

# The Cosmological Cheshire Cat

Recent progress to explain dark matter halos

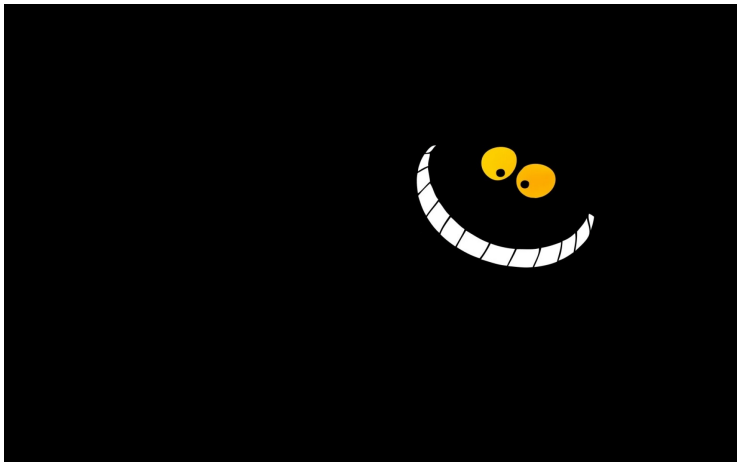
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# Dark matter: only a mischievous grin in our universe

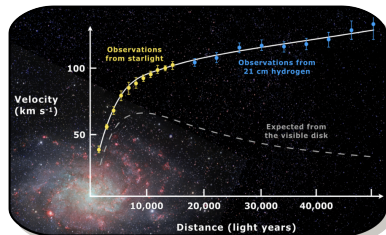


# Observational evidence for dark matter

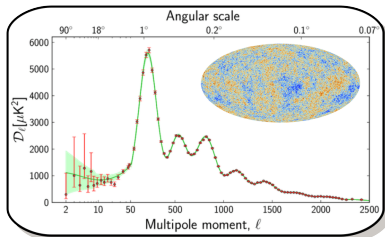
strong gravitational lensing



galaxy rotation curves



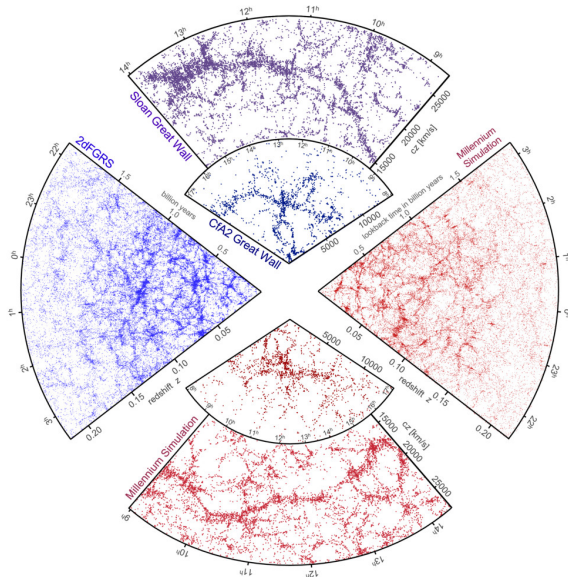
cosmic microwave background



multi-band observations



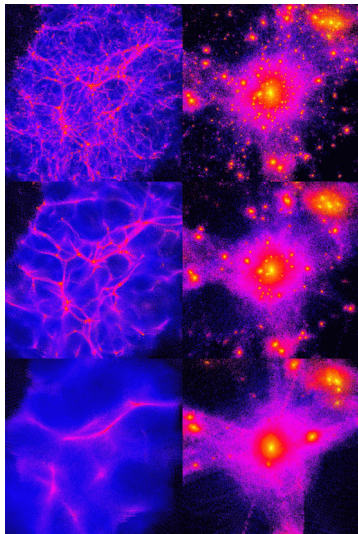
# Dark matter structures in cosmological simulations



