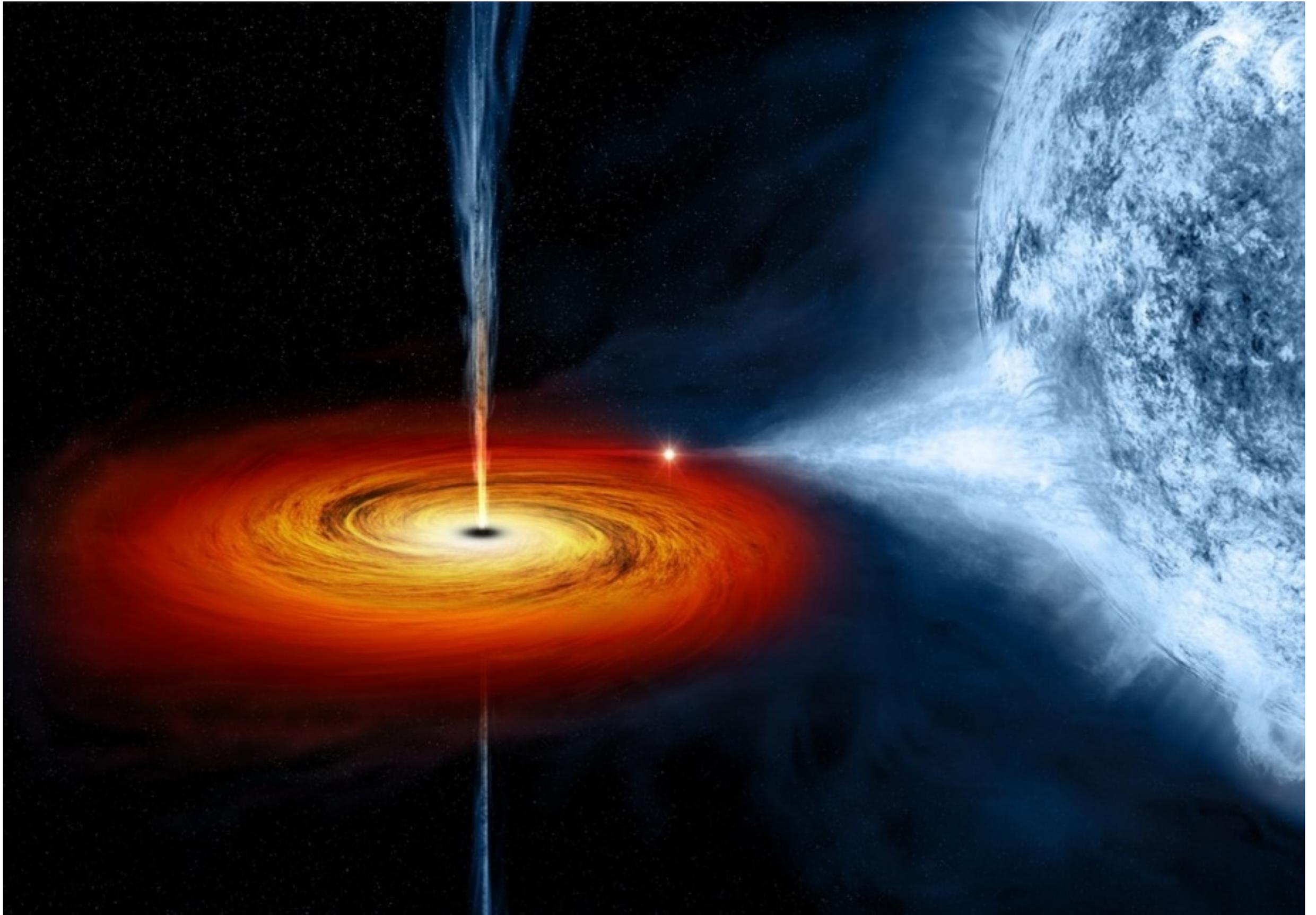
# The missing pulsar problem

Prior to the detection of a very luminous magnetar one tenth of a parsec away from the ( $10^8 \odot$  mass) black hole at the galactic center, it was assumed that pulsars had not been detected there, because a charged screen of material at the galactic center was broadening pulse signals. Pulsars are expected because of a large population of high mass progenitor stars and X-ray binary systems.

# The missing pulsar problem

However, measurements of radio pulses from the newly discovered galactic center magnetar indicate a much cleaner path for radio pulses than was supposed. In addition, measurements of the radio pulses' angular broadening match those of SgA\*. This is evidence that the scattering screen is homogeneous.

