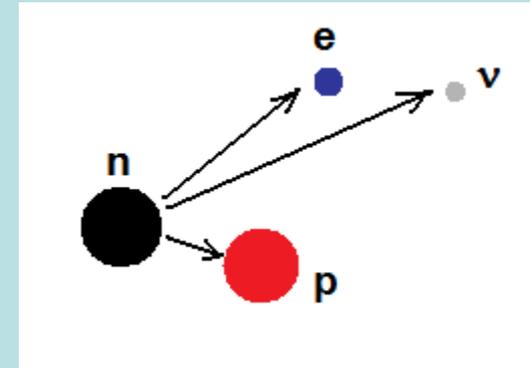
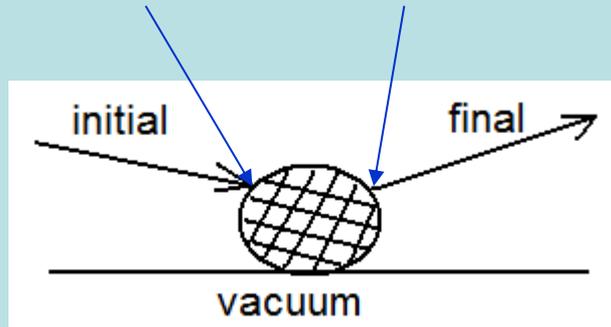


# Extending the frontiers of Particle Physics

Lecture from the course  
« **Cosmoparticle physics** »

# Elementary Particles

- Classical mechanics, Quantum mechanics – an elementary particle is **eternal**
  - The problems with *Energy* and *Angular momentum* conservation in  $\beta$ -decay (1930) – electron and neutrino are **created** in this process
- => **annihilation** and **creation** of particles



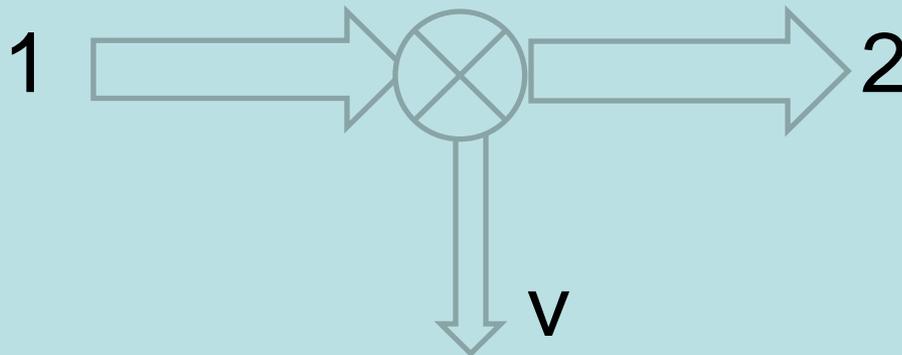
=> New formalism – Quantum field theory

Note: in 1922 A.Friedman discovered non-stationary solutions of Einstein equations for the Universe, proved in the end of 20-s by E.Hubble. Static Universe and its eternal constituents were changed by the modern concepts of particles and expanding Universe nearly simultaneously.

# Particle Interactions

- Charge of an eternal particle converts in QFT in transition in which initial particle annihilates and the final particle is created with emission or absorption of  $\nu$  - quantum of the field of interaction (in QED – photon)

$$e\psi \Rightarrow e|\Psi|^2 \quad \psi \Rightarrow e\Psi_2^* \Psi_1 \quad \psi$$



Elementary act of particle interaction, mediated by boson  $\nu$ , in which initial particle 1 converts into final particle 2

# The Standard Model (SM) of Particle Physics

- Gauge principle (Quantum Electrodynamics (QED)):

$$L = i\bar{\psi}\hat{\partial}\psi + eA_{\mu}J_{\mu} - \frac{1}{4}F_{\mu\nu}F^{\mu\nu}, \quad J_{\mu} = \bar{\psi}\gamma_{\mu}\psi, \quad F_{\mu\nu} = \partial_{\mu}A_{\nu} - \partial_{\nu}A_{\mu}$$

$$\psi \rightarrow \psi \exp(ie\chi(x)) \quad \longleftarrow \quad \text{U(1)-symmetry transformation}$$

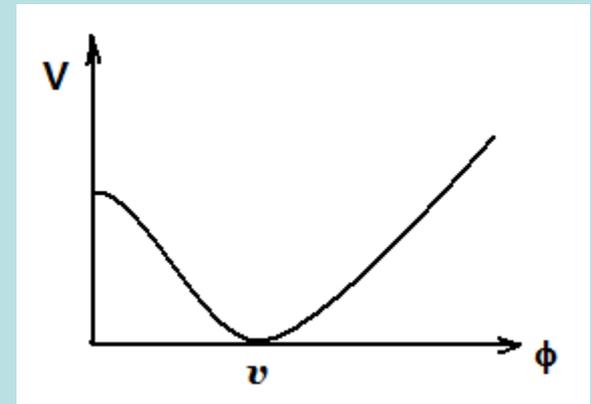
$$A_{\mu} \rightarrow A_{\mu} + \partial_{\mu}\chi$$

- Generalization on *weak* and *strong* interactions:  $\mathbf{SU(3)_c \times SU(2)_L \times U(1)_Y = SM}$
- Higgs mechanism of spontaneous symmetry breaking

$$m_f = \frac{fv}{\sqrt{2}}, \quad m_W = \frac{fv}{2}, \quad m_H = \sqrt{2\lambda}v$$

=>in case of *global* symmetry breaking, a Nambu-Goldstone boson  $\alpha$  appears which interacts as

$$L_{\text{int}} \propto \frac{1}{v} (\partial_{\mu}\alpha)J_{\mu}$$



# Problems of SM (1)

$$SU(3)_c \times SU(2)_L \times U(1)_Y$$

1) Excessive number of parameters:  $g, g', v, \lambda, \{f_{ij}\}, \dots$

2) Problem of number of generations

$$\begin{pmatrix} \nu_e \\ e \end{pmatrix} \begin{pmatrix} \nu_\mu \\ \mu \end{pmatrix} \begin{pmatrix} \nu_\tau \\ \tau \end{pmatrix}$$
$$\begin{pmatrix} u \\ d \end{pmatrix} \begin{pmatrix} c \\ s \end{pmatrix} \begin{pmatrix} t \\ b \end{pmatrix}$$

3) The problem of neutrino mass.

$$m(\nu_e) < 1 \text{ eV} \ll m(e) = 511 \text{ keV}$$

4) CKM mixing matrix  $\begin{pmatrix} d' \\ s' \\ b' \end{pmatrix} = \begin{pmatrix} & & \\ & 3 & \\ & & 3 \end{pmatrix} \begin{pmatrix} d \\ s \\ b \end{pmatrix}$

5) The problem of C-, P- and CP-violation observed (in weak interactions and K-mesons).