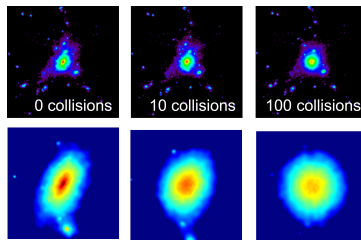


# Constraints on dark matter properties from simulations

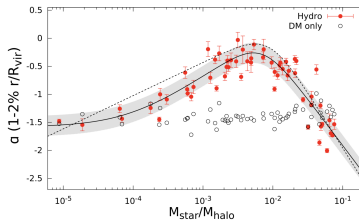
velocity of dark matter



self-interactions of dark matter

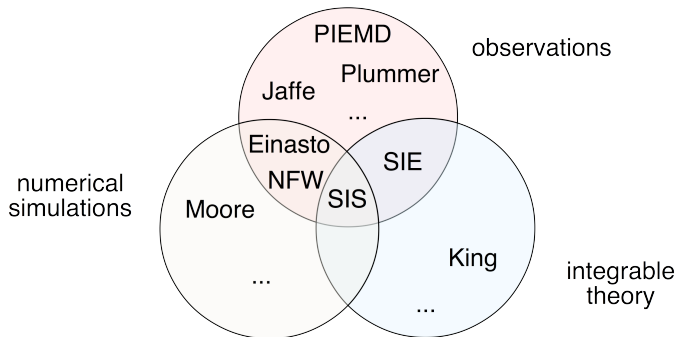


luminous matter feedback



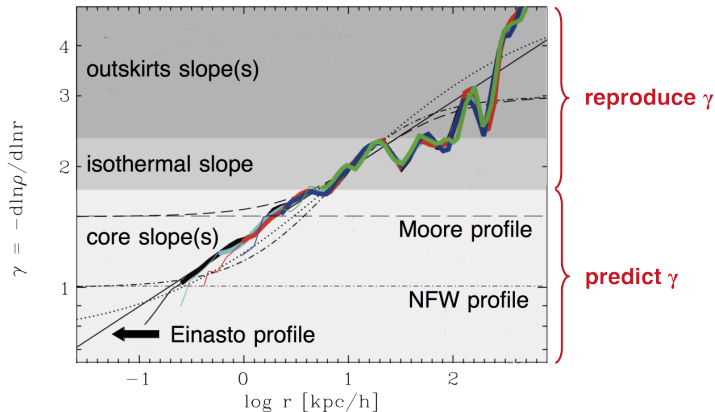
# Motivation: one theory to explain them all

Goal: one model for individual dark matter halo mass densities  $\rho(r)$



# Prerequisites for a testable, evidence-based model

almost universal mass density profiles

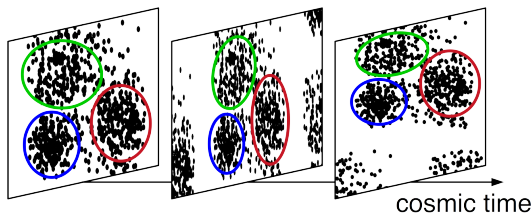


# Prerequisites for an effective, macroscopic model

## scales of interest:

- simulations:  $M \geq \mathcal{O}(10^4 M_\odot)$ ,  $L \geq \mathcal{O}(10 \text{ kpc})$
- observations:  $M \geq \mathcal{O}(10^8 M_\odot)$ ,  $L \geq \mathcal{O}(10 \text{ kpc})$

## decoupling of structure and dynamics



→ individual particle-based halo model at each time

# Prerequisites for a simplified, abstract model

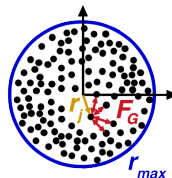
## physical picture

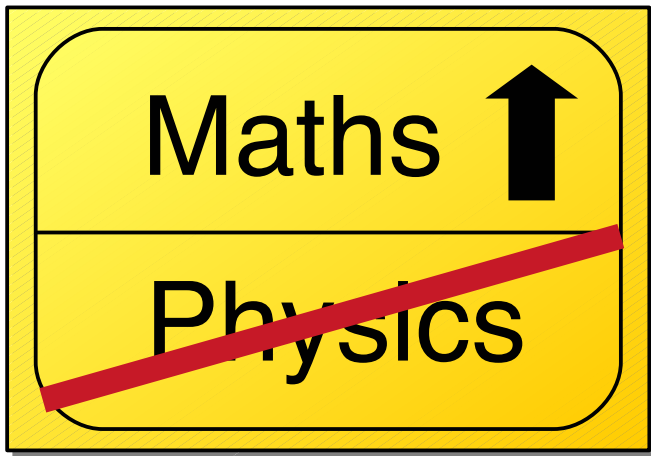
- $n_p$  dark matter particles
- halo = sphere of radius  $r_{\max}$
- Newtonian gravity as the only interaction
- collisionless particles