(Older) millisecond pulsars form from x-ray binaries

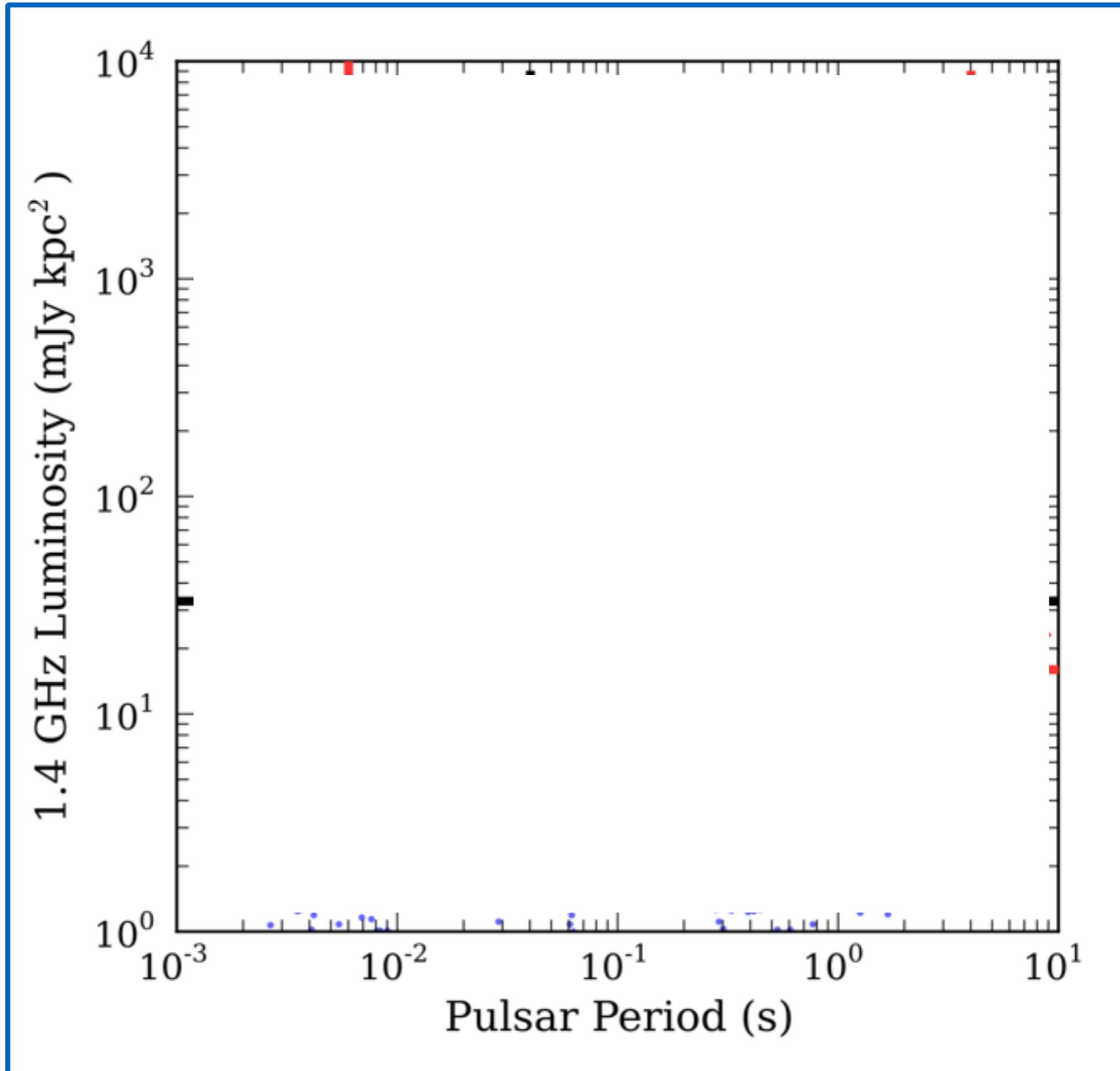


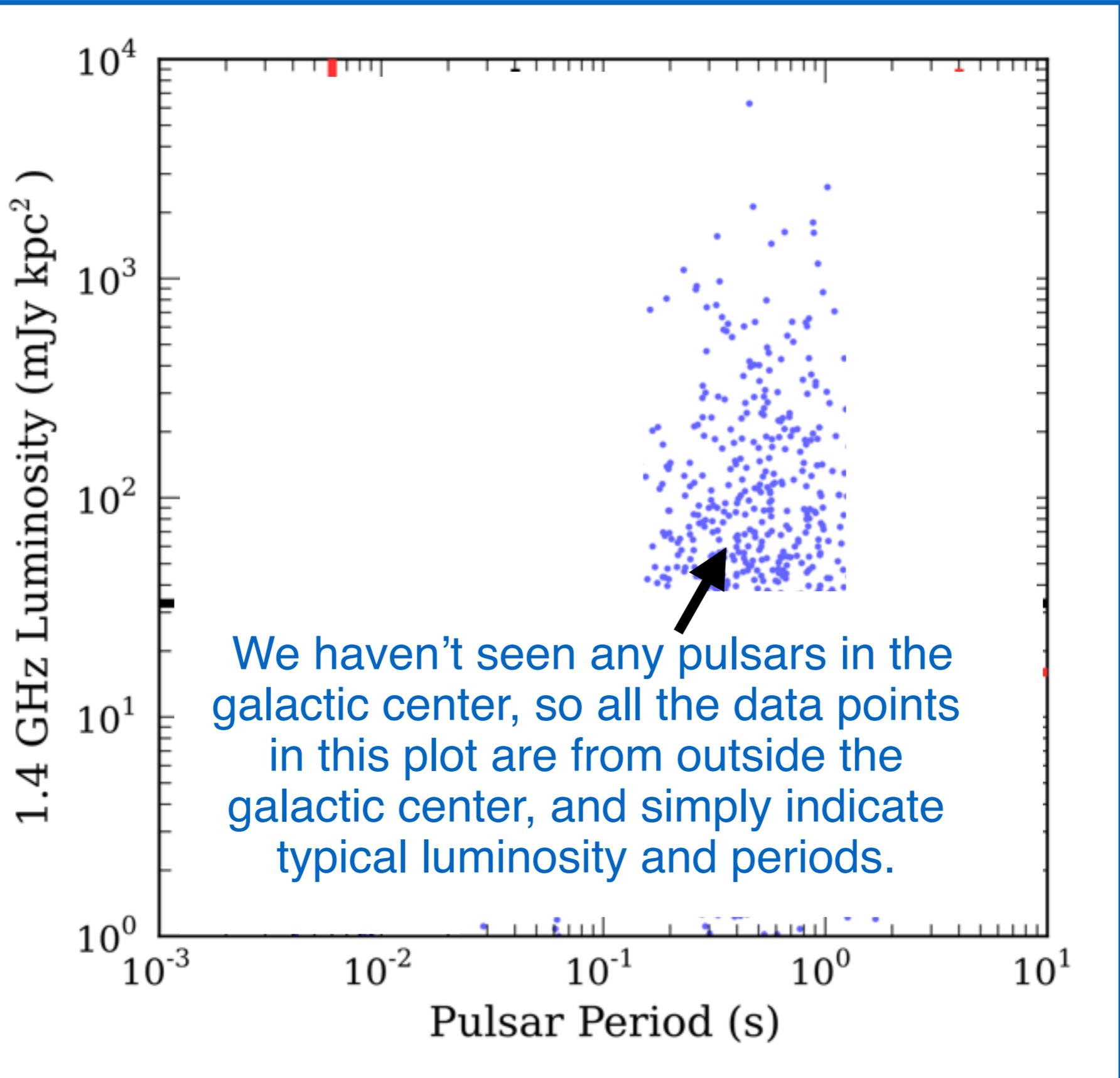
# The missing pulsar problem

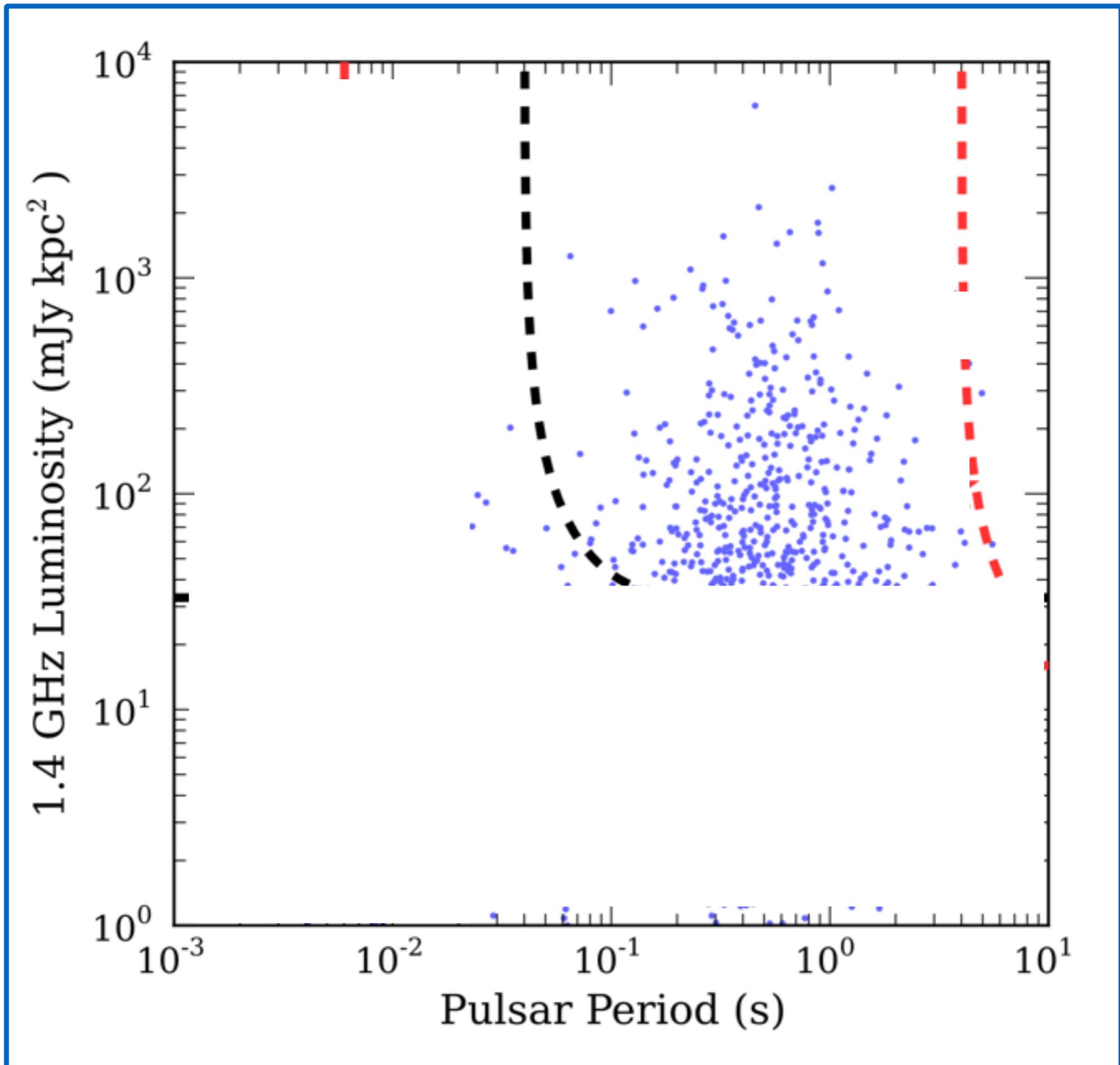
However, measurements of radio pulses from the newly discovered galactic center magnetar indicate a much cleaner path for radio pulses than was supposed. In addition, measurements of the radio pulses' angular broadening match those of SgA\*. This is evidence that the scattering screen is homogeneous.

This creates two missing pulsar problems. Both young and old millisecond pulsars seem to be absent. There are 50-500 missing millisecond pulsars.