# Problems of SM (2)

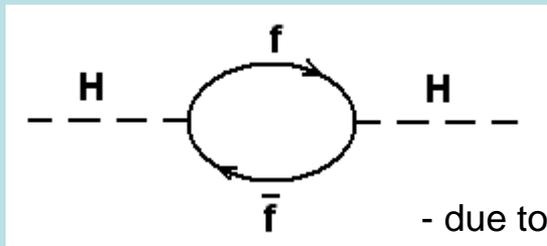
$$SU(3)_c \times SU(2)_L \times U(1)_Y$$

## 6) Problem of CP-violation in QCD

$$\Delta L_{\text{QCD}} = \frac{\alpha_s^2}{16\pi} \theta \cdot \varepsilon_{\alpha\beta\mu\nu} G^{a\alpha\beta} G^{a\mu\nu} \quad \Rightarrow \quad \text{Strong CP-violation}$$

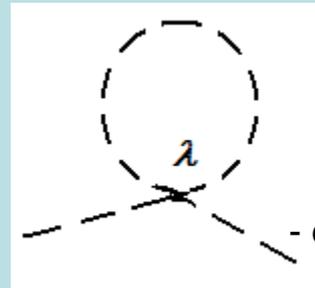
$$\Rightarrow \quad d_n \approx e\theta m_u / m_n^2 \sim \theta \cdot 10^{-16} e \cdot \text{cm} \quad \text{From experiment one has } |\theta| < 10^{-9}$$

## 7) Divergence of Higgs boson mass (due to radiation corrections), and, as consequence, instability of Higgs potential against radiation corrections



$$\sim \Lambda^2$$

- due to  $g_f \phi \bar{f} f$ .



$$\sim \Lambda^2$$

- due to  $\lambda \phi^4$ .

# Other arguments for extension of SM

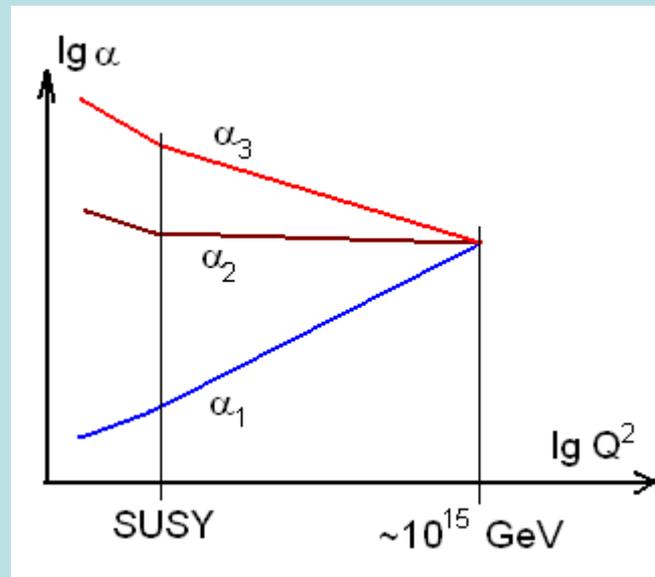
## $SU(3)_c \times SU(2)_L \times U(1)_Y$

Unifying all interactions:

for aesthetical reasons

for reasons of similarity in their descriptions:  $L \sim g A_\mu J_\mu$

convergence of gauge constants (SUSY changes the lines and gives it naturally)



# Trends in extension of SM (1)

“See-saw” mechanism of generation of neutrino mass

$$\Delta L_m = m_D (\bar{\nu}_L \nu_R + \bar{\nu}_R \nu_L) + m_R (\bar{\nu}_R^c \nu_R + \bar{\nu}_R \nu_R^c) + m_L (\bar{\nu}_L^c \nu_L + \bar{\nu}_L \nu_L^c)$$

To avoid problems with SM parameters

at  $m_D \ll m_R$

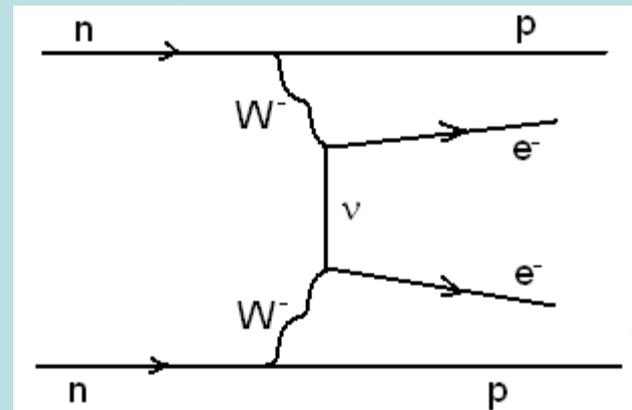
$$\Delta L_m = m_1 (\bar{\nu}_1^c \nu_1 + \bar{\nu}_1 \nu_1^c) + m_2 (\bar{\nu}_2^c \nu_2 + \bar{\nu}_2 \nu_2^c)$$