## mathematical picture

- identical, identically distributed
- finite ensemble props  $(n_p, r_{\max})$
- linear superposition, spatial power-law PDF
- independent particles





# Individual power-law probability densities

For each of the  $n_p$  particles holds:

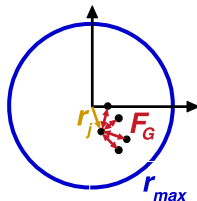
- probability for a particle to be at  $r_j$ :

$$p(r_j) = N \left( \frac{r_j}{r_\sigma} \right)^{-\gamma}, \quad \gamma > 0$$

- normalisation  $N$  given by finite volume:

$$\int_{V_{\max}} dV p(r_j) \stackrel{!}{=} 1$$

$$4\pi N \int_{r_{\min}}^{r_{\max}} dr r^2 \left( \frac{r}{r_\sigma} \right)^{-\gamma} \stackrel{!}{=} 1 \quad \rightarrow \quad N = N(\gamma, r_\sigma, r_{\max}, r_{\min})$$



$$p(r_j) = N(\gamma, r_\sigma, r_{\max}, r_{\min}) \left( \frac{r_j}{r_\sigma} \right)^{-\gamma}$$

## Joining the independent, identically distributed particles:

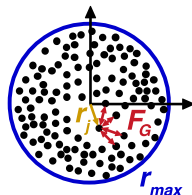
- identically distributed = same PDF for each particle

$$p(r) = N(\gamma, r_\sigma, r_{\max}, r_{\min}) \left( \frac{r}{r_\sigma} \right)^{-\gamma}$$

- independent = joint PDF is product of individual PDFs

$$p_E(r_1, \dots, r_{n_p}) = \prod_{j=1}^{n_p} p(r_j)$$

- (in)distinguishable particles ?



$$p_E(r_1, \dots, r_{n_p}) = \left( \frac{1}{n_p!} \right) \prod_{j=1}^{n_p} p(r_j)$$



# Most likely configuration of the joint ensemble

- joint PDF

$$p_E(r_1, \dots, r_{n_p}) = \left( \frac{1}{n_p!} \right) \prod_{j=1}^{n_p} p(r_j) = \left( \frac{1}{n_p!} \right) \prod_{j=1}^{n_p} N(\gamma) \left( \frac{r_j}{r_\sigma} \right)^{-\gamma}$$

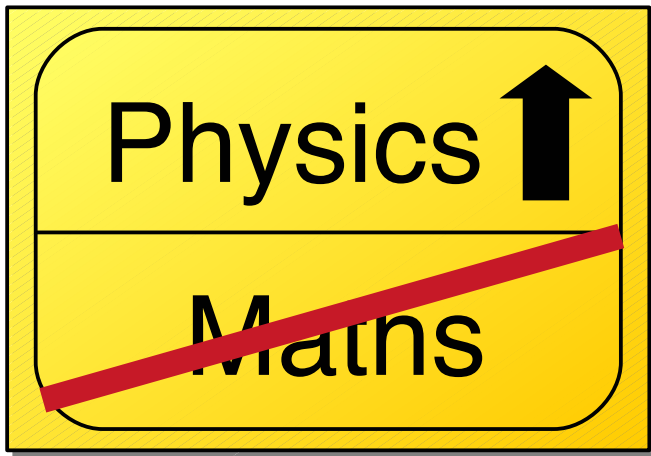
- log-likelihood

$$\begin{aligned} \mathcal{L}(r_1, \dots, r_{n_p}) &= \log(p_E(r_1, \dots, r_{n_p})) = -\log(n_p!) + \sum_{j=1}^{n_p} \log(p(r_j)) \\ &= -\log(n_p!) + n_p \log(N(\gamma)) - \gamma \sum_{j=1}^{n_p} \log\left(\frac{r_j}{r_\sigma}\right) \end{aligned}$$