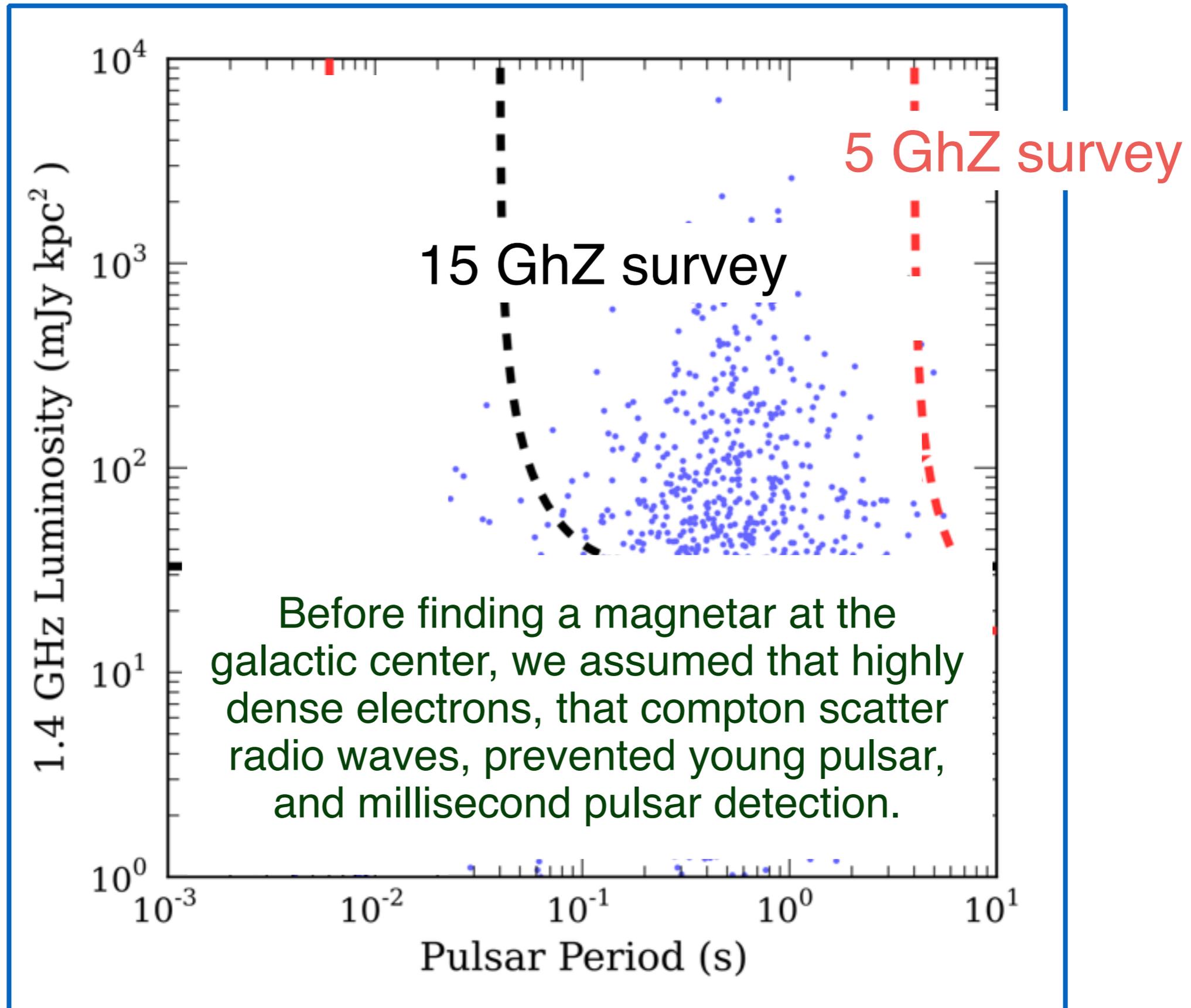
# The missing pulsar problem

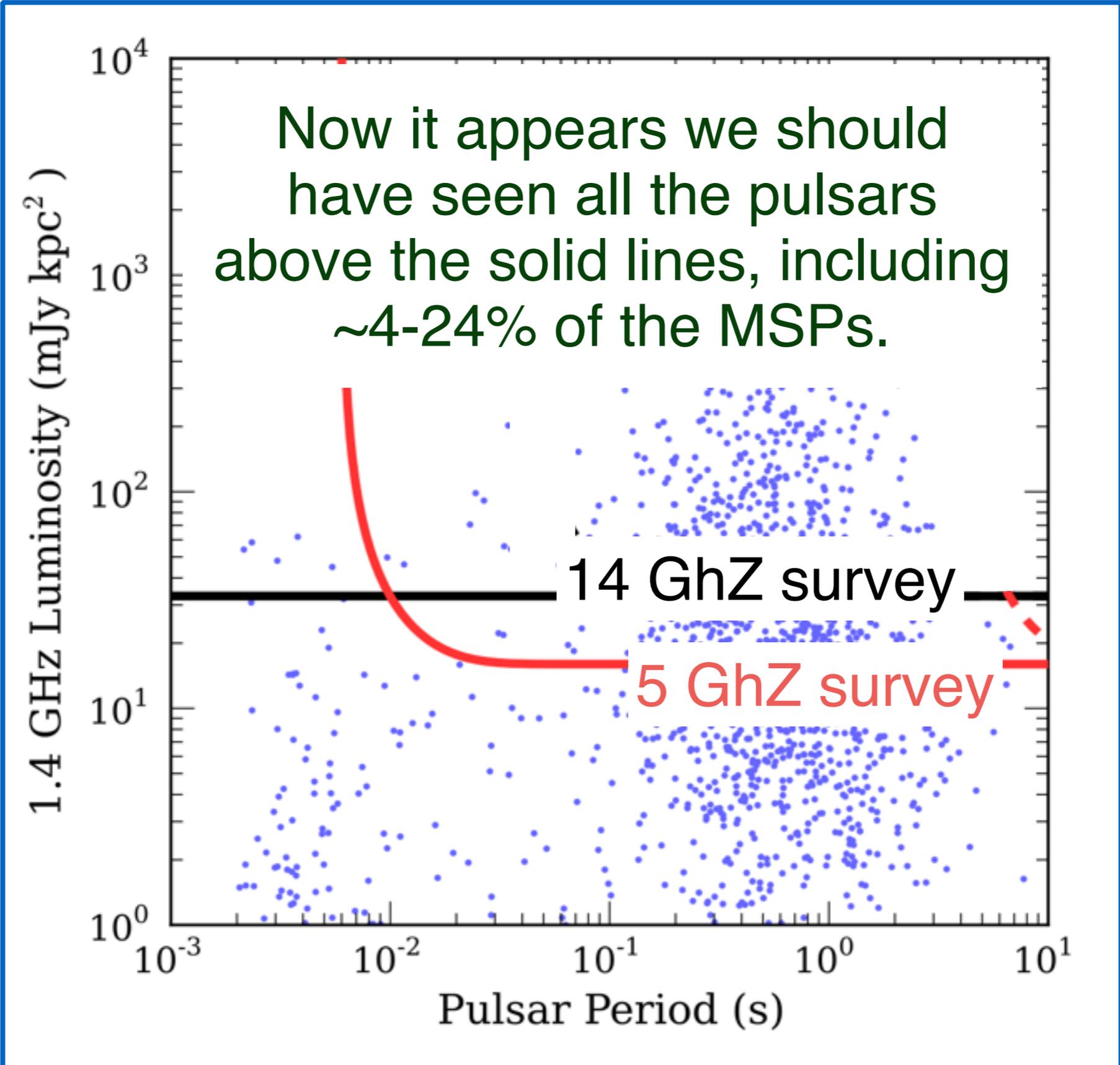




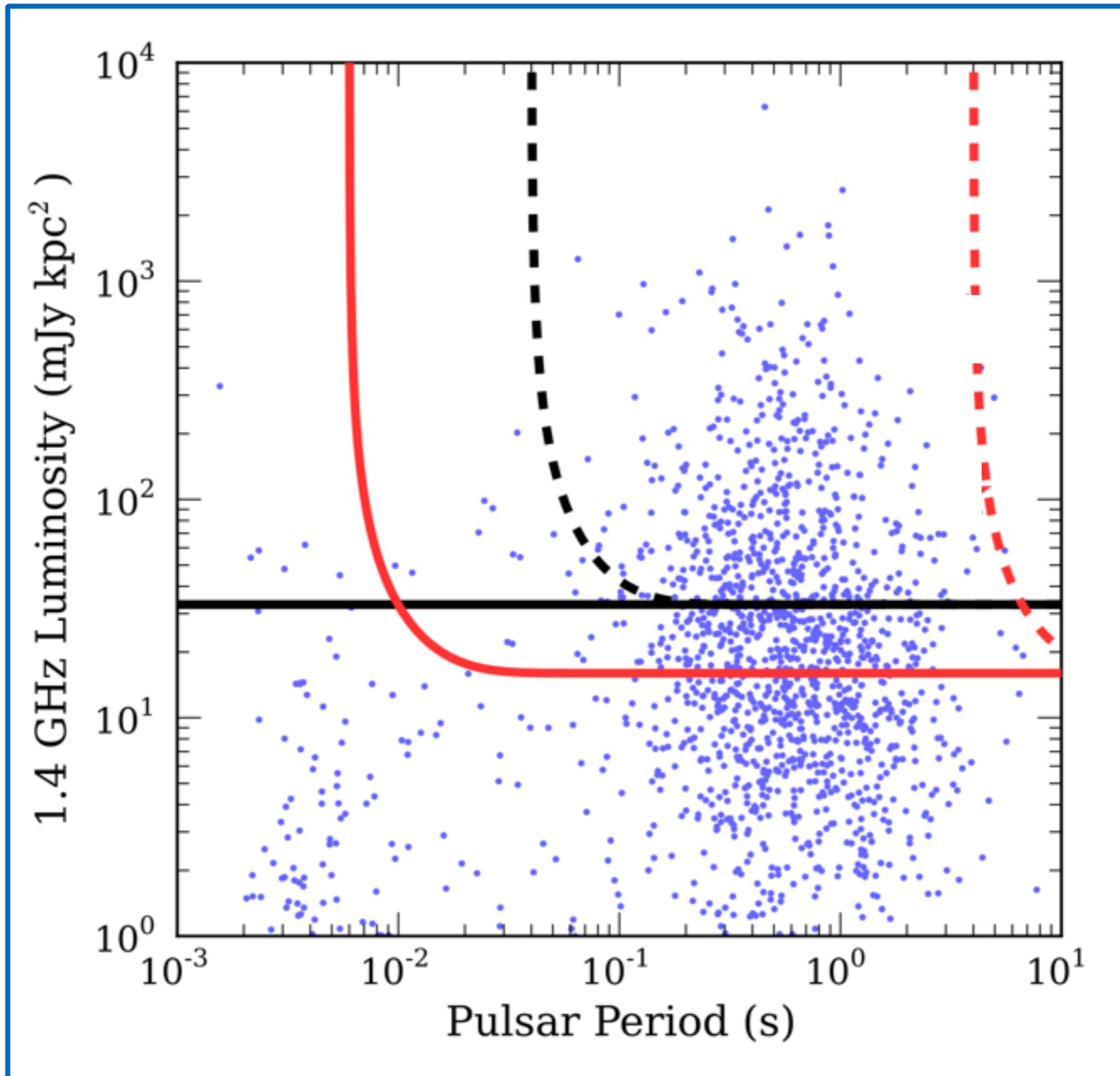


# Survey sensitivity before magnetar





# So where are the galactic center pulsars?



# Asymmetric Dark Matter Imploding Pulsars

The capture rate for dark matter on pulsars scales inversely with velocity dispersion and linearly with the local dark matter density.

# Asymmetric Dark Matter Imploding Pulsars

The capture rate for dark matter on pulsars scales inversely with velocity dispersion and linearly with the local dark matter density.

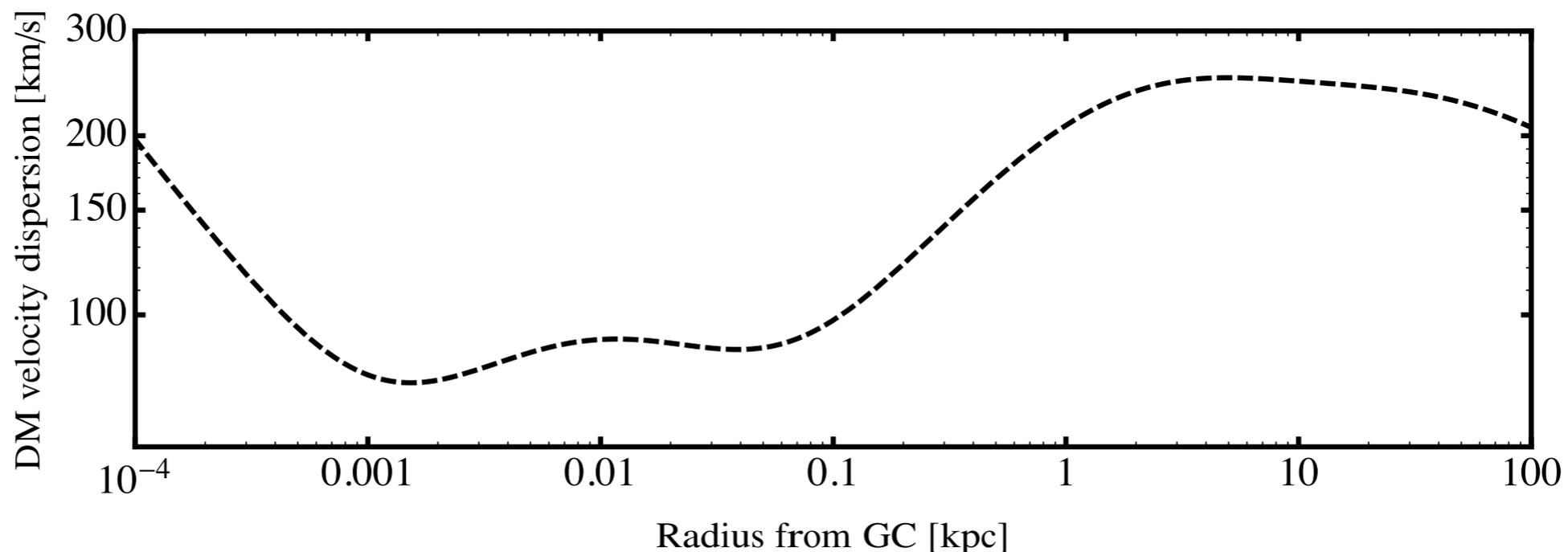
$$C_X \propto \frac{\rho_X}{\bar{v}} \sigma_{nX}$$

# Asymmetric Dark Matter Imploding Pulsars

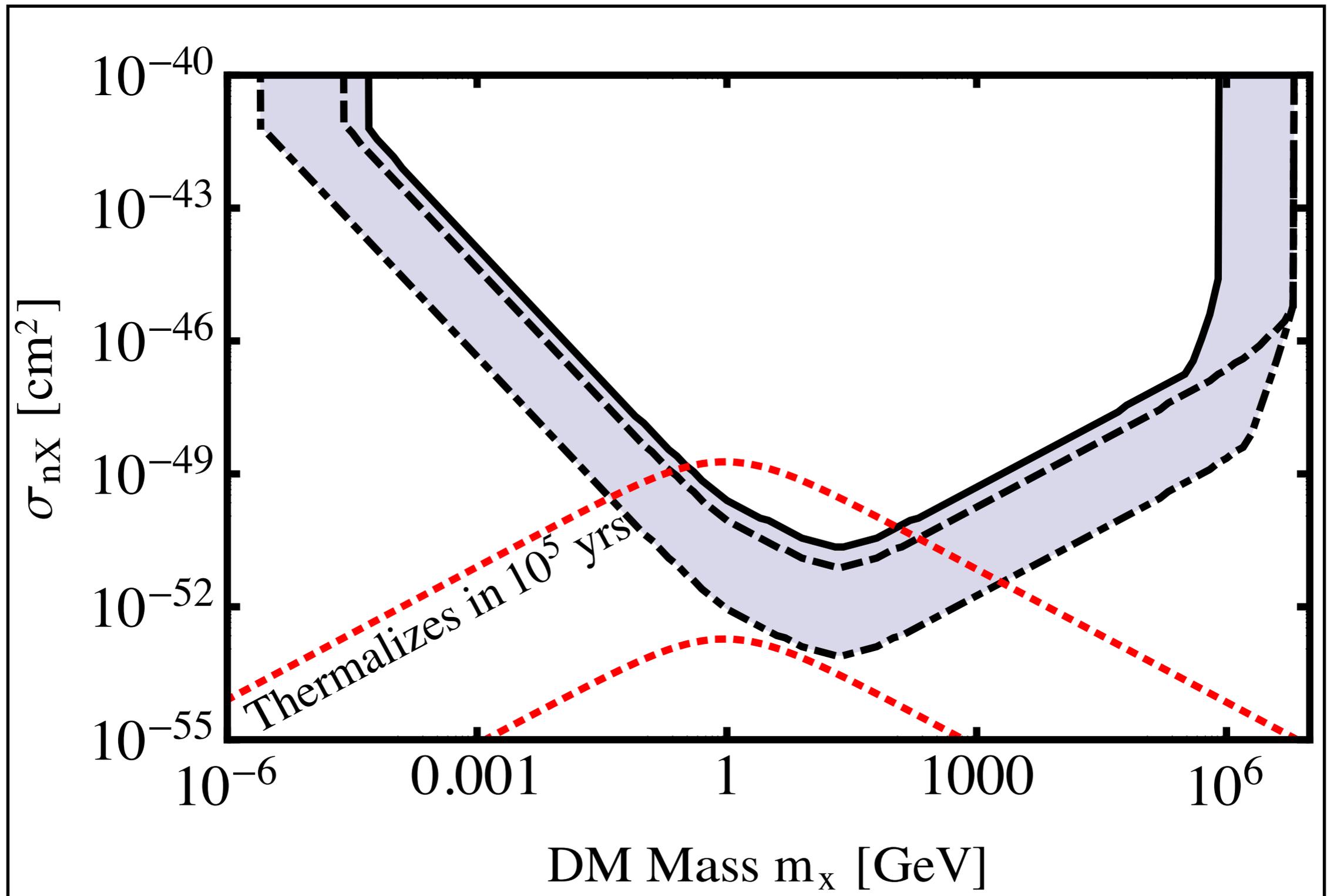
The capture rate for dark matter on pulsars scales inversely with velocity dispersion and linearly with the local dark matter density.