$$m_1 \approx \frac{m_D^2}{m_R} \ll m_D \sim m_l, \quad m_2 \approx m_R$$

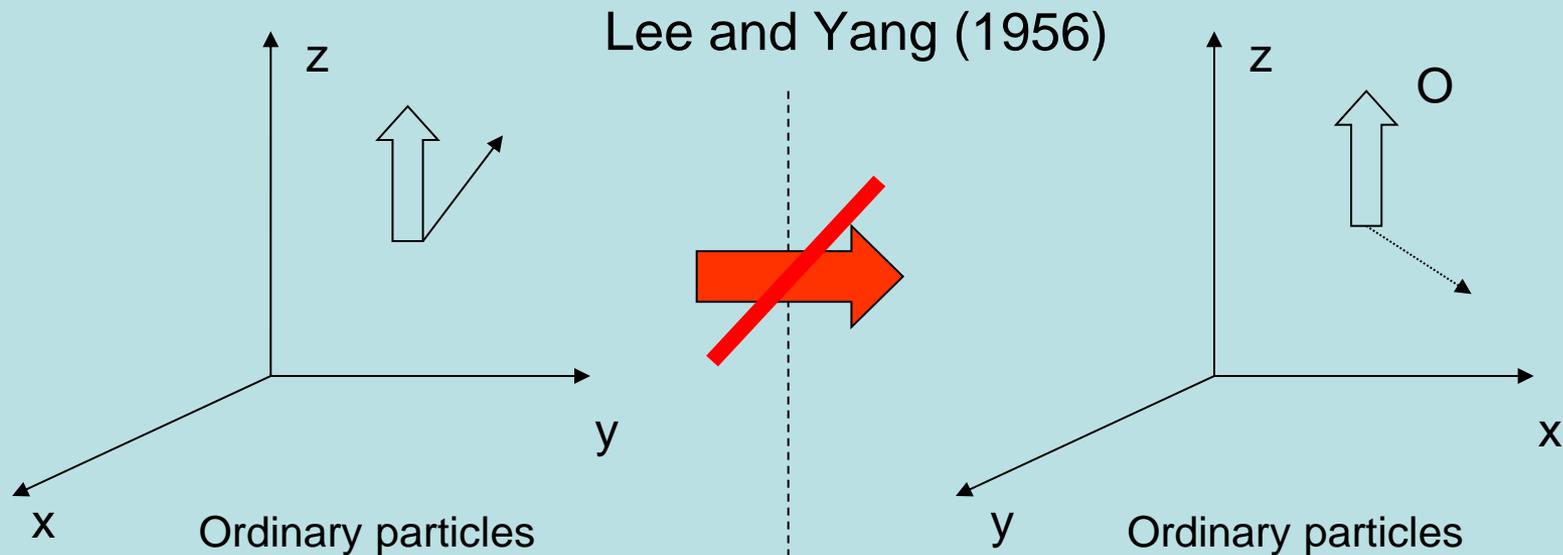
Majorana mass term violates lepton number, which will lead to processes with as  $\Delta L=2$ .

Consequence: neutrinoless double  $\beta$ -decay

Spontaneous breaking of symmetry of lepton charge implies the existence of Nambu-Goldstone boson – Majoron.

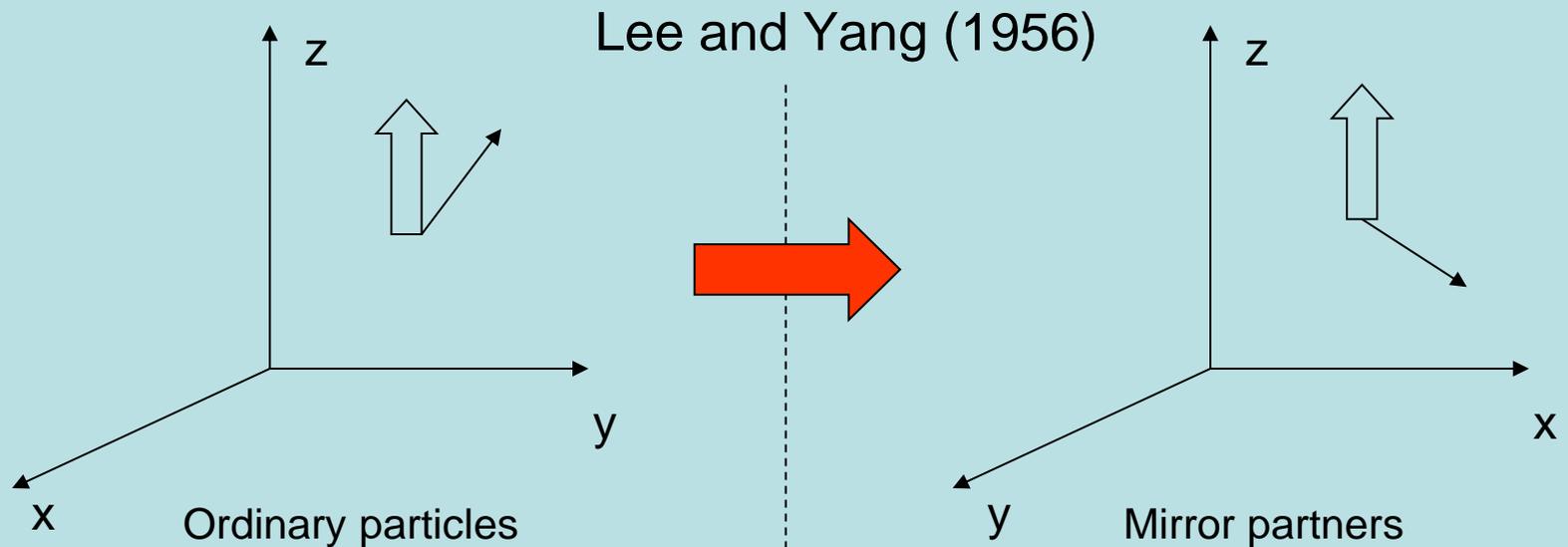


# P-violation



Parity (P) violation means that the process, reflected in mirror does not exist in Nature. It means non-equivalence of left- and right-handed coordinate systems. Beta decay of polarized nucleus, reflected in mirror, does not exist. To restore the equivalence of left- and right-handed coordinate systems P-transformation should be generalized. Together with mirror reflections particles should be changed by their mirror partners.