- most likely configuration to determine  $\gamma$

$$\partial_\gamma \mathcal{L}(r_1, \dots, r_{n_p}) \stackrel{!}{=} 0$$

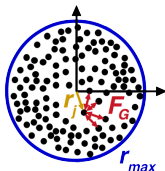
$$\frac{\partial_\gamma N(\gamma, r_\sigma, r_{\max}, r_{\min})}{N(\gamma, r_\sigma, r_{\max}, r_{\min})} - \frac{1}{n_p} \sum_{j=1}^{n_p} \ln\left(\frac{r_j}{r_\sigma}\right) \stackrel{!}{=} 0$$



# Transfer from particle ensemble to mass density

- Density of a dark matter halo

$$\rho(r) = m_p n(r) = m_p n_p p(r) = m_p n_p N(\gamma, r_\sigma, r_{\max}, r_{\min}) \left( \frac{r}{r_\sigma} \right)^{-\gamma}$$



→  $\rho(r)$  from a finite number of spatially bounded particles

→ most likely spatial configuration determines  $\gamma$

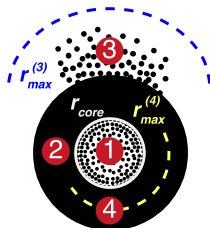
Choose  $r_\sigma \ll r_j$  in

$$\frac{\partial_\gamma N(\gamma, r_\sigma, r_{\max}, r_{\min})}{N(\gamma, r_\sigma, r_{\max}, r_{\min})} - \frac{1}{n_p} \sum_{j=1}^{n_p} \ln \left( \frac{r_j}{r_\sigma} \right) \stackrel{!}{=} 0$$

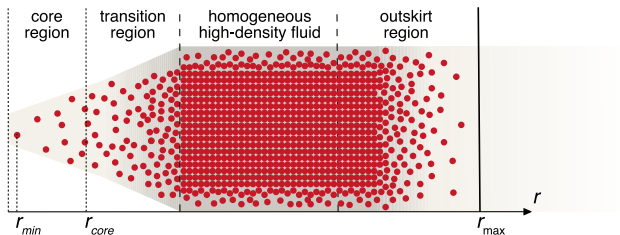
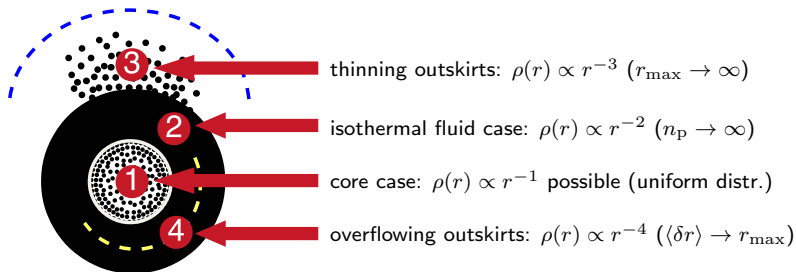
to obtain

$$\gamma = 2 + \frac{n_p}{\sum_{j=1}^{n_p} \ln \left( \frac{r_j}{r_{\max}} \right)}$$

- ❶ core case:  $0 < r_{\min} \leq r_j \leq r_{\text{core}}$
- ❷ isothermal fluid case:  $n_p \gg 1$
- ❸ thinning outskirts:  $r_j \ll r_{\max}$
- ❹ overflowing outskirts:  $r_{\max} \leq r_j \approx 2r_{\max}$



# Limiting approximations



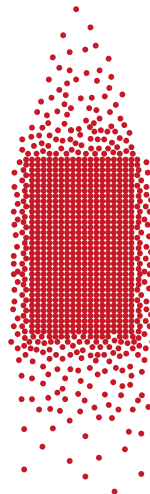
# Summary: dark matter halo profiles from a fundamental theory

- **Our approach so far**

- purely probabilistic (mathematical) approach
  - **separation of structure & dynamics**
  - **independence of cosmic assembly history**
- simple, analytic approximations as solutions
  - **clear understanding of halo shapes**
  - **optimisation of simulations & observations**
- translations of mathematical terms to physical ones exist
  - **connection to thermodynamics**
  - **reduction to necessary prerequisites**

- **Further reading**

- Essay (hon. mention of the Gravity Research Foundation)
  - **arXiv:2005.08975**
- Full paper (in press in General Relativity & Gravitation)
  - **arXiv:2002.00960**



# Thank you for your attention



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**Further information:**

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