$$C_X \propto \frac{\rho_X}{\bar{v}} \sigma_{nX}$$

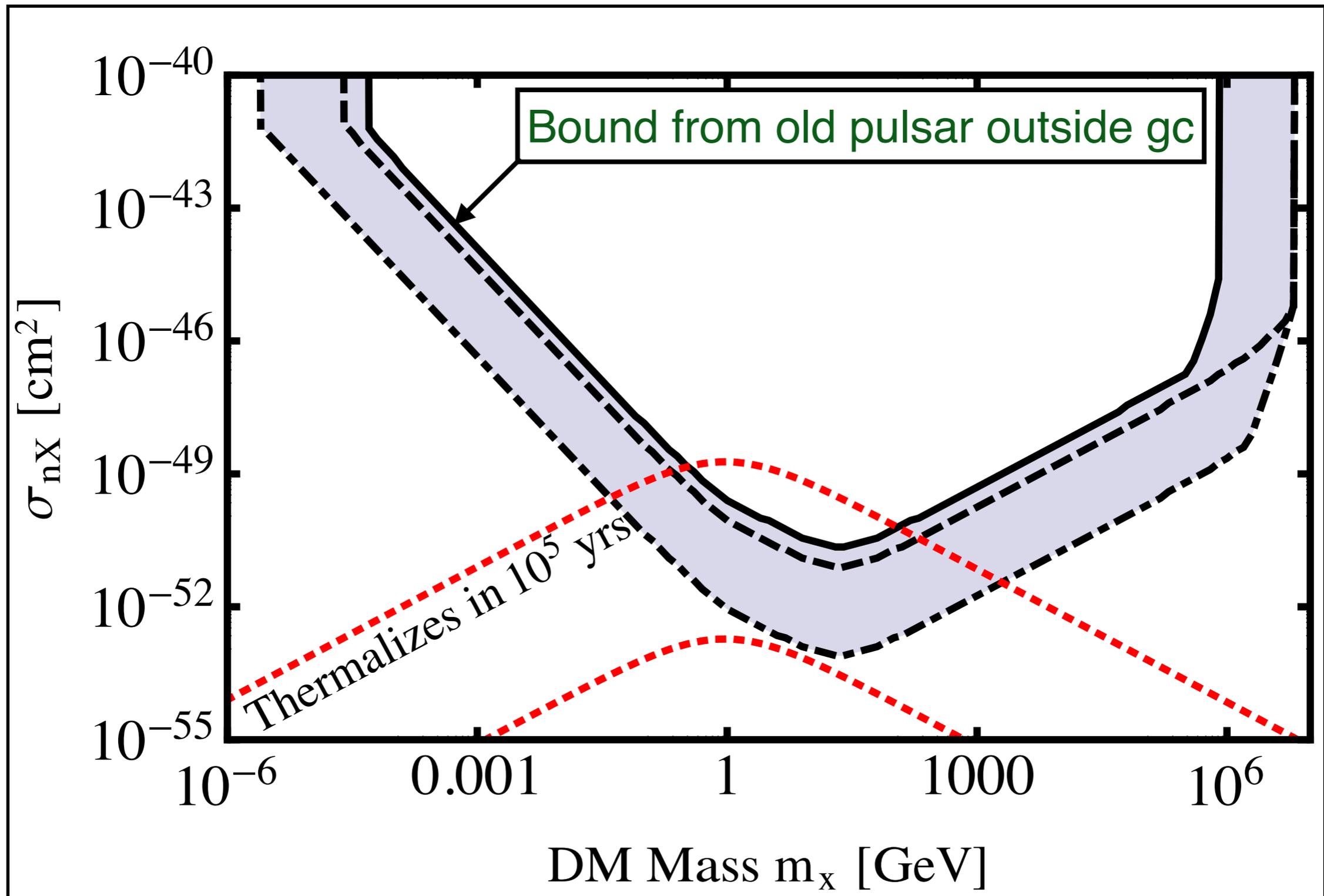
Model our galaxy's dark matter profile using disc and galactic center star velocity measurements (NFW).