# Mirror partners?

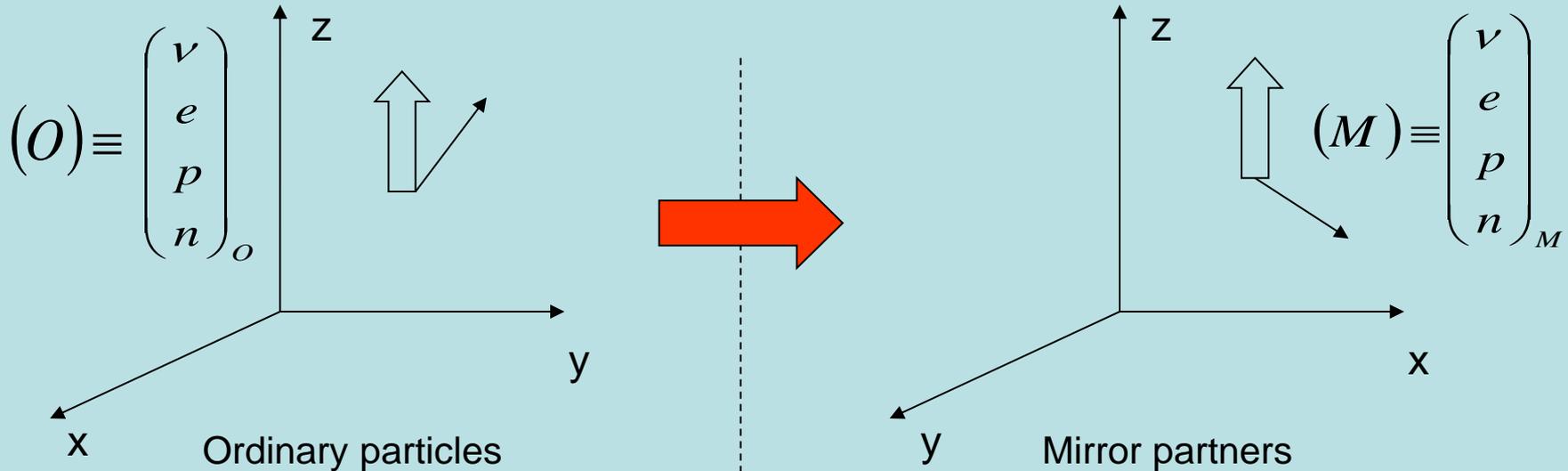


The equivalence between left- and right-handed coordinate systems is restored, if reflection in mirror is accompanied by change of ordinary particles by their mirror partners. Lee, Landau,... (1957) offered an economic solution: CP invariance assumes that antiparticles play the role of mirror partners. Discovery of CP-violation in 1964 put again the question of proper choice for the set of mirror partners.

# Trends in extension of SM (2)

## “Mirror” world

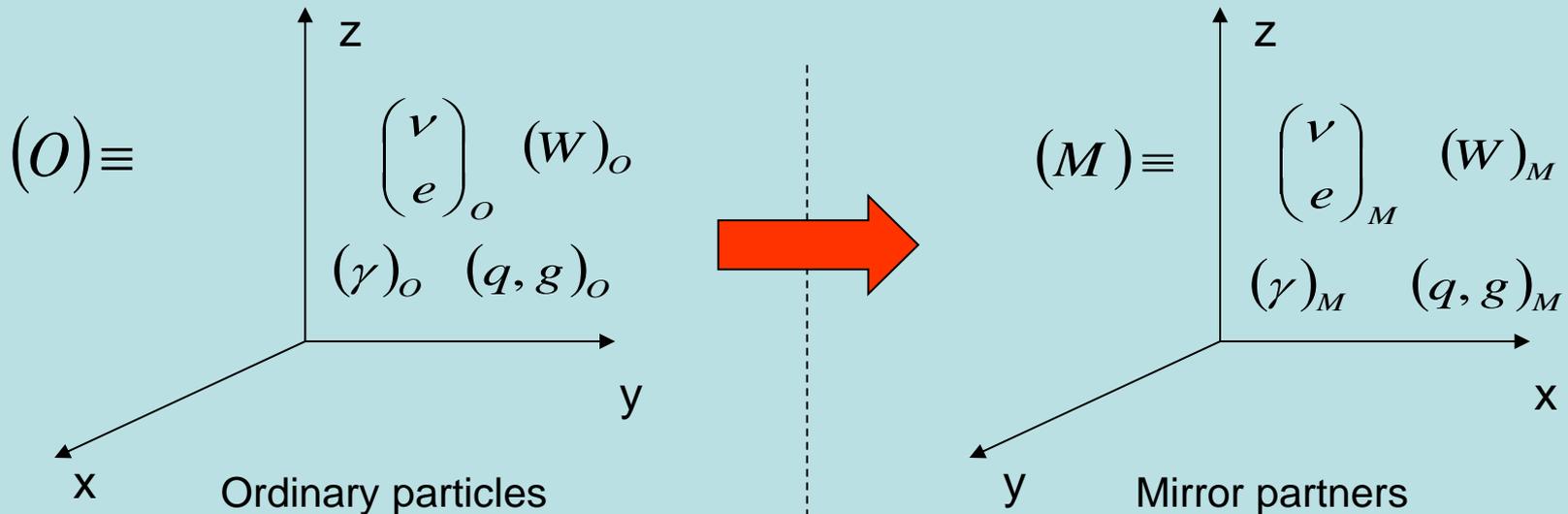
Kobzarev, Okun, Pomeranchuk (1966)



The equivalence between left- and right-handed coordinate systems is restored, if reflection in mirror is accompanied by change of ordinary particles by their mirror partners. Mirror partners are strictly symmetric to ordinary particles. Therefore they can not have ordinary electromagnetic and strong interactions (doubling of atomic levels, or pion states). Successive analysis have shown that (O) and (M) also can not share W and Z boson mediated weak interaction. 11

# Mirror world

Blinnikov, Khlopov (1980, 1982, 1984)



Assume that there is no common interaction between ordinary particles and their mirror partners, except for gravity. All the masses and coupling constants of mirror particles are strictly symmetric to the ordinary ones. The initial conditions are also assumed strictly symmetric.

# Trends in extension of SM (3)

## SU(5) GUT model

$$\text{SU}(5) \supset \text{SU}(3)_c \times \text{SU}(2)_L \times \text{U}(1)_Y$$

- There are 24 gauge bosons, which contain 12 known ones ( $g, Z, W, \gamma$ ) and 12 new ones ( $X, Y$ )
- All known fermions (15, excluding  $\nu_R$ ) are distributed into SU(5) 5-plets and 10-plets.
- $\Lambda_{\text{GUT}} \sim 10^{15} \text{ GeV}$

Example of 5-plet:  $(\bar{d}_{red} \quad \bar{d}_{yellow} \quad \bar{d}_{blue} \quad e^- \quad \nu_e)_L$

Proton decays

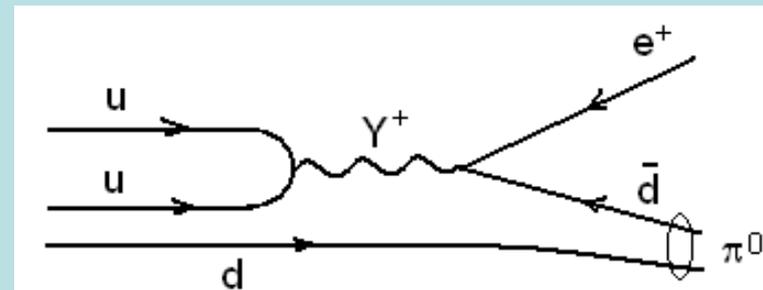
Note:  $\Delta B = \Delta L$

$$\tau_p^{\text{theor}} \sim \frac{\Lambda_{\text{GUT}}^5}{m_p^4} \sim 10^{31} \text{ years}$$

vs

$$\tau_{p \rightarrow e\pi}^{\text{exp}} > 10^{33} \text{ years}$$

It strongly restricts the model.



# Trends in extension of SM (4)

## Symmetry of families and Familons

To resolve **the problem of quark-lepton families** a new global symmetry  $SU(3)_H$  can be introduced. It's spontaneous breaking leads to 8 new scalar fields Nambu-Goldstone bosons - *familons*.

Familons interact with quarks and leptons at energy scale  $v_a$  as

$$L_{aff} = ic_{aff} \frac{m_f}{v_a} \bar{a}f \gamma_5 f$$

Their interactions lead to various transitions between fermions of different families.

In the model of singlet familon there is only 1 such Nambu-Goldstone boson.

# Trends in extension of SM (5)

## Model of Peccei-Quinn

To resolve **the problem of strong CP-violation** a new global symmetry  $U(1)_{PQ}$  with new scalar field is introduced.

Scalar field is associated with spontaneous  $U(1)_{PQ}$  violation at energy scale  $v_a$ , and provides mechanism of dynamical suppression of  $\theta$  (due to its Nambu-Goldstone boson – **axion**).

Axion interacts with quarks as

$$L_{aff} = ic_{aff} \frac{m_f}{v_a} a \bar{f} \gamma_5 f$$

and have potential  $V_{eff} \sim m_u \Lambda_{QCD}^3 \left(1 - \cos(\theta + a / v_a)\right)$

minimizing at  $a = \langle a \rangle = -\theta v_a$


$$\Delta L_{QCD} = \frac{\alpha_s^2}{16\pi} \left( \theta + \frac{a}{v_a} \right) \cdot \epsilon_{\alpha\beta\mu\nu} G^{a\alpha\beta} G^{a\mu\nu} \rightarrow 0$$

# Trends in extension of SM (5)

## Model of Peccei-Quinn: properties of axion

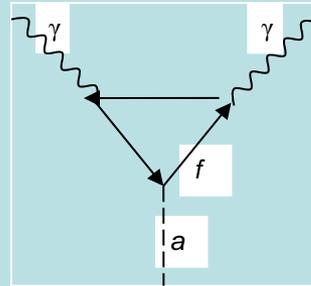
The axion has the mass, which is related with its energy scale

$$m_a \approx 0,6 \text{ meV} \frac{10^{10} \text{ GeV}}{v_a}$$

hep-ph/0002154

It decays into two photons at

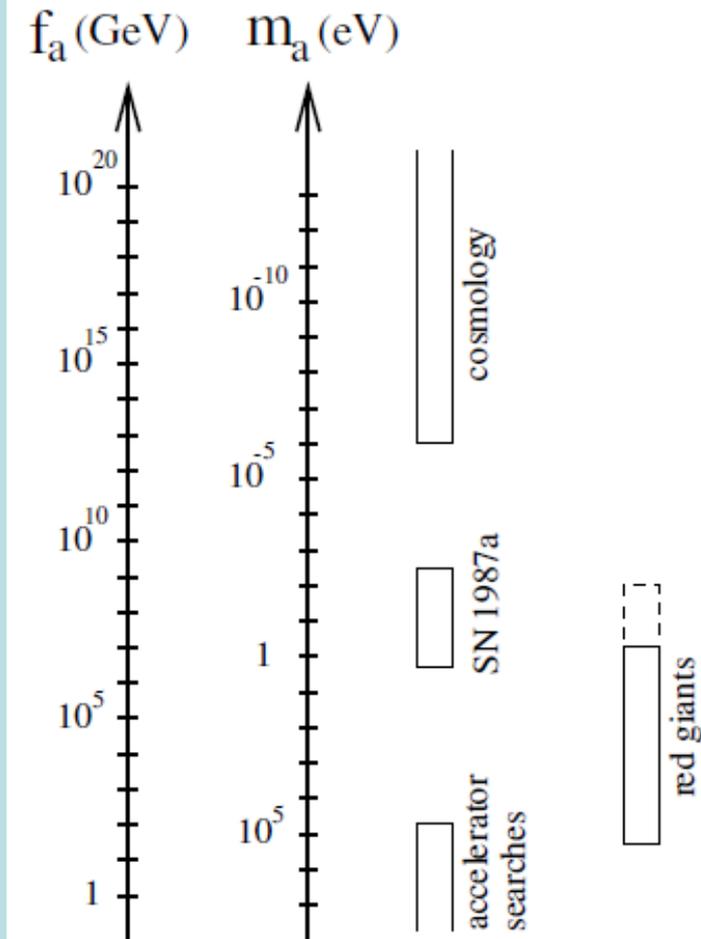
$$\begin{aligned} \tau(a \rightarrow \gamma\gamma) &= \frac{64\pi^3}{c_{a\gamma\gamma}^2 \alpha^2} \frac{v_a^2}{m_a^3} \approx \frac{2 \cdot 10^{17} \text{ years}}{c_{a\gamma\gamma}^2} \left( \frac{\text{eV}}{m_a} \right)^5 \approx \\ &\approx \frac{2 \cdot 10^{33} \text{ years}}{c_{a\gamma\gamma}^2} \left( \frac{v_a}{10^{10} \text{ GeV}} \right)^5. \end{aligned}$$