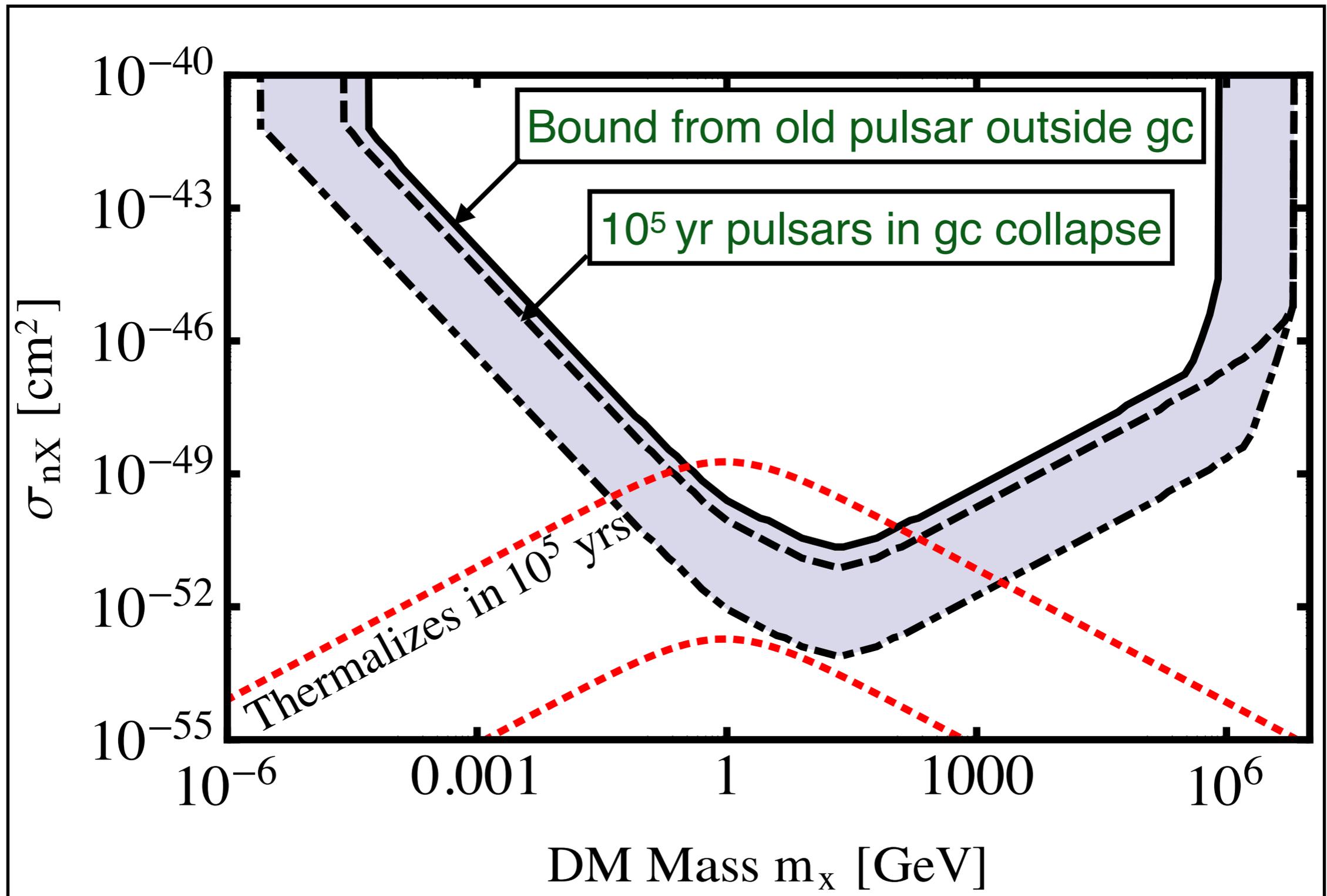
# Bosons Imploding Pulsars



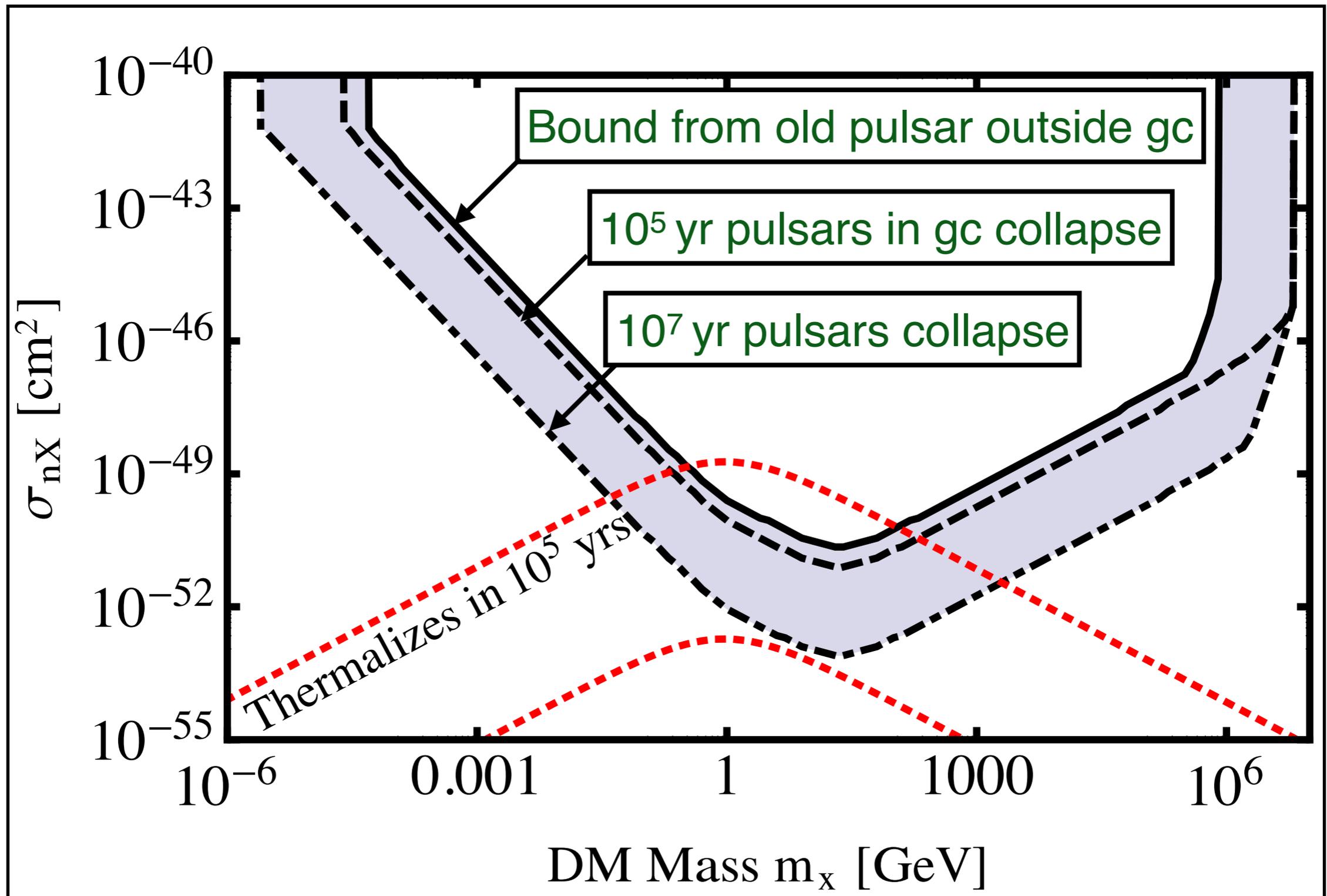
# Bosons Imploding Pulsars



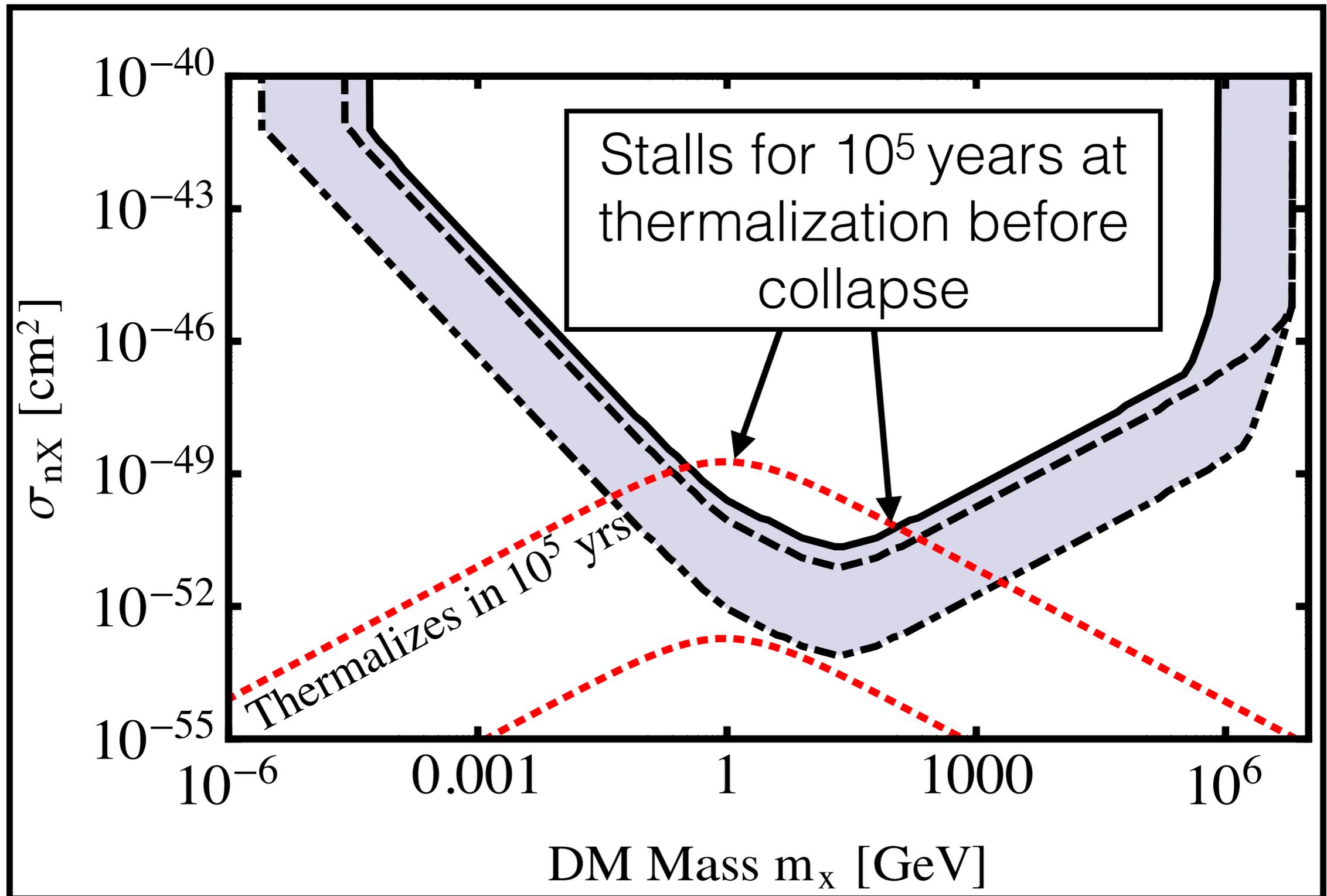
# Bosons Imploding Pulsars



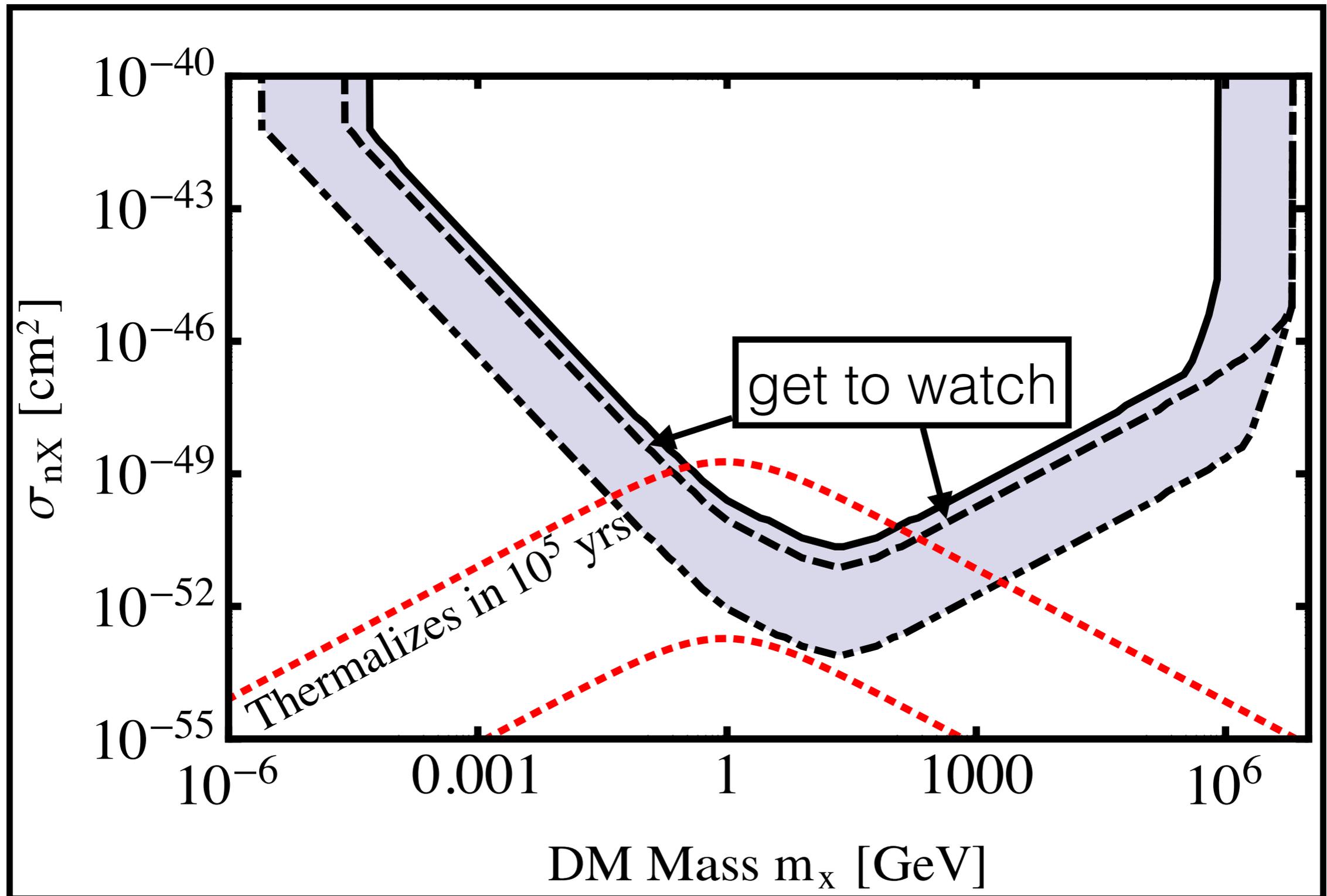
# Bosons Imploding Pulsars



# Bosons Imploding Pulsars

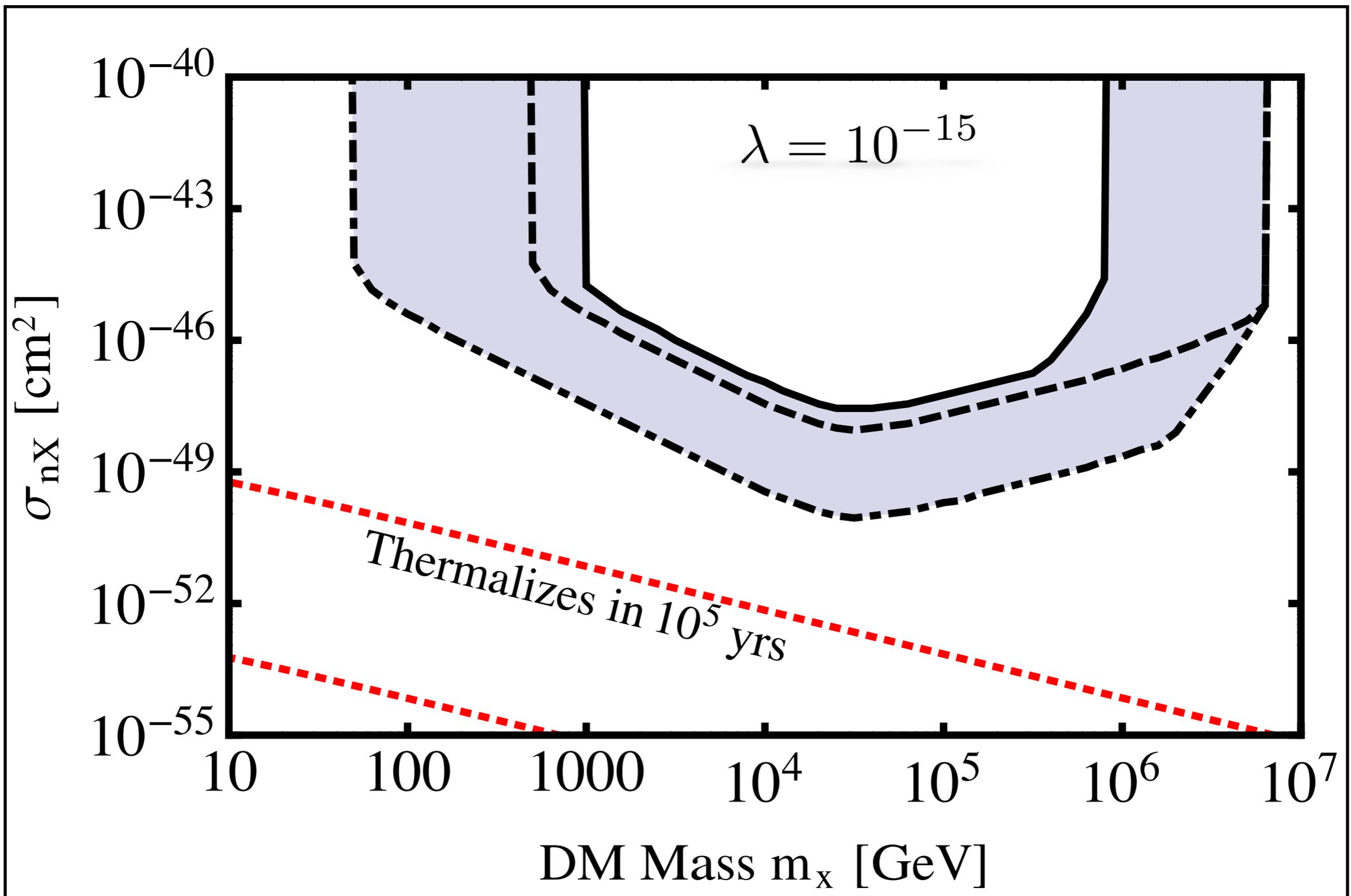


# Bosons Imploding Pulsars

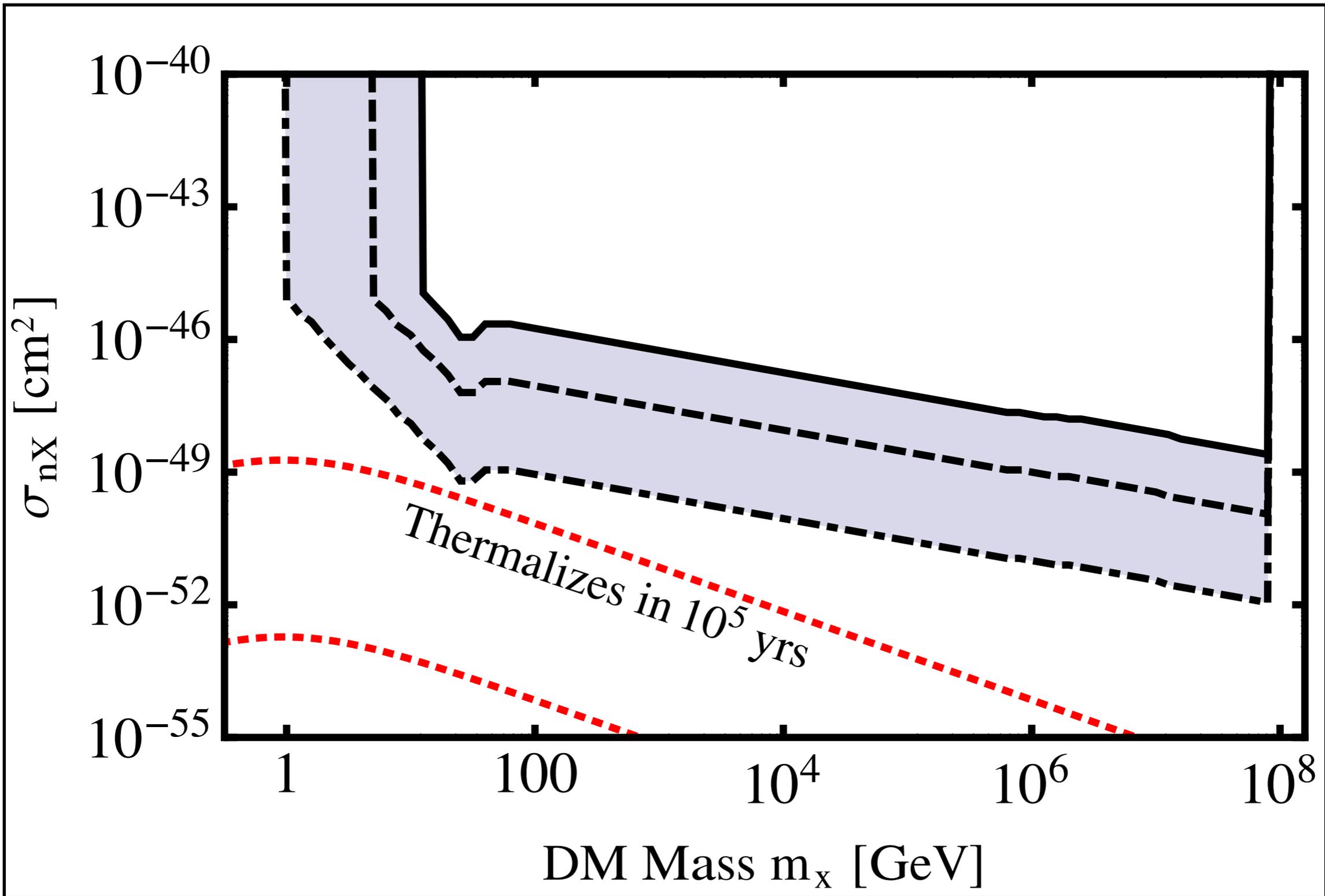


# Bosons Imploding Pulsars

Highly Sensitive to Self-Interactions

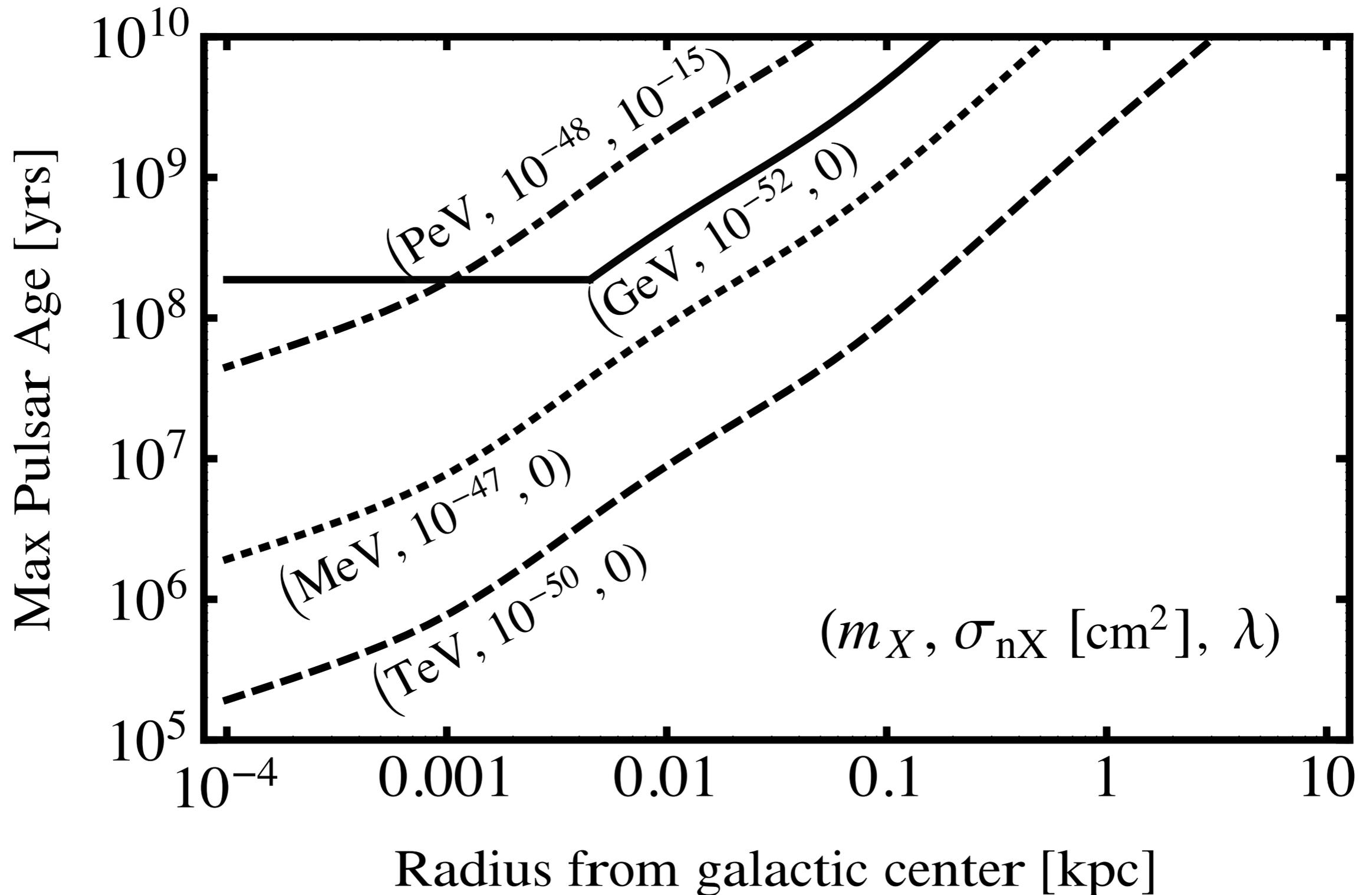


# Fermions Imploding Pulsars

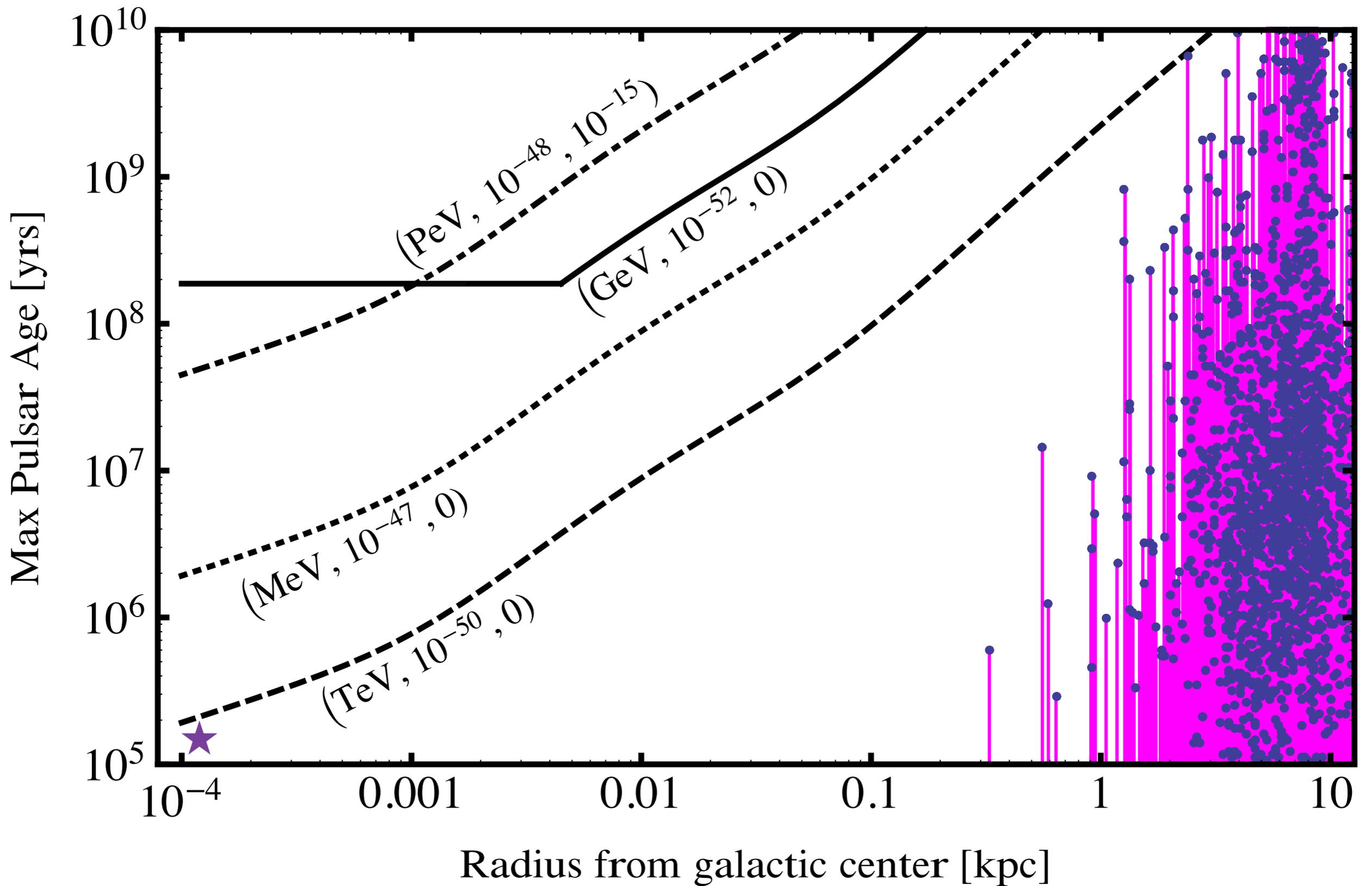