According existing constraints its mass and energy scale must be

$$0.5 \cdot 10^{-5} \text{ eV} < m_a < 0,5 \cdot 10^{-2} \text{ eV}$$

$$2 \cdot 10^9 \text{ GeV} < v_a < 10^{12} \text{ GeV}$$



# Gauge approach to the problem of quark-lepton families

- There are free families of quarks and leptons with apparent symmetry of their interactions and apparent hierarchy of their mass states.

$$\begin{pmatrix} \nu_e \\ e \\ u \\ d \end{pmatrix} \begin{pmatrix} \nu_\mu \\ \mu \\ c \\ s \end{pmatrix} \begin{pmatrix} \nu_\tau \\ \tau \\ t \\ b \end{pmatrix}$$

Masses of fermions are determined by EW scale.

Masses of H-gauge bosons must be much heavier than this scale, due to the absence of flavor changing neutral currents (FCNC)

# Gauge model of broken family symmetry

- To avoid symmetric mass terms L- and R-handed states of fermions should belong to different representations. It excludes orthogonal and vector-like groups of family symmetry and reduces the choice to SU(3) for 3 generations [8 H gauge bosons].
- Heavy partners F of ordinary fermions f acquire mass by Yukawa coupling with Higgs fields [3 multiplets  $\langle \xi_i \rangle$  and singlet  $h \langle \eta \rangle = \mu$ ]. Mixing of F and f induces SU(3) symmetry breaking pattern in mass pattern of quarks and leptons.
- To compensate anomalies heavy partners N of neutrinos are necessary. It provides the mechanism of neutrino mass.
- Natural choice of Higgs potential leads to additional global U(1) symmetry. It links physics of axion to physics of broken family symmetry.

# Fermion masses

- Dirac see-saw mechanism of mass generation for quarks and charged leptons

$$\begin{array}{ccccccc}
 \longrightarrow & \otimes & \longrightarrow & \otimes & \longrightarrow & \otimes & \longrightarrow \\
 f_L & g\langle\varphi\rangle & F_R & G\langle\xi_i\rangle & F_L & h\langle\eta\rangle & f_R
 \end{array}
 \qquad
 m_f = \frac{h\langle\eta\rangle}{G\langle\xi_i\rangle} g_f \langle\varphi\rangle$$

- Maiorana see-saw mechanism for neutrino mass.

$$\begin{array}{ccccccc}
 \longrightarrow & \otimes & \longrightarrow & \otimes & \longrightarrow & \otimes & \longrightarrow \\
 \nu_L & g\langle\varphi\rangle & N_R & G\langle\xi_i\rangle_{sym} & \bar{N}_L & g\langle\varphi\rangle & \bar{\nu}_R
 \end{array}
 \qquad
 m_\nu = \frac{g_f \langle\varphi\rangle}{G\langle\xi_i\rangle} g_f \langle\varphi\rangle = m_f \frac{g_f \langle\varphi\rangle}{h\langle\eta\rangle}$$

- Pattern of symmetry breaking

$$SU(3) \otimes U(1) \xrightarrow{\langle\xi_0\rangle} SU(2) \otimes U'(1) \xrightarrow{\langle\xi_1\rangle} U''(1) \xrightarrow{\langle\xi_2\rangle \equiv V} I$$

- Results in mass pattern

$$m_i \propto \langle\xi_i\rangle^2$$

# Reduction of number of parameters

Hierarchy of masses is not given by hands, but follows  
from pattern of family symmetry breaking

	$e$	$\mu$	$\tau$		$e$	$\mu$	$\tau$
$e$	$m_e$	$s_{12}$	$s_{13}$	$e$	0	$p$	0
$\mu$	$s_{12}$	$m_\mu$	$s_{23}$	$\mu$	$p$	0	$q$
$\tau$	$s_{13}$	$s_{23}$	$m_\tau$	$\tau$	0	$q$	$r$