$m_\phi = 10 \text{ MeV}, \quad \alpha = 0.1$

# Prediction of Pulsar Age Curves



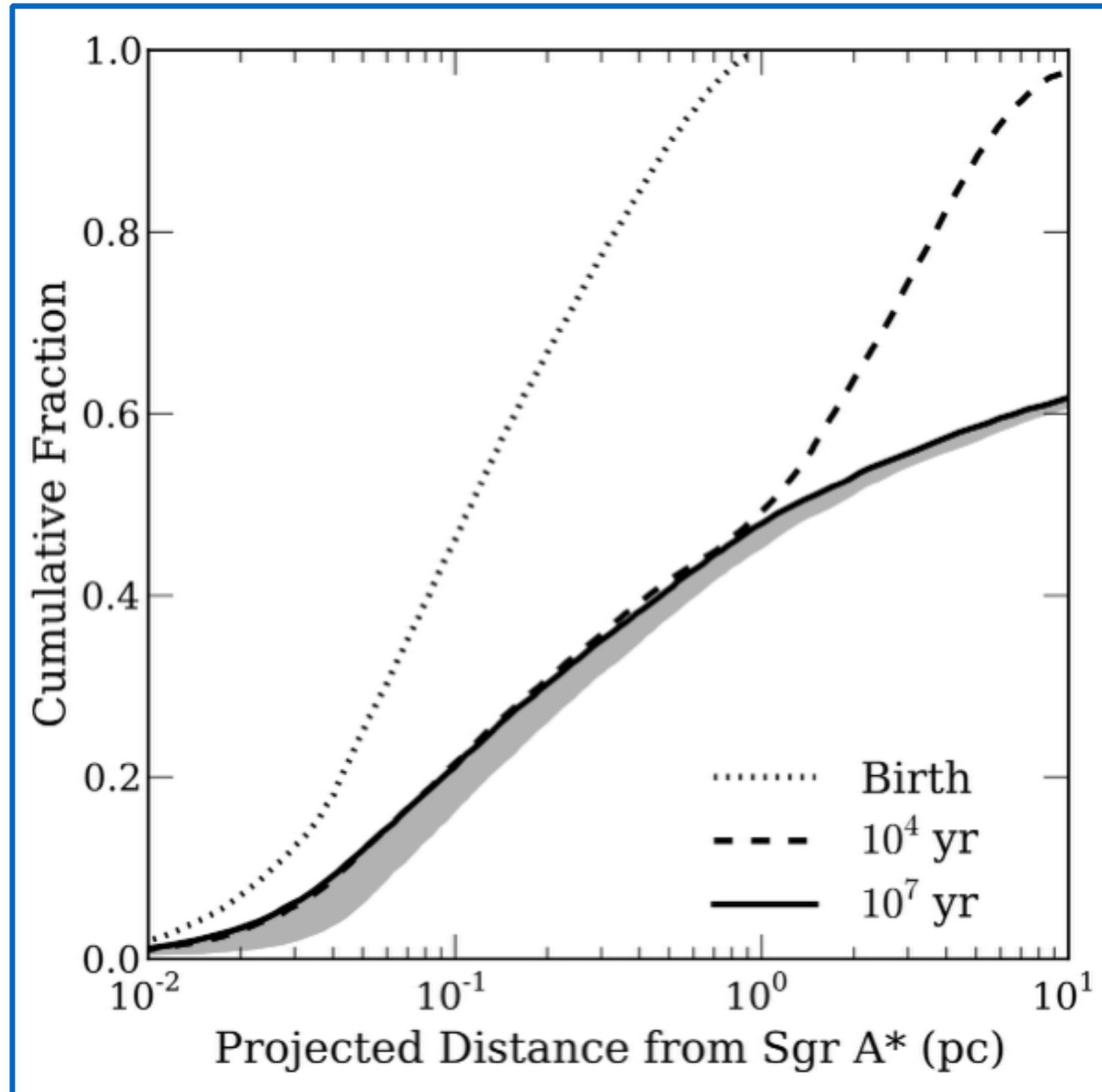
# Prediction of Pulsar Age Curves



# Conclusions

There is an exciting method to distinguish asymmetric from majorana or mixed dark matter. Pulsar age distributions could also be applied to bounds on and searches for multicomponent dark matter.

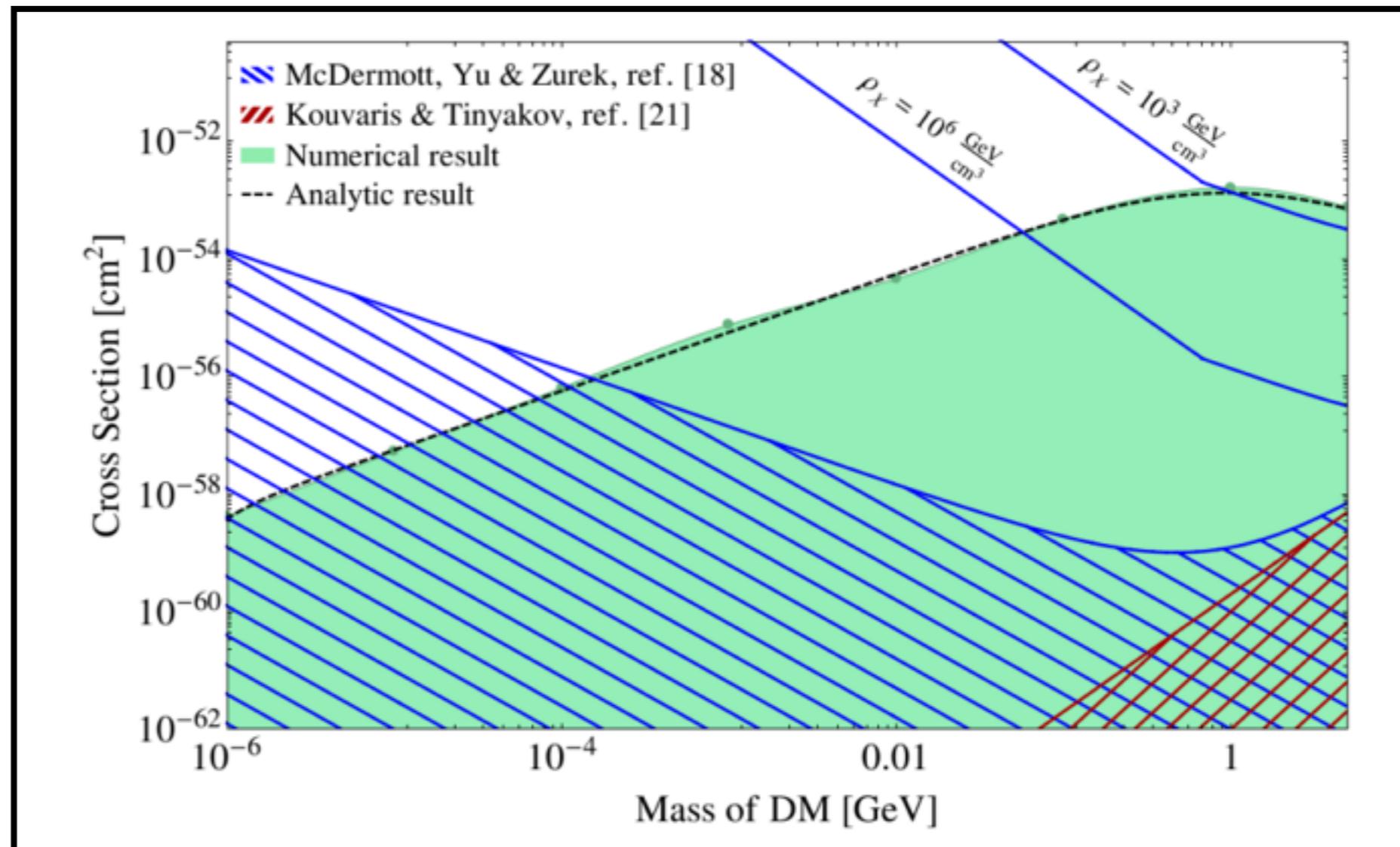
# The missing pulsar problem



# More Precise DM-Neutron Star Thermalization Times

Assumes a 100 GeV heavy mediator and axial/vector nucleon coupling to dark matter.

In a full treatment, thermalization time may be shorter for heavier DM.



# Self-interactions, bullet clusters

- Bullet clusters provide an upper bound on dark matter self-interactions.
- X-ray-emitting ionized IGM slowed by ram pressure as the subcluster slams through a megacuster.

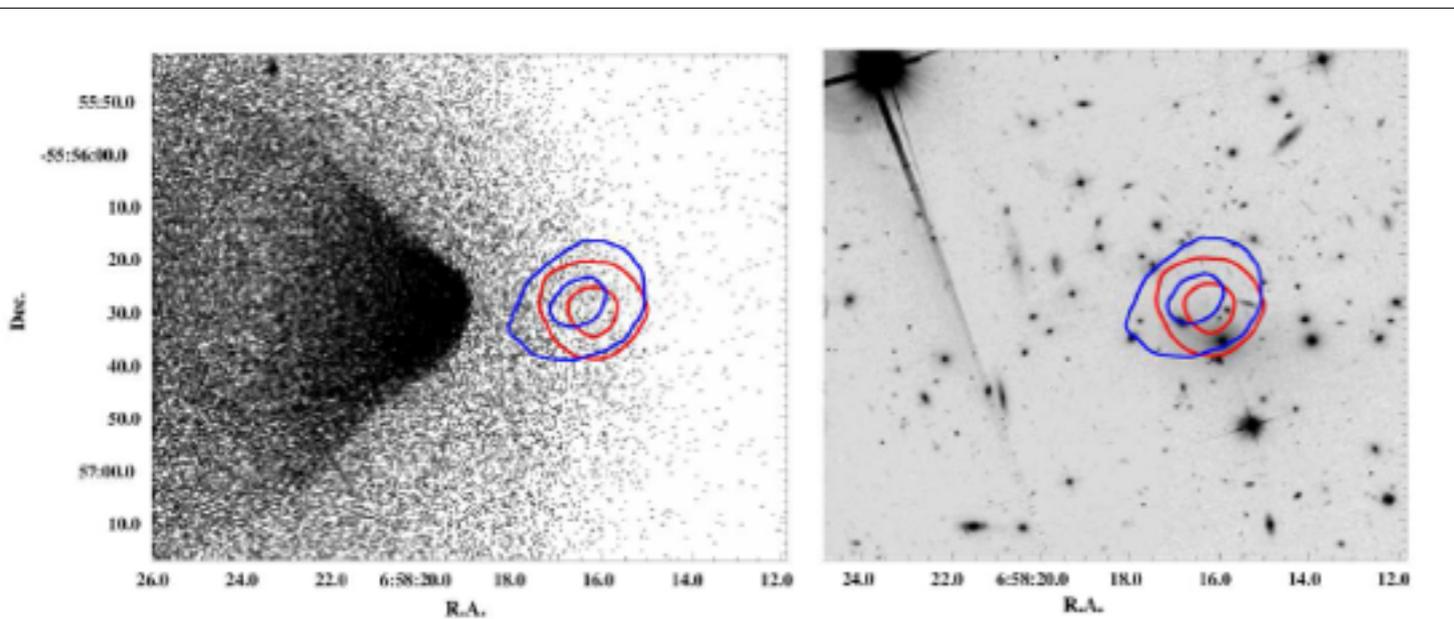


Fig. 2.— Close up of the subcluster bullet region, with the DM (blue) and galaxy (red) centroid error contours overlain. The contours show the 68.3% and 99.7% error regions. The left panel shows the X-ray *Chandra* image, while the right shows the optical *HST* image.