3 parameters instead of 9

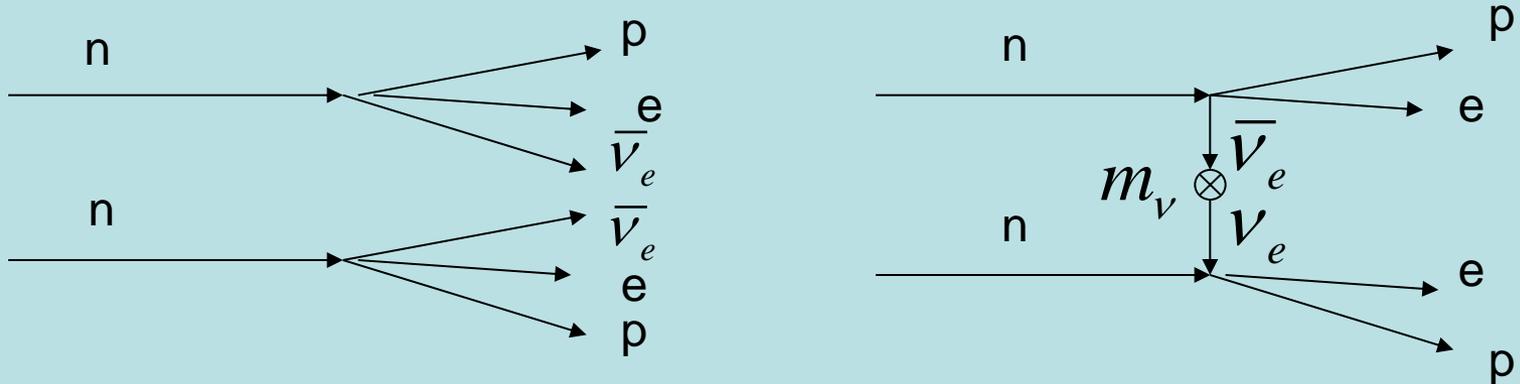
# Archion

- The assumption that Higgs potential contains only terms, which can be generated by radiative effects of gauge and Yukawa couplings excludes the term  $\propto p \xi_i \xi_j \xi_k$ .
- It leads to additional global U(1) symmetry. Breaking of this symmetry results in the existence of Goldstone boson  $\alpha$ .
- This boson shares the properties of Majoron, familon and axion and was called archion.
- Archion couplings are proportional to  $1/V$  and its mass is given by

$$m_\alpha = C \frac{f_\pi}{V} m_\pi$$

# Physical consequences

- FCNC, due to H-bosons  $K \rightarrow \mu e; \bar{D}^0 \leftrightarrow D^0; \dots$
- Mass of neutrino  $m_\nu \propto V^{-1}$ , neutrino oscillations, double neutrinoless beta-decay due to Majorana mass



- Archion decays  $K \rightarrow \pi \alpha; \mu \rightarrow e \alpha; \dots$
- Archion decays of massive neutrinos

$$\nu_H \rightarrow \nu_L \alpha$$

$$\tau = \frac{V^2}{a_{HL}^2 m_H^3} \propto V^5$$

# Trends in extension of SM (6)

## Supersymmetry (SUSY)

All particles with spin  $s$  get supersymmetric partner – particle with spin  $s' = |s - 1/2|$  and all other identical quantum numbers (except the mass in case of broken SUSY).

$$e_{L,R}(s=0) \leftrightarrow \tilde{e}_{L,R}(s=1/2)$$

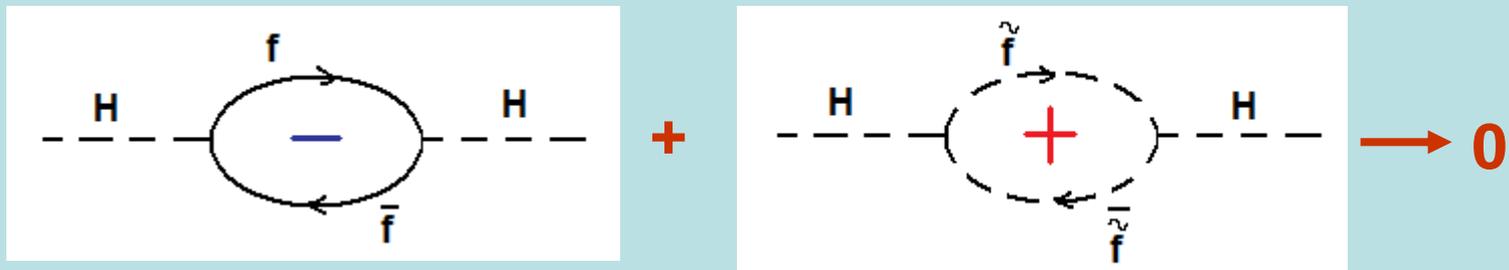
$$\gamma(s=1) \leftrightarrow \tilde{\gamma}(s=1/2)$$

$$\text{or } B(s=1) \leftrightarrow \tilde{B}(s=1/2)$$

....

Photino (bino) is Majorana particle

- 1) SUSY helps to resolve the problem of large radiation correction in Higgs sector.



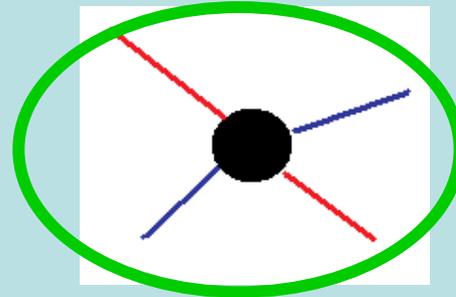
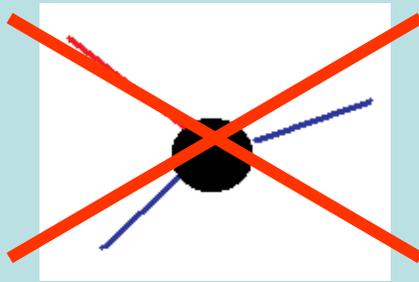
- 2) SUSY improves convergence of the gauge constants (in GUT).
- 3) SUSY may serve as a basic symmetry for construction of quantum gravity.

SUSY must be, evidently, broken. A “hidden” sector is introduced which induces soft SUSY breaking in our sector by means of 1) gravity (**SUGRA**), 2) special gauge fields (**GMSB**). 23

# Trends in extension of SM (6)

## Supersymmetry (SUSY)

To avoid fast decay of proton in GUT, a multiplicative quantum number – *R-parity* is introduced.  $R=+1$  is assigned for all ordinary particles, and  $R=-1$  – for SUSY partners.



**Consequence:** the lightest SUSY particles (**LSP**) is stable.

**Neutralino** is often suggested for a role of LSP. Neutralino is a superposition of four SUSY partners of bino, neutral wino, two neutral higgsino.

$$\chi_{1,2,3,4}^0 = N_{1,2,3,4}^{(B)} \tilde{B} + N_{1,2,3,4}^{(W)} \tilde{W}^3 + N_{1,2,3,4}^{(H_1)} \tilde{H}_1^0 + N_{1,2,3,4}^{(H_2)} \tilde{H}_2^0$$

Also, **gravitino** or **axino** are sometimes suggested for a role of LSP.

# Trends in extension of SM (7)

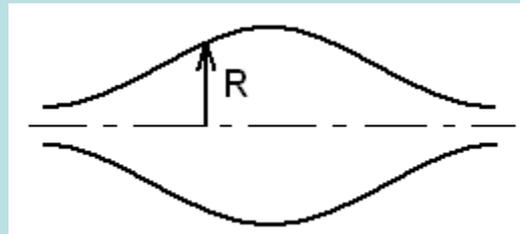
## Kaluza-Klein (KK) models

$$g_{MN} = \begin{pmatrix} g_{\mu\nu} & g_{\mu n} \\ g_{m\nu} & g_{mn} \end{pmatrix}$$

$M, N = 0, \dots, D$   
 $\mu, \nu = 0, 1, \dots, 3$   
 $m, n = 4, \dots, D$   
 $\underbrace{\hspace{10em}}_d$

Describes other interactions

If extra  $d$  dimensions are compactified with  $R = \text{const}$ , then they are manifested in form of the **KK-mass states** for all particles, if  $R \neq \text{const}$  – in form of **interaction**.

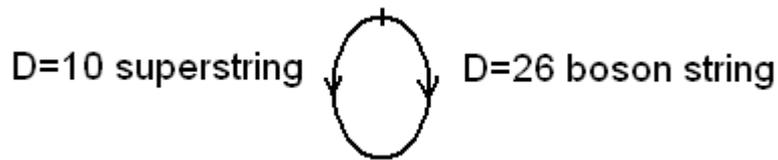


# Trends in extension of SM (9)

## Superstring theory, M-theory

Each species of particles is a certain vibration mode of a string – fundamental element of zero thickness and of Planckian length. Strings may be closed and open.

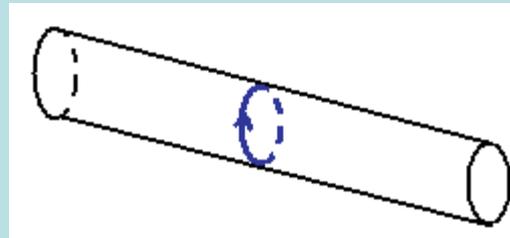
Most popular string model is a so called *heterotic string* (closed) – combines the models of *superstring* and of *bosonic string*.



Fixed gauge groups:  $E_8 \times E_8$ ,  $SO(32)$

Each contains 248 bosons and 248 fermion fields

Possible *consequence*: existence of homotopically stable particle



Superstring U supergravity  $\subset$  M-theory (?)

D=10

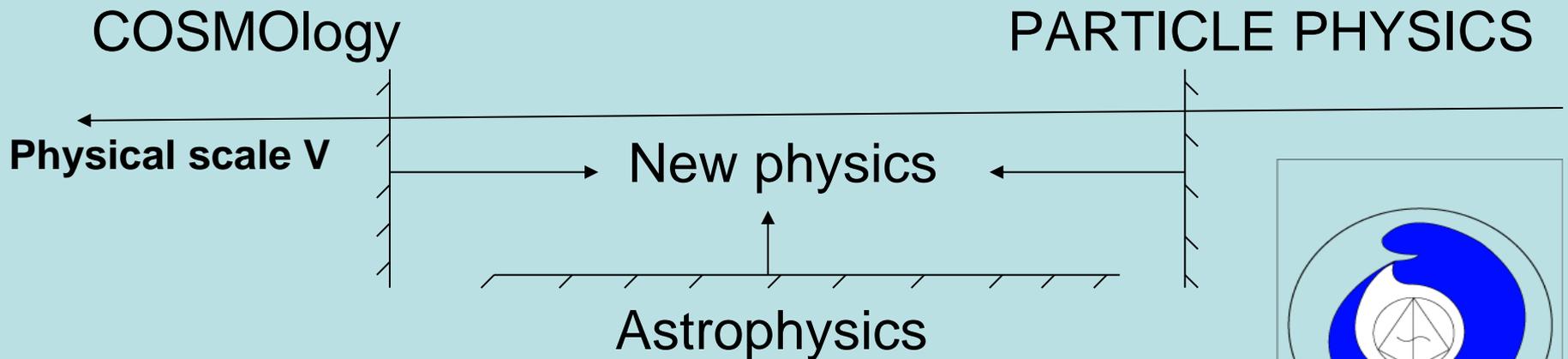
D $\leq$ 11

# Conclusions

- The Standard model contains internal inconsistencies and is not complete. The first experimental evidence of its incompleteness gives nonzero mass of neutrino.
- The problems of the Standard Model imply its extensions for their solution.
- There are only few experimental probes for these extensions, most of which involve physics of very high energy.
- One should extend the set of such probes by astrophysical and cosmological tests.
- We'll see in the successive lectures that the modern cosmology, in its turn, involves physics beyond the Standard model.
- Mutual relationship of problems of frontiers of physics at extreme microscopic and macroscopic scales is the subject of cosmoparticle physics

# Basic ideas of cosmoparticle physics

- Physics beyond the Standard model can be studied in combination of indirect physical, astrophysical and cosmological effects
- New symmetries imply new conserved charges. Strictly conserved charge implies stability of the lightest particle, possessing it.
- New **stable particles** should be present in the Universe. Breaking of new symmetries implies cosmological **phase transitions**. Cosmological and astrophysical constraints are supplementary to direct experimental search and probe the fundamental structure of particle theory at the scale  $V$
- Combination of physical, cosmological and astrophysical effects provide an over-determined system of equations for parameters of particle theory



Extremes of physical knowledge converge in the mystical Uroboros vicious circle of problems, which can be resolved by methods of Cosmoparticle physics