# Self-interactions, bullet clusters

- Bullet clusters provide an upper bound on dark matter self-interactions.
- X-ray-emitting ionized IGM slowed by ram pressure as the subcluster slams through a megacuster.
- **Galaxies**, **DM** not slowed – so compare their separation ( $\Delta x$ ) this will bound self-interacting **DM**!

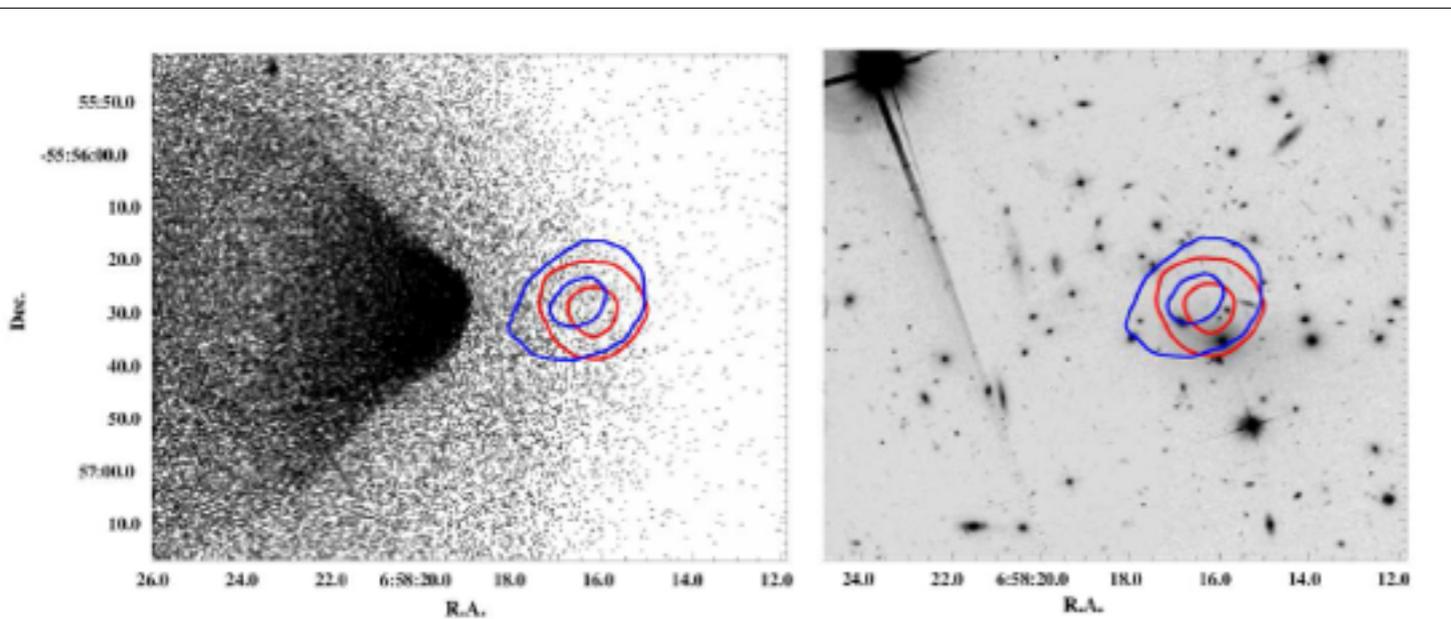


Fig. 2.— Close up of the subcluster bullet region, with the DM (blue) and galaxy (red) centroid error contours overlain. The contours show the 68.3% and 99.7% error regions. The left panel shows the X-ray *Chandra* image, while the right shows the optical *HST* image.

# Self-interactions, bullet clusters

- Bullet clusters provide an upper bound on dark matter self-interactions.
- X-ray-emitting ionized IGM slowed by ram pressure as the subcluster slams through a megacuster.
- **Galaxies**, **DM** not slowed – so compare their separation ( $\Delta x$ ) this will bound self-interacting **DM**!

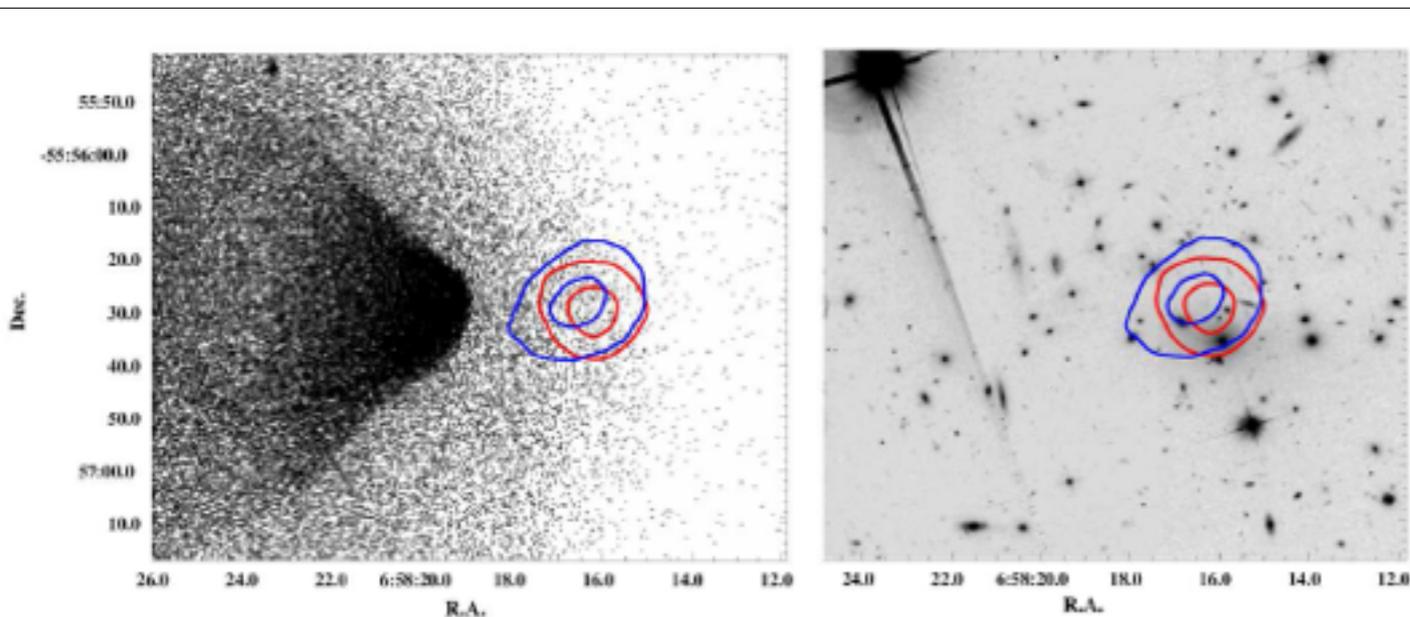


Fig. 2.— Close up of the subcluster bullet region, with the DM (blue) and galaxy (red) centroid error contours overlain. The contours show the 68.3% and 99.7% error regions. The left panel shows the X-ray *Chandra* image, while the right shows the optical *HST* image.

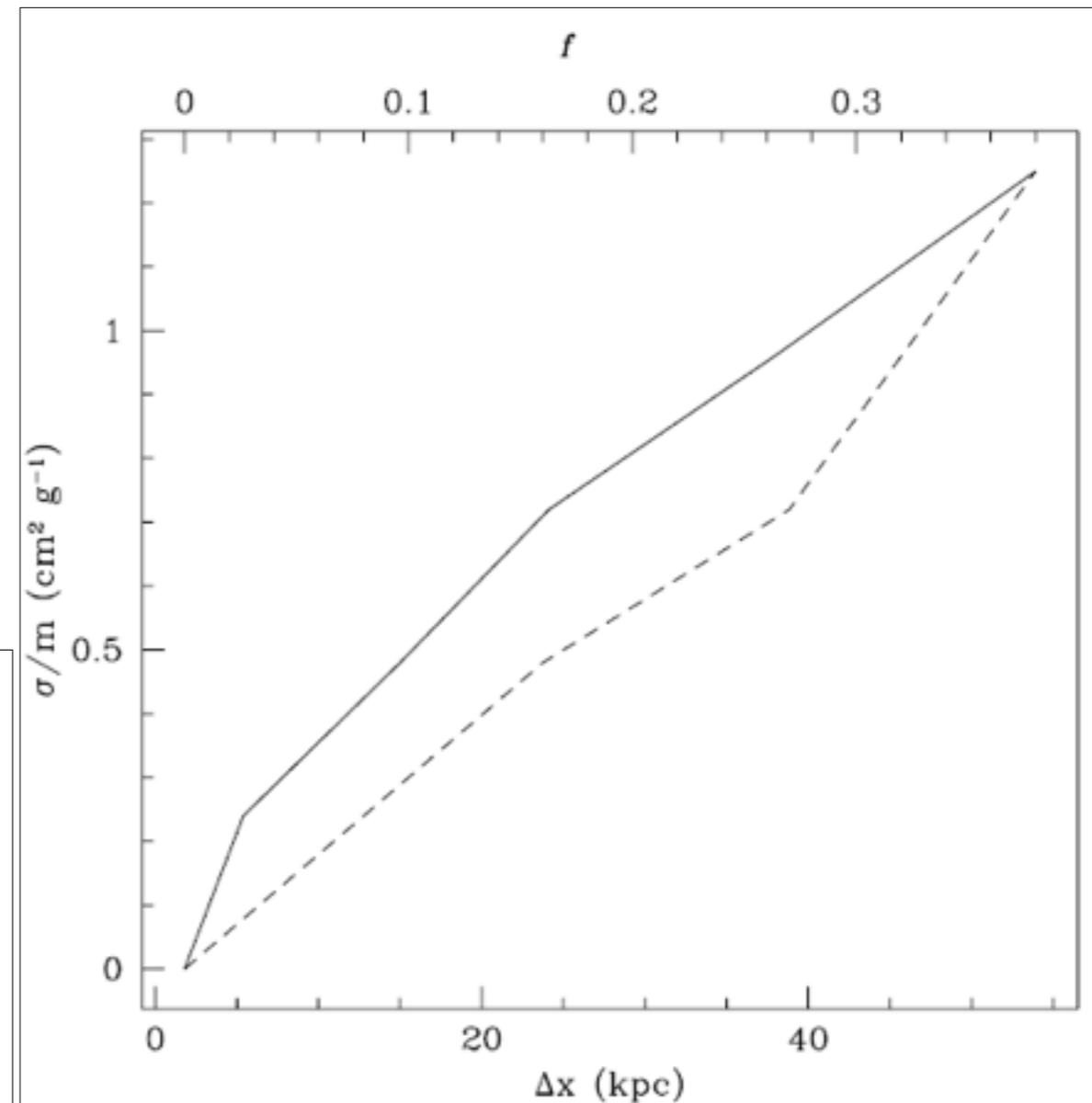


Fig. 5.— The dependence of the subcluster galaxy and total mass centroid offset ( $\Delta x$ , solid line) and the fractional change in the subcluster  $M/L$  ratio ( $f$ , dashed line) on  $\sigma/m$ . Based on the values given in Table 2